

Discretisations of rough stochastic PDEs

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Abstract

We develop a general framework for spatial discretisations of parabolic stochastic PDEs whose solutions are provided in the framework of the theory of regularity structures and which are functions in time. As an application, we show that the dynamical Φ_3^4 model on the dyadic grid converges after renormalisation to its continuous counterpart. This result in particular implies that, as expected, the Φ_3^4 measure is invariant for this equation and that the lifetime of its solutions is almost surely infinite for almost every initial condition.

Contents

1	Introduction	1
2	Regularity structures	6
3	Solutions to parabolic stochastic PDEs	15
4	Discrete models and modelled distributions	20
5	Analysis of discrete stochastic PDEs	31
6	Inhomogeneous Gaussian models	37
7	Convergence of the discrete dynamical Φ_3^4 model	40

1 Introduction

The aim of this article is to develop a general framework for spatial discretisations of the parabolic stochastic PDEs of the form

$$\partial_t u = Au + F(u, \xi) ,$$

where A is an elliptic differential operator, ξ is a rough noise, and F is a non-linear function in u which is affine in ξ . The class of spatial discretisations we work with are of the form

$$\partial_t u^\varepsilon = A^\varepsilon u^\varepsilon + F^\varepsilon(u^\varepsilon, \xi^\varepsilon) ,$$

with the spatial variable taking values in the dyadic grid with mesh size $\varepsilon > 0$, where A^ε , ξ^ε and F^ε are discrete approximations of A , ξ and F respectively.

A particular example prototypical for the class of equations we are interested in is the dynamical Φ^4 model in dimension 3, which can be formally described by the equation

$$\partial_t \Phi = \Delta \Phi + \infty \Phi - \Phi^3 + \xi, \quad \Phi(0, \cdot) = \Phi_0(\cdot), \quad (\Phi_3^4)$$

on the torus $\mathbf{T}^3 \stackrel{\text{def}}{=} (\mathbf{R}/\mathbf{Z})^3$ and for $t \geq 0$, where Δ is the Laplace operator on \mathbf{T}^3 , Φ_0 is some initial data, and ξ is the space-time white noise over $L^2(\mathbf{R} \times \mathbf{T}^3)$, see [PW81].

Here, ∞ denotes an ‘‘infinite constant’’: (Φ_3^4) should be interpreted as the limit of solutions to the equation obtained by mollifying ξ and replacing ∞ by a constant which diverges in a suitable way as the mollifier tends to the identity. It was shown in [Hai14] that this limit exists and is independent of the choice of mollifier. The reason for the appearance of this infinite constant is that solutions are random Schwartz distributions (this is already the case for the linear equation, see [DPZ14]), so that their third power is undefined. The above notation also correctly suggests that solutions to (Φ_3^4) still depend on one parameter, namely the ‘‘finite part’’ of the infinite constant, but this will not be relevant here and we consider this as being fixed from now on.

In two spatial dimensions, a solution theory for (Φ_3^4) was given in [AR91, DPD03], see also [JLM85] for earlier work on a closely related model. In three dimensions, alternative approaches to (Φ_3^4) were recently obtained in [CC13] (via paracontrolled distributions, see [GIP15] for the development of that approach), and in [Kup15] (via renormalisation group techniques à la Wilson).

It is natural to consider finite difference approximations to (Φ_3^4) for a number of reasons. Our main motivation goes back to the seminal article [BFS83], where the authors provide a very clean and relatively compact argument showing that lattice approximations μ_ε to the Φ_3^4 measure are tight as the mesh size goes to 0. Since these measures are invariant for the natural finite difference approximation of (Φ_3^4) , showing that these converge to (Φ_3^4) straightforwardly implies that any accumulation point of μ_ε is invariant for the solutions of (Φ_3^4) . These accumulation points are known to coincide with the Φ_3^4 measure μ [Par75], thus showing that μ is indeed invariant for (Φ_3^4) , as one might expect. Another reason why discretisations of (Φ_3^4) are interesting is because they can be related to the behaviour of Ising-type models under Glauber dynamics near their critical temperature, see [SG73, GRS75]. See also the related result [MW14] where the dynamical Φ_2^4 model is obtained from the Glauber dynamic for a Kac-Ising model in a more direct way, without going through lattice approximations. Similar results are expected to hold in three spatial dimensions, see e.g. the review article [GLP99].

We will henceforth consider discretisations of (Φ_3^4) of the form

$$\frac{d}{dt} \Phi^\varepsilon = \Delta^\varepsilon \Phi^\varepsilon + C_\varepsilon \Phi^\varepsilon - (\Phi^\varepsilon)^3 + \xi^\varepsilon, \quad \Phi^\varepsilon(0, \cdot) = \Phi_0^\varepsilon(\cdot), \quad (\Phi_{3,\varepsilon}^4)$$

on the dyadic discretisation \mathbf{T}_ε^3 of \mathbf{T}^3 with mesh size $\varepsilon = 2^{-N}$ for $N \in \mathbf{N}$, where $\Phi_0^\varepsilon \in \mathbf{R}^{\mathbf{T}_\varepsilon^3}$, Δ^ε is the nearest-neighbour approximation of the Laplacian Δ , and ξ^ε is a spatial discretisation of ξ . We construct these discretisations on a common probability space by setting

$$\xi^\varepsilon(t, x) \stackrel{\text{def}}{=} \varepsilon^{-3} \langle \xi(t, \cdot), \mathbf{1}_{|\cdot - x| \leq \varepsilon/2} \rangle, \quad (t, x) \in \mathbf{R} \times \mathbf{T}_\varepsilon^3, \quad (1.1)$$

where $|x|$ denotes the supremum norm of $x \in \mathbf{R}^3$. Our results are however flexible enough to easily accommodate a variety of different approximations to the noise and the Laplacian.

Existence and uniqueness of global solutions to $(\Phi_{3,\varepsilon}^4)$ for any fixed $\varepsilon > 0$ follows immediately from standard results for SDEs [Has80, IW89]. Our main approximation result is the following, where we take the initial conditions Φ_0^ε to be random variables defined on a common probability space, independent of the noise ξ . (We could of course simply take them deterministic, but this formulation will be how it will then be used in our proof of existence of global solutions.)

Theorem 1.1. *Let ξ be a space-time white noise over $L^2(\mathbf{R} \times \mathbf{T}^3)$ on a probability space $(\Omega, \mathcal{F}, \mathbf{P})$, let $\Phi_0 \in \mathcal{C}^\eta(\mathbf{R}^d)$ almost surely, for some $\eta > -\frac{2}{3}$, and let Φ be the unique maximal solution of (Φ_3^4) on $[0, T^*]$. Let furthermore Δ^ε be the nearest-neighbour approximation of Δ , let $\Phi_0^\varepsilon \in \mathbf{R}^{\mathbf{T}_\varepsilon^3}$ be a random variable on the same probability space, let ξ^ε be given by (1.1), and let Φ^ε be the unique global solution of $(\Phi_{3,\varepsilon}^4)$. If the initial data satisfy almost surely*

$$\lim_{\varepsilon \rightarrow 0} \|\Phi_0; \Phi_0^\varepsilon\|_{\mathcal{C}^\eta}^{(\varepsilon)} = 0,$$

then for every $\alpha < -\frac{1}{2}$ there is a sequence of renormalisation constants $C^{(\varepsilon)} \sim \varepsilon^{-1}$ in $(\Phi_{3,\varepsilon}^4)$ and a sequence of stopping times T_ε satisfying $\lim_{\varepsilon \rightarrow 0} T_\varepsilon = T^*$ in probability such that, for every $\bar{\eta} < \eta \wedge \alpha$, and for any $\delta > 0$ small enough, one has the limit in probability

$$\lim_{\varepsilon \rightarrow 0} \|\Phi; \Phi^\varepsilon\|_{\mathcal{C}_{\bar{\eta}, T_\varepsilon}^{\delta, \alpha}}^{(\varepsilon)} = 0. \quad (1.2)$$

As a corollary of this convergence result and an argument along the lines of [Bou94], we have the following result, where we denote by μ the Φ_3^4 measure on the torus.

Corollary 1.2. *For μ -almost every initial condition Φ_0 and for every $T > 0$, the solution of (Φ_3^4) constructed in [Hai14] belongs to $\mathcal{C}_{\bar{\eta}}^{\delta, \alpha}([0, T], \mathbf{T}^3)$, for δ, α and $\bar{\eta}$ as in (1.2). In particular, this yields a Markov process on $\mathcal{C}^\alpha(\mathbf{T}^3)$ which admits μ as an invariant measure.*

In order to prove this result, we will use regularity structures, as introduced in [Hai14], to obtain uniform bounds (in ε) on solutions to $(\Phi_{3,\varepsilon}^4)$ by describing the right hand side via a type of generalised ‘‘Taylor expansion’’ in the neighbourhood

of any space-time point. The problem of obtaining uniform bounds is then split into the problem of on the one hand obtaining uniform bounds on the objects playing the role of Taylor monomials (these require subtle stochastic cancellations, but are given by explicit formulae), and on the other hand obtaining uniform regularity estimates on the “Taylor coefficients” (these are described implicitly as solutions to a fixed point problem but can be controlled by standard Banach fixed point arguments).

In order to treat the discretised equation $(\Phi_{3,\varepsilon}^4)$, we introduce a discrete analogue to the concept of “model” introduced in [Hai14] and we show that the corresponding “reconstruction map” satisfies uniform bounds analogous to the ones available in the continuous case. One technical difficulty we encounter with this approach is that the set-up is somewhat asymmetric since time is continuous while space is discrete. Instead of considering a fixed model as in [Hai14], we will consider a family of models indexed by the time parameter and satisfying a suitable regularity property. This idea requires some modification of the original theory, in particular of the “abstract integration” operation [Hai14, Sec. 5] and of the corresponding Schauder-type estimates.

As this article was nearing its completion, Zhu and Zhu [ZZ15] independently obtained the convergence of solutions to $(\Phi_{3,\varepsilon}^4)$ to those of (Φ_3^4) using different methods. Additionally, Gubinelli and Perkowski [GP15] recently obtained a similar result for the KPZ equation. One advantage of the approach pursued here is that it is quite systematic and that many of our intermediate results do not specifically refer to the Φ_3^4 model. This lays the foundations of a systematic approximation theory which can in principle be applied to many other singular SPDEs, e.g. stochastic Burgers-type equations [Hai11, HMW14, HM14], the KPZ equation [KPZ86, BG97, Hai13], or the continuous parabolic Anderson model [Hai14, HL15].

Acknowledgements

Structure of the article

In Section 2 we introduce regularity structures and inhomogeneous models (i.e. models which are functions in the time variable). Furthermore, we prove here the key results of the theory in our present framework, namely the reconstruction theorem and the Schauder estimates. In Section 3 we provide a solution theory for a general parabolic stochastic PDE, whose solution is a function in time. Section 4 is devoted to the development of a discrete analogue of inhomogeneous models, which we use in Section 5 to analyse solutions of discretised stochastic equations. In Section 6 we analyse models, built from a Gaussian noise. Finally, in Section 7, we prove Theorem 1.1 and Corollary 1.2.

Notations and conventions

Throughout this article, we will work in \mathbf{R}^{d+1} where d is the dimension of space and 1 is the dimension of time. Moreover, we consider the time-space scaling $\mathfrak{s} = (\mathfrak{s}_0, 1, \dots, 1)$ of \mathbf{R}^{d+1} , where $\mathfrak{s}_0 > 0$ is an integer time scaling and $\mathfrak{s}_i = 1$, for $i = 1, \dots, d$, is the scaling in each spatial direction. We set $|\mathfrak{s}| \stackrel{\text{def}}{=} \sum_{i=0}^d \mathfrak{s}_i$,

denote by $|x|$ the ℓ^∞ -norm of a point $x \in \mathbf{R}^d$, and define $\|z\|_s \stackrel{\text{def}}{=} |t|^{1/s_0} \vee |x|$ to be the s -scaled ℓ^∞ -norm of $z = (t, x) \in \mathbf{R}^{d+1}$. For a multiindex $k \in \mathbf{N}^{d+1}$ we define $|k|_s \stackrel{\text{def}}{=} \sum_{i=0}^d s_i k_i$, and for $k \in \mathbf{N}^d$ with the scaling $(1, \dots, 1)$ we denote the respective norm by $|k|$. (Our natural numbers \mathbf{N} include 0.)

For $r > 0$, we denote by $\mathcal{C}^r(\mathbf{R}^d)$ the usual Hölder space on \mathbf{R}^d , by $\mathcal{C}_0^r(\mathbf{R}^d)$ we denote the space of compactly supported \mathcal{C}^r -functions and by $\mathcal{B}_0^r(\mathbf{R}^d)$ we denote the set of \mathcal{C}^r -functions, compactly supported in $B(0, 1)$ (the unit ball centered at the origin) and with the \mathcal{C}^r -norm bounded by 1.

For $\varphi \in \mathcal{B}_0^r(\mathbf{R}^d)$, $\lambda > 0$ and $x, y \in \mathbf{R}^d$ we define $\varphi_x^\lambda(y) \stackrel{\text{def}}{=} \lambda^{-d} \varphi(\lambda^{-1}(y - x))$. For $\alpha < 0$, we define the space $\mathcal{C}^\alpha(\mathbf{R}^d)$ to consist of $\zeta \in \mathcal{S}'(\mathbf{R}^d)$, belonging to the dual space of the space of \mathcal{C}_0^r -functions, with $r > -\lfloor \alpha \rfloor$, and such that

$$\|\zeta\|_{\mathcal{C}^\alpha} \stackrel{\text{def}}{=} \sup_{\varphi \in \mathcal{B}_0^r} \sup_{x \in \mathbf{R}^d} \sup_{\lambda \in (0, 1]} \lambda^{-\alpha} |\langle \zeta, \varphi_x^\lambda \rangle| < \infty. \quad (1.3)$$

Furthermore, for a function $\mathbf{R} \ni t \mapsto \zeta_t$ we define the operator $\delta^{s,t}$ by

$$\delta^{s,t} \zeta \stackrel{\text{def}}{=} \zeta_t - \zeta_s, \quad (1.4)$$

and for $\delta > 0$, $\eta \leq 0$ and $T > 0$, we define the space $\mathcal{C}_\eta^{\delta, \alpha}([0, T], \mathbf{R}^d)$ to consist of the functions $(0, T] \ni t \mapsto \zeta_t \in \mathcal{C}^\alpha(\mathbf{R}^d)$, such that the following norm is finite

$$\|\zeta\|_{\mathcal{C}_\eta^{\delta, \alpha}} \stackrel{\text{def}}{=} \sup_{t \in (0, T]} |t|_0^{-\eta} \|\zeta_t\|_{\mathcal{C}^\alpha} + \sup_{s \neq t \in (0, T]} |t, s|_0^{-\eta} \frac{\|\delta^{s,t} \zeta\|_{\mathcal{C}^{\alpha-\delta}}}{|t - s|^{\delta/s_0}}, \quad (1.5)$$

where $|t|_0 \stackrel{\text{def}}{=} |t|^{1/s_0} \wedge 1$ and $|t, s|_0 \stackrel{\text{def}}{=} |t|_0 \wedge |s|_0$.

Sometimes we will need to work with space-time distributions with scaling s . In order to describe their regularities, we define, for a test function φ on \mathbf{R}^{d+1} , for $\lambda > 0$ and $z, \bar{z} \in \mathbf{R}^{d+1}$,

$$\varphi_z^{\lambda, s}(\bar{z}) \stackrel{\text{def}}{=} \lambda^{-|s|} \varphi(\lambda^{-s_0}(\bar{z}_0 - z_0), \lambda^{-1}(\bar{z}_1 - z_1), \dots, \lambda^{-1}(\bar{z}_d - z_d)), \quad (1.6)$$

and we define the space $\mathcal{C}_s^\alpha(\mathbf{R}^{d+1})$ similarly to $\mathcal{C}^\alpha(\mathbf{R}^d)$, but using the scaled functions (1.6) in (1.3).

In this article we will also work with discrete functions $\zeta^\varepsilon \in \mathbf{R}^{\Lambda_\varepsilon^d}$ on the dyadic grid $\Lambda_\varepsilon^d \subset \mathbf{R}^d$ with the mesh size $\varepsilon = 2^{-N}$ for $N \in \mathbf{N}$. In order to compare them with their continuous counterparts $\zeta \in \mathcal{C}^\alpha(\mathbf{R}^d)$ with $\alpha \leq 0$, we introduce the following “distance”

$$\|\zeta; \zeta^\varepsilon\|_{\mathcal{C}^\alpha}^{(\varepsilon)} \stackrel{\text{def}}{=} \sup_{\varphi \in \mathcal{B}_0^r} \sup_{x \in \mathbf{R}^d} \sup_{\lambda \in [\varepsilon, 1]} \lambda^{-\alpha} |\langle \zeta, \varphi_x^\lambda \rangle - \langle \zeta^\varepsilon, \varphi_x^\lambda \rangle_\varepsilon|,$$

where $\langle \cdot, \cdot \rangle_\varepsilon$ is the discrete analogue of the duality pairing on the grid, i.e.

$$\langle \zeta^\varepsilon, \varphi_x^\lambda \rangle_\varepsilon \stackrel{\text{def}}{=} \int_{\Lambda_\varepsilon^d} \zeta^\varepsilon(y) \varphi_x^\lambda(y) dy \stackrel{\text{def}}{=} \varepsilon^d \sum_{y \in \Lambda_\varepsilon^d} \zeta^\varepsilon(y) \varphi_x^\lambda(y). \quad (1.7)$$

For space-time distributions / functions ζ and ζ^ε , for $\delta > 0$ and $\eta \leq 0$, we define

$$\|\zeta; \zeta^\varepsilon\|_{\mathcal{C}_{\eta,T}^{\delta,\alpha}}^{(\varepsilon)} \stackrel{\text{def}}{=} \sup_{t \in (0,T]} |t|_\varepsilon^{-\eta} \|\zeta_t; \zeta_t^\varepsilon\|_{\mathcal{C}^\alpha} + \sup_{s \neq t \in (0,T]} |s, t|_\varepsilon^{-\eta} \frac{\|\delta^{s,t} \zeta; \delta^{s,t} \zeta^\varepsilon\|_{\mathcal{C}^{\alpha-\delta}}}{(|t-s|^{1/s_0} \vee \varepsilon)^\delta}, \quad (1.8)$$

where $|t|_\varepsilon \stackrel{\text{def}}{=} |t|_0 \vee \varepsilon$ and $|s, t|_\varepsilon \stackrel{\text{def}}{=} |s|_\varepsilon \wedge |t|_\varepsilon$. Furthermore, we define the norm $\|\zeta^\varepsilon\|_{\mathcal{C}_{\eta,T}^{\delta,\alpha}}^{(\varepsilon)}$ in the same way as in (1.3) and (1.5), but using the discrete pairing (1.7).

Finally, we denote by \star and \star_ε the convolutions on \mathbf{R}^{d+1} and $\mathbf{R} \times \Lambda_\varepsilon^d$ respectively, and by $x \lesssim y$ we mean that there exists a constant C independent of the relevant quantities such that $x \leq Cy$.

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2 Regularity structures

In this section we recall the definition of a regularity structure and we introduce the inhomogeneous models used in this article, which are maps from \mathbf{R} (the time coordinate) to the usual space of models as in [Hai14, Def. 2.17], endowed with a norm enforcing some amount of time regularity. Furthermore, we define inhomogeneous modelled distributions and prove the respective reconstruction theorem and Schauder estimates. Throughout this section, we work with the scaling $\mathfrak{s} = (\mathfrak{s}_0, 1, \dots, 1)$ of \mathbf{R}^{d+1} , but all our results can easily be generalised to any non-Euclidean scaling in space, similarly to [Hai14].

2.1 Regularity structures and inhomogeneous models

The purpose of regularity structures, introduced in [Hai14] and motivated by [Lyo98, Gub04], is to generalise Taylor expansions using essentially arbitrary functions/distributions instead of polynomials. The precise definition is as follows.

Definition 2.1. A *regularity structure* $\mathcal{T} = (\mathcal{T}, \mathcal{G})$ consists of two objects:

- A *model space* \mathcal{T} , which is a graded vector space $\mathcal{T} = \bigoplus_{\alpha \in \mathcal{A}} \mathcal{T}_\alpha$, where each \mathcal{T}_α is a (finite dimensional in our case) Banach space and $\mathcal{A} \subset \mathbf{R}$ is a finite set of “homogeneities”.
- A *structure group* \mathcal{G} of linear transformations of \mathcal{T} , such that for every $\Gamma \in \mathcal{G}$, every $\alpha \in \mathcal{A}$ and every $\tau \in \mathcal{T}_\alpha$ one has $\Gamma\tau - \tau \in \mathcal{T}_{<\alpha}$, with $\mathcal{T}_{<\alpha} \stackrel{\text{def}}{=} \bigoplus_{\beta < \alpha} \mathcal{T}_\beta$.

In [Hai14, Def. 2.1], the set \mathcal{A} was only assumed to be locally finite and bounded from below. Our assumption is more strict, but does not influence anything in the analysis of the equations we consider. In addition, our definition rules out the ambiguity of topologies on \mathcal{T} .

Remark 2.2. One of the simplest non-trivial examples of a regularity structure is given by the “abstract polynomials” in $d + 1$ indeterminates X_i , with $i = 0, \dots, d$. The set \mathcal{A} in this case consists of the values $\alpha \in \mathbf{N}$ such that $\alpha \leq r$, for some $r < \infty$ and, for each $\alpha \in \mathcal{A}$, the space \mathcal{T}_α contains all monomials in the X_i of scaled degree α . The structure group $\mathcal{G}_{\text{poly}}$ is then simply the group of translations in \mathbf{R}^{d+1} acting on X^k by $h \mapsto (X - h)^k$.

We now fix $r > 0$ to be sufficiently large and denote by $\mathcal{T}_{\text{poly}}$ the space of such polynomials of scaled degree r and by $\mathcal{F}_{\text{poly}}$ the set $\{X^k : |k|_s \leq r\}$. We will only ever consider regularity structures containing $\mathcal{T}_{\text{poly}}$ as a subspace. In particular, we always assume that there’s a natural morphism $\mathcal{G} \rightarrow \mathcal{G}_{\text{poly}}$ compatible with the action of $\mathcal{G}_{\text{poly}}$ on $\mathcal{T}_{\text{poly}} \hookrightarrow \mathcal{T}$.

Remark 2.3. For $\tau \in \mathcal{T}$ we will write $\mathcal{Q}_\alpha \tau$ for its canonical projection onto \mathcal{T}_α , and define $\|\tau\|_\alpha \stackrel{\text{def}}{=} \|\mathcal{Q}_\alpha \tau\|$. We also write $\mathcal{Q}_{<\alpha}$ for the projection onto $\mathcal{T}_{<\alpha}$, etc.

Another object in the theory of regularity structures is a model. Given an abstract expansion, the model converts it into a concrete distribution describing its local behaviour around every point. We modify the original definition of model in [Hai14], in order to be able to describe time-dependent distributions.

Definition 2.4. Given a regularity structure $\mathcal{S} = (\mathcal{T}, \mathcal{G})$, an *inhomogeneous model* (Π, Γ, Σ) consists of the following three elements:

- A collection of maps $\Gamma^t : \mathbf{R}^d \times \mathbf{R}^d \rightarrow \mathcal{G}$, parametrised by $t \in \mathbf{R}$, such that

$$\Gamma_{xx}^t = 1, \quad \Gamma_{xy}^t \Gamma_{yz}^t = \Gamma_{xz}^t, \quad (2.1)$$

for any $x, y, z \in \mathbf{R}^d$ and $t \in \mathbf{R}$, and the action of Γ_{xy}^t on polynomials is given as in Remark 2.2 with $h = (0, y - x)$.

- A collection of maps $\Sigma_x : \mathbf{R} \times \mathbf{R} \rightarrow \mathcal{G}$, parametrized by $x \in \mathbf{R}^d$, such that, for any $x \in \mathbf{R}^d$ and $s, r, t \in \mathbf{R}$, one has

$$\Sigma_x^{tt} = 1, \quad \Sigma_x^{sr} \Sigma_x^{rt} = \Sigma_x^{st}, \quad \Sigma_x^{st} \Gamma_{xy}^t = \Gamma_{xy}^s \Sigma_y^{st}, \quad (2.2)$$

and the action of Σ_x^{st} on polynomials is given as in Remark 2.2 with $h = (t - s, 0)$.

- A collection of linear maps $\Pi_x^t : \mathcal{T} \rightarrow \mathcal{S}'(\mathbf{R}^d)$, such that

$$\Pi_y^t = \Pi_x^t \Gamma_{xy}^t, \quad (\Pi_x^t X^{(0, \bar{k})})(y) = (y - x)^{\bar{k}}, \quad (\Pi_x^t X^{(k_0, \bar{k})})(y) = 0, \quad (2.3)$$

for all $x, y \in \mathbf{R}^d$, $t \in \mathbf{R}$, $\bar{k} \in \mathbf{N}^d$, $k_0 \in \mathbf{N}$ such that $k_0 > 0$.

Moreover, for any $\gamma > 0$ and every $T > 0$, there is a constant C for which the analytic bounds

$$|\langle \Pi_x^t \tau, \varphi_x^\lambda \rangle| \leq C \|\tau\| \lambda^l, \quad \|\Gamma_{xy}^t \tau\|_m \leq C \|\tau\| |x - y|^{l-m}, \quad (2.4a)$$

$$\|\Sigma_x^{st} \tau\|_m \leq C \|\tau\| |t - s|^{(l-m)/s_0}, \quad (2.4b)$$

hold uniformly over all $\tau \in \mathcal{T}_l$, with $l \in \mathcal{A}$ and $l < \gamma$, all $m \in \mathcal{A}$ such that $m < l$, all $\lambda \in (0, 1]$, all $\varphi \in \mathcal{B}_0^r(\mathbf{R}^d)$ with $r > -\lfloor \min \mathcal{A} \rfloor$, and all $t, s \in [-T, T]$ and $x, y \in \mathbf{R}^d$ such that $|t - s| \leq 1$ and $|x - y| \leq 1$.

In addition, we say that the map Π has time regularity $\delta > 0$, if the bound

$$|\langle (\Pi_x^t - \Pi_x^s)\tau, \varphi_x^\lambda \rangle| \leq C \|\tau\| |t - s|^{\delta/s_0} \lambda^{l-\delta}, \quad (2.5)$$

holds for all $\tau \in \mathcal{T}_l$ and the other parameters as before.

Remark 2.5. For a model $Z = (\Pi, \Gamma, \Sigma)$, we denote by $\|\Pi\|_{\gamma;T}$, $\|\Gamma\|_{\gamma;T}$ and $\|\Sigma\|_{\gamma;T}$ the smallest constants C such that the bounds on Π , Γ and Σ in (2.4a) and (2.4b) hold. Furthermore, we define

$$\|Z\|_{\gamma;T} \stackrel{\text{def}}{=} \|\Pi\|_{\gamma;T} + \|\Gamma\|_{\gamma;T} + \|\Sigma\|_{\gamma;T}.$$

If $\bar{Z} = (\bar{\Pi}, \bar{\Gamma}, \bar{\Sigma})$ is another model, then we also define the ‘‘distance’’ between two models

$$\|Z; \bar{Z}\|_{\gamma;T} \stackrel{\text{def}}{=} \|\Pi - \bar{\Pi}\|_{\gamma;T} + \|\Gamma - \bar{\Gamma}\|_{\gamma;T} + \|\Sigma - \bar{\Sigma}\|_{\gamma;T}. \quad (2.6)$$

We note that the norms on the right-hand side still make sense with Γ and Σ viewed as linear maps on \mathcal{T} . We also set $\|\Pi\|_{\delta,\gamma;T} \stackrel{\text{def}}{=} \|\Pi\|_{\gamma;T} + C$, where C is the smallest constant such that the bound (2.5) holds, and we define

$$\|Z\|_{\delta,\gamma;T} \stackrel{\text{def}}{=} \|\Pi\|_{\delta,\gamma;T} + \|\Gamma\|_{\gamma;T} + \|\Sigma\|_{\gamma;T}.$$

Finally, we define the ‘‘distance’’ $\|Z; \bar{Z}\|_{\delta,\gamma;T}$ as in (2.6).

Remark 2.6. In [Hai14, Def. 2.17] the analytic bounds on a model were assumed to hold locally uniformly. In the problems which we aim to consider, the models are periodic in space, which allows us to require the bounds to hold globally.

Remark 2.7. For a given model (Π, Γ, Σ) we can define the following two objects

$$(\tilde{\Pi}_{(t,x)}\tau)(s, y) = (\Pi_x^s \Sigma_x^{st}\tau)(y), \quad \tilde{\Gamma}_{(t,x),(s,y)} = \Gamma_{xy}^t \Sigma_y^{ts} = \Sigma_x^{ts} \Gamma_{xy}^s, \quad (2.7)$$

for $\tau \in \mathcal{T}$. (Of course, in general we cannot fix the spatial point y in the definition of $\tilde{\Pi}$, and we should really write $((\tilde{\Pi}_{(t,x)}\tau)(s, \cdot))(\varphi) = (\Pi_x^s \Sigma_x^{st}\tau)(\varphi)$ instead, for any test function φ , but the notation (2.7) is more suggestive.) One can then easily verify that the pair $(\tilde{\Pi}, \tilde{\Gamma})$ is a model in the original sense of [Hai14, Def. 2.17].

2.2 Inhomogeneous modelled distributions

Modelled distributions represent abstract expansions in the basis of a regularity structure. In order to be able to describe the singularity coming from the behaviour of our solutions near time 0, we introduce inhomogeneous modelled distributions which admit a certain blow-up as time goes to zero.

Given a regularity structure $\mathcal{T} = (\mathcal{T}, \mathcal{G})$ with a model $Z = (\Pi, \Gamma, \Sigma)$, values $\gamma, \eta \in \mathbf{R}$ and a final time $T > 0$, we consider maps $H : (0, T] \times \mathbf{R}^d \rightarrow \mathcal{T}_{<\gamma}$ and define

$$\begin{aligned} \|H\|_{\gamma, \eta; T} \stackrel{\text{def}}{=} & \sup_{t \in (0, T]} \sup_{x \in \mathbf{R}^d} \sup_{l < \gamma} |t|_0^{(l-\eta) \vee 0} \|H_t(x)\|_l \\ & + \sup_{t \in (0, T]} \sup_{\substack{x \neq y \in \mathbf{R}^d \\ |x-y| \leq 1}} \sup_{l < \gamma} \frac{\|H_t(x) - \Gamma_{xy}^t H_t(y)\|_l}{|t|_0^{\eta-\gamma} |x-y|^{\gamma-l}}, \end{aligned} \quad (2.8)$$

where $l \in \mathcal{A}$ in the third supremum. Then the space $\mathcal{D}_T^{\gamma, \eta}$ consists of all such functions H , for which one has

$$\|H\|_{\gamma, \eta; T} \stackrel{\text{def}}{=} \|H\|_{\gamma, \eta; T} + \sup_{\substack{s \neq t \in (0, T] \\ |t-s| \leq |t, s|_0^{s_0}}} \sup_{x \in \mathbf{R}^d} \sup_{l < \gamma} \frac{\|H_t(x) - \Sigma_x^{ts} H_s(x)\|_l}{|t, s|_0^{\eta-\gamma} |t-s|^{(\gamma-l)/s_0}} < \infty. \quad (2.9)$$

The quantities $|t|_0$ and $|t, s|_0$ used in these definitions were introduced in (1.5). Elements of these spaces will be called *inhomogeneous modelled distributions*.

Remark 2.8. The norm in (2.9) depends on Γ and Σ , but does *not* depend on Π ; this fact will be crucial in the sequel. When we want to stress the dependency on the model, we will also write $\mathcal{D}_T^{\gamma, \eta}(Z)$.

Remark 2.9. In contrast to the singular modelled distributions from [Hai14, Def. 6.2], we do not require the restriction $|x-y| \leq |t, s|_0$ in the second term in (2.8). This is due to the fact that we consider the space and time variables separately (see the proof of Theorem 2.21, where this fact is used).

Remark 2.10. Since our spaces $\mathcal{D}_T^{\gamma, \eta}$ are almost identical to those of [Hai14], the multiplication and differentiation results from [Hai14, Sec. 6] hold also for our definition.

To be able to compare two modelled distributions $H \in \mathcal{D}_T^{\gamma, \eta}(Z)$ and $\bar{H} \in \mathcal{D}_T^{\gamma, \eta}(\bar{Z})$, we define the quantities

$$\begin{aligned} \|H; \bar{H}\|_{\gamma, \eta; T} \stackrel{\text{def}}{=} & \sup_{t \in (0, T]} \sup_{x \in \mathbf{R}^d} \sup_{l < \gamma} |t|_0^{(l-\eta) \vee 0} \|H_t(x) - \bar{H}_t(x)\|_l \\ & + \sup_{t \in (0, T]} \sup_{\substack{x \neq y \in \mathbf{R}^d \\ |x-y| \leq 1}} \sup_{l < \gamma} \frac{\|H_t(x) - \Gamma_{xy}^t H_t(y) - \bar{H}_t(x) + \bar{\Gamma}_{xy}^t \bar{H}_t(y)\|_l}{|t|_0^{\eta-\gamma} |x-y|^{\gamma-l}}, \\ \|H; \bar{H}\|_{\gamma, \eta; T} \stackrel{\text{def}}{=} & \|H; \bar{H}\|_{\gamma, \eta; T} \\ & + \sup_{\substack{s \neq t \in (0, T] \\ |t-s| \leq |t, s|_0^{s_0}}} \sup_{x \in \mathbf{R}^d} \sup_{l < \gamma} \frac{\|H_t(x) - \Sigma_x^{ts} H_s(x) - \bar{H}_t(x) + \bar{\Sigma}_x^{ts} \bar{H}_s(x)\|_l}{|t, s|_0^{\eta-\gamma} |t-s|^{(\gamma-l)/s_0}}. \end{aligned}$$

The ‘‘reconstruction theorem’’ is one of the key results of the theory of regularity structures. Here is its statement in our current framework.

Theorem 2.11. *Let $\mathcal{F} = (\mathcal{T}, \mathcal{G})$ be a regularity structure with $\alpha \stackrel{\text{def}}{=} \min \mathcal{A} < 0$ and $Z = (\Pi, \Gamma, \Sigma)$ be a model. Then, for every $\eta \in \mathbf{R}$, $\gamma > 0$ and $T > 0$, there is a unique family of linear operators $\mathcal{R}_t : \mathcal{D}_T^{\gamma, \eta}(Z) \rightarrow \mathcal{C}^\alpha(\mathbf{R}^d)$, parametrised by $t \in (0, T]$, such that the bound*

$$|\langle \mathcal{R}_t H_t - \Pi_x^t H_t(x), \varphi_x^\lambda \rangle| \lesssim \lambda^\gamma |t|_0^{\eta-\gamma} \|H\|_{\gamma, \eta; T} \|\Pi\|_{\gamma; T}, \quad (2.11)$$

holds uniformly in $H \in \mathcal{D}_T^{\gamma, \eta}(Z)$, $t \in (0, T]$, $x \in \mathbf{R}^d$, $\lambda \in (0, 1]$ and $\varphi \in \mathcal{B}_0^r(\mathbf{R}^d)$ with $r > -\lfloor \alpha \rfloor$.

If furthermore the map Π has time regularity $\delta > 0$, then, for any $\tilde{\delta} \in (0, \delta]$ such that $\tilde{\delta} \leq (m - \zeta)$ for all $\zeta, m \in ((-\infty, \gamma) \cap \mathcal{A}) \cup \{\gamma\}$ such that $\zeta < m$, the function $t \mapsto \mathcal{R}_t H_t$ satisfies

$$\|\mathcal{R}H\|_{\mathcal{C}_{\eta-\gamma, T}^{\tilde{\delta}, \alpha}} \lesssim \|\Pi\|_{\delta, \gamma; T} (1 + \|\Sigma\|_{\gamma; T}) \|H\|_{\gamma, \eta; T}. \quad (2.12)$$

Let $\bar{Z} = (\bar{\Pi}, \bar{\Gamma}, \bar{\Sigma})$ be another model for the same regularity structure, and let $\bar{\mathcal{R}}_t$ be the operator as above, but for the model \bar{Z} . Moreover, let the maps Π and $\bar{\Pi}$ have time regularities $\delta > 0$. Then, for every $H \in \mathcal{D}_T^{\gamma, \eta}(Z)$ and $\bar{H} \in \mathcal{D}_T^{\gamma, \eta}(\bar{Z})$, the maps $t \mapsto \mathcal{R}_t H_t$ and $t \mapsto \bar{\mathcal{R}}_t \bar{H}_t$ satisfy

$$\|\mathcal{R}H - \bar{\mathcal{R}}\bar{H}\|_{\mathcal{C}_{\eta-\gamma, T}^{\tilde{\delta}, \alpha}} \lesssim \|H; \bar{H}\|_{\gamma, \eta; T} + \|Z; \bar{Z}\|_{\delta, \gamma; T}, \quad (2.13)$$

for any $\tilde{\delta}$ as above, and where the proportionality constant depends on $\|H\|_{\gamma, \eta; T}$, $\|\bar{H}\|_{\gamma, \eta; T}$, $\|Z\|_{\delta, \gamma; T}$ and $\|\bar{Z}\|_{\delta, \gamma; T}$.

Proof. Existence and uniqueness of the maps \mathcal{R}_t , as well as the bound (2.11), follow from [Hai14, Thm. 3.10]. The uniformity in time in (2.11) follows from the uniformity of the corresponding bounds in [Hai14, Thm. 3.10].

To prove that $t \mapsto \mathcal{R}_t H_t$ belongs to $\mathcal{C}_{\eta-\gamma}^{\tilde{\delta}, \alpha}([0, T], \mathbf{R}^d)$, we will first bound $\langle \mathcal{R}_t H_t, \varrho_x^\lambda \rangle$, for $\lambda \in (0, 1]$, $x \in \mathbf{R}^d$ and $\varrho \in \mathcal{B}_0^r(\mathbf{R}^d)$. Using (2.11) and properties of Π and H we get

$$\begin{aligned} |\langle \mathcal{R}_t H_t, \varrho_x^\lambda \rangle| &\leq |\langle \mathcal{R}_t H_t - \Pi_x^t H_t(x), \varrho_x^\lambda \rangle| + |\langle \Pi_x^t H_t(x), \varrho_x^\lambda \rangle| \\ &\lesssim \lambda^\gamma |t|_0^{\eta-\gamma} + \sum_{\zeta \in \lfloor \alpha, \gamma \rfloor \cap \mathcal{A}} \lambda^\zeta |t|_0^{(\eta-\zeta) \wedge 0} \lesssim \lambda^\alpha |t|_0^{\eta-\gamma}, \end{aligned} \quad (2.14)$$

where the proportionality constant is affine in $\|H\|_{\gamma, \eta; T} \|\Pi\|_{\gamma; T}$, and α is the minimal homogeneity in \mathcal{A} .

In order to obtain the time regularity of $t \mapsto \mathcal{R}_t H_t$, we show that the distribution $\zeta_x^{st} \stackrel{\text{def}}{=} \Pi_x^t H_t(x) - \Pi_x^s H_s(x)$ satisfies the bound

$$|\langle \zeta_x^{st} - \zeta_y^{st}, \varrho_x^\lambda \rangle| \lesssim |t - s|^{\tilde{\delta}/s_0} |s, t|_0^{\eta-\gamma} |x - y|^{\gamma - \tilde{\delta} - \alpha} \lambda^\alpha, \quad (2.15)$$

uniformly over all $x, y \in \mathbf{R}^d$ such that $\lambda \leq |x - y| \leq 1$, all $s, t \in \mathbf{R}$, and for any value of $\tilde{\delta}$ as in the statement of the theorem. To this end, we consider two regimes: $|x - y| \leq |t - s|^{1/s_0}$ and $|x - y| > |t - s|^{1/s_0}$.

In the first case, when $|x - y| \leq |t - s|^{1/s_0}$, we write, using Definition 2.4,

$$\zeta_x^{st} - \zeta_y^{st} = \Pi_x^t (H_t(x) - \Gamma_{xy}^t H_t(y)) - \Pi_x^s (H_s(x) - \Gamma_{xy}^s H_s(y)), \quad (2.16)$$

and bound these two terms separately. From the properties (2.4a) and (2.9) we get

$$\begin{aligned} |\langle \Pi_x^t (H_t(x) - \Gamma_{xy}^t H_t(y)), \varrho_x^\lambda \rangle| &\lesssim \sum_{\zeta \in [\alpha, \gamma] \cap \mathcal{A}} \lambda^\zeta \|H_t(x) - \Gamma_{xy}^t H_t(y)\|_\zeta \\ &\lesssim \sum_{\zeta \in [\alpha, \gamma] \cap \mathcal{A}} \lambda^\zeta |x - y|^{\gamma - \zeta} |t|_0^{\eta - \gamma} \lesssim \lambda^\alpha |x - y|^{\gamma - \alpha} |t|_0^{\eta - \gamma}, \end{aligned} \quad (2.17)$$

where we have exploited the condition $|x - y| \geq \lambda$. Recalling now the case we consider, we can bound the last expression by the right-hand side of (2.15). The same estimate holds for the second term in (2.16).

Now, we will consider the case $|x - y| > |t - s|^{1/s_0}$. In this regime we use the definition of model and write

$$\begin{aligned} \zeta_x^{st} - \zeta_y^{st} &= (\Pi_x^t - \Pi_x^s)(H_t(x) - \Gamma_{xy}^t H_t(y)) + \Pi_x^s(1 - \Sigma_x^{st})(H_t(x) - \Gamma_{xy}^t H_t(y)) \\ &\quad - \Pi_x^s(H_s(x) - \Sigma_x^{st} H_t(x)) + \Pi_y^s(H_s(y) - \Sigma_y^{st} H_t(y)). \end{aligned} \quad (2.18)$$

The first term can be bounded exactly as (2.17), but using this time (2.5), i.e.

$$|\langle (\Pi_x^t - \Pi_x^s)(H_t(x) - \Gamma_{xy}^t H_t(y)), \varrho_x^\lambda \rangle| \lesssim \lambda^{\alpha - \delta} |x - y|^{\gamma - \alpha} |t|_0^{\eta - \gamma} |t - s|^{\delta/s_0}.$$

In order to estimate the second term in (2.18), we first notice that from (2.4b) and (2.9) we get

$$\begin{aligned} \|(1 - \Sigma_x^{st})(H_t(x) - \Gamma_{xy}^t H_t(y))\|_\zeta &\lesssim \sum_{\zeta < m < \gamma} |t - s|^{(m - \zeta)/s_0} \|H_t(x) - \Gamma_{xy}^t H_t(y)\|_m \\ &\lesssim \sum_{\zeta < m < \gamma} |t - s|^{(m - \zeta)/s_0} |x - y|^{\gamma - m} |t|_0^{\eta - \gamma} \lesssim |t - s|^{\tilde{\delta}/s_0} |x - y|^{\gamma - \tilde{\delta} - \zeta} |t|_0^{\eta - \gamma}, \end{aligned} \quad (2.19)$$

for any $\tilde{\delta} \leq \min_{m > \zeta \in \mathcal{A}}(m - \zeta)$, where we have used the assumption on the time variables. Hence, for the second term in (2.18) we have

$$\begin{aligned} |\langle \Pi_x^s(1 - \Sigma_x^{st})(H_t(x) - \Gamma_{xy}^t H_t(y)), \varrho_x^\lambda \rangle| \\ \lesssim |t - s|^{\tilde{\delta}/s_0} |t|_0^{\eta - \gamma} \sum_{\zeta < \gamma} \lambda^\zeta |x - y|^{\gamma - \tilde{\delta} - \zeta}. \end{aligned}$$

Since $|x - y| \geq \lambda$ and $\zeta \geq \alpha$, the estimate (2.15) holds for this expression.

The third term in (2.18) we bound using the properties (2.4a) and (2.9) by

$$\begin{aligned} |\langle \Pi_x^s(H_s(x) - \Sigma_x^{st} H_t(x)), \varrho_x^\lambda \rangle| &\lesssim \sum_{\zeta < \gamma} \lambda^\zeta \|H_s(x) - \Sigma_x^{st} H_t(x)\|_\zeta \\ &\lesssim \sum_{\zeta < \gamma} \lambda^\zeta |t - s|^{(\gamma - \zeta)/s_0} |t, s|_0^{\eta - \gamma}. \end{aligned} \quad (2.20)$$

It follows from $|x - y| \geq \lambda$, $|x - y| > |t - s|^{1/s_0}$ and $\zeta \geq \alpha$, that the latter can be estimated as in (2.15), when $\tilde{\delta} \leq \min\{\gamma - \zeta : \zeta \in \mathcal{A}, \zeta < \gamma\}$. The same bound holds for the last term in (2.18), and this finishes the proof of (2.15).

In view of the bound (2.15) and [Hai14, Prop. 3.25], we conclude that

$$|\langle \mathcal{R}_t H_t - \mathcal{R}_s H_s - \zeta_x^{st}, \varrho_x^\lambda \rangle| \lesssim |t - s|^{\tilde{\delta}/s_0} \lambda^{\gamma - \tilde{\delta}} |s, t|_0^{\eta - \gamma}, \quad (2.21)$$

uniformly over $s, t \in \mathbf{R}$ and the other parameters as in (2.11). Thus, we can write

$$\langle \mathcal{R}_t H_t - \mathcal{R}_s H_s, \varrho_x^\lambda \rangle = \langle \mathcal{R}_t H_t - \mathcal{R}_s H_s - \zeta_x^{st}, \varrho_x^\lambda \rangle + \langle \zeta_x^{st}, \varrho_x^\lambda \rangle,$$

where the first term is bounded in (2.21). The second term we can write as

$$\begin{aligned} \langle \zeta_x^{st}, \varrho_x^\lambda \rangle &= \langle (\Pi_x^t - \Pi_x^s) H_t(x), \varrho_x^\lambda \rangle + \langle \Pi_x^s (H_t(x) - \Sigma_x^{ts} H_s(x)), \varrho_x^\lambda \rangle \\ &\quad + \langle \Pi_x^s (\Sigma_x^{ts} - 1) H_s(x), \varrho_x^\lambda \rangle, \end{aligned}$$

which can be bounded by $|t - s|^{\tilde{\delta}/s_0} \lambda^{\alpha - \tilde{\delta}} |s, t|_0^{\eta - \gamma}$, using (2.5), (2.20) and (2.4b). Here, in order to estimate the last term, we act similarly to (2.19). Combining all these bounds together, we conclude that

$$|\langle \mathcal{R}_t H_t - \mathcal{R}_s H_s, \varrho_x^\lambda \rangle| \lesssim |t - s|^{\tilde{\delta}/s_0} \lambda^{\alpha - \tilde{\delta}} |s, t|_0^{\eta - \gamma}, \quad (2.22)$$

which finishes the proof of the claim.

The bound (2.13) can be shown in a similar way. More precisely, similarly to (2.14) and using [Hai14, Eq. 3.4], we can show that

$$|\langle \mathcal{R}_t H_t - \bar{\mathcal{R}}_t \bar{H}_t, \varrho_x^\lambda \rangle| \lesssim \lambda^\alpha |t|_0^{\eta - \gamma} (\|\Pi\|_{\gamma; T} \|H; \bar{H}\|_{\gamma, \eta; T} + \|\Pi - \bar{\Pi}\|_{\gamma; T} \|\bar{H}\|_{\gamma, \eta; T}).$$

Denoting $\bar{\zeta}_x^{st} \stackrel{\text{def}}{=} \bar{\Pi}_x^t \bar{H}_t(x) - \bar{\Pi}_x^s \bar{H}_s(x)$ and acting as above, we can prove an analogue of (2.21):

$$\begin{aligned} &|\langle \mathcal{R}_t H_t - \bar{\mathcal{R}}_t \bar{H}_t - \mathcal{R}_s H_s + \bar{\mathcal{R}}_s \bar{H}_s - \zeta_x^{st} + \bar{\zeta}_x^{st}, \varrho_x^\lambda \rangle| \\ &\lesssim |t - s|^{\tilde{\delta}/s_0} \lambda^{\gamma - \tilde{\delta}} |s, t|_0^{\eta - \gamma} (\|H; \bar{H}\|_{\gamma, \eta; T} + \|Z; \bar{Z}\|_{\delta, \gamma; T}), \end{aligned}$$

with the values of $\tilde{\delta}$ as before. Finally, similarly to (2.22) we get

$$\begin{aligned} &|\langle \mathcal{R}_t H_t - \bar{\mathcal{R}}_t \bar{H}_t - \mathcal{R}_s H_s + \bar{\mathcal{R}}_s \bar{H}_s, \varrho_x^\lambda \rangle| \lesssim |t - s|^{\tilde{\delta}/s_0} \lambda^{\alpha - \tilde{\delta}} |s, t|_0^{\eta - \gamma} \\ &\quad \times (\|H; \bar{H}\|_{\gamma, \eta; T} + \|Z; \bar{Z}\|_{\delta, \gamma; T}), \end{aligned}$$

which finishes the proof. \square

Definition 2.12. We will call the map \mathcal{R} , introduced in Theorem 2.11, the *reconstruction operator*, and we will always postulate in what follows that $\mathcal{R}_t = 0$, for $t \leq 0$.

Remark 2.13. One can see that the map $\tilde{\mathcal{R}}(t, \cdot) \stackrel{\text{def}}{=} \mathcal{R}_t(\cdot)$ is the reconstruction operator for the model (2.7) in the sense of [Hai14, Thm. 3.10].

2.3 Convolutions with singular kernels

In the definition of a mild solution to a parabolic stochastic PDE, convolutions with singular kernels are involved. In particular Schauder estimates plays a key role. To describe this on the abstract level, we introduce the abstract integration map.

Definition 2.14. Given a regularity structure $\mathcal{F} = (\mathcal{T}, \mathcal{G})$, a linear map $\mathcal{I} : \mathcal{T} \rightarrow \mathcal{T}$ is said to be an *abstract integration map* of order $\beta > 0$ if it satisfies the following properties:

- One has $\mathcal{I} : \mathcal{T}_m \rightarrow \mathcal{T}_{m+\beta}$, for every $m \in \mathcal{A}$ such that $m + \beta \in \mathcal{A}$.
- For every $\tau \in \mathcal{T}_{\text{poly}}$, one has $\mathcal{I}\tau = 0$, where $\mathcal{T}_{\text{poly}} \subset \mathcal{T}$ contains the polynomial part of \mathcal{T} and was introduced in Remark 2.2.
- One has $\mathcal{I}\Gamma\tau - \Gamma\mathcal{I}\tau \in \mathcal{T}_{\text{poly}}$, for every $\tau \in \mathcal{T}$ and $\Gamma \in \mathcal{G}$.

Remark 2.15. The second and third properties are dictated by the special role played by polynomials in the Taylor expansion. One can find a more detailed motivation for this definition in [Hai14, Sec. 5]. In general, we also allow for the situation where \mathcal{I} has a domain which isn't all of \mathcal{T} .

Now, we will define the singular kernels, convolutions with which we are going to describe.

Definition 2.16. A function $K : \mathbf{R}^{d+1} \setminus \{0\} \rightarrow \mathbf{R}$ is regularising of order $\beta > 0$, if there is a constant $r > 0$ such that we can decompose

$$K = \sum_{n \geq 0} K^{(n)}, \quad (2.23)$$

in such a way that each term $K^{(n)}$ is supported in $\{z \in \mathbf{R}^{d+1} : \|z\|_s \leq c2^{-n}\}$ for some $c > 0$, satisfies

$$|D^k K^{(n)}(z)| \lesssim 2^{(|s| - \beta + |k|_s)n}, \quad (2.24)$$

for every multiindex k with $|k|_s \leq r$, and annihilates every polynomial of scaled degree r . (See [Hai14, Sec. 5] for more details.)

Now, we will describe the action of a model on the abstract integration map. When it is convenient for us, we will write $K_t(x) = K(z)$, for $z = (t, x)$.

Definition 2.17. Let \mathcal{I} be an abstract integration map of order β for a regularity structure $\mathcal{F} = (\mathcal{T}, \mathcal{G})$, let $Z = (\Pi, \Gamma, \Sigma)$ be a model and let K be regularising of order β with $r > -\lfloor \min \mathcal{A} \rfloor$. We say that Z realises K for \mathcal{I} , if for every $\alpha \in \mathcal{A}$ and every $\tau \in \mathcal{T}_\alpha$ one has the identity

$$\Pi_x^t (\mathcal{I}\tau + \mathcal{J}_{t,x}\tau)(y) = \int_{\mathbf{R}} \langle \Pi_x^s \Sigma_x^{st} \tau, K_{t-s}(y - \cdot) \rangle ds, \quad (2.25)$$

where $\mathcal{J}_{t,x}\tau$ is defined by

$$\mathcal{J}_{t,x}\tau \stackrel{\text{def}}{=} \sum_{|k|_s < \alpha + \beta} \frac{X^k}{k!} \int_{\mathbf{R}} \langle \Pi_x^s \Sigma_x^{st} \tau, D^k K_{t-s}(x - \cdot) \rangle ds, \quad (2.26)$$

where $k \in \mathbf{N}^{d+1}$ and the derivative D^k is in time-space. Moreover, we require that

$$\begin{aligned}\Gamma_{xy}^t(\mathcal{I} + \mathcal{J}_{t,y}) &= (\mathcal{I} + \mathcal{J}_{t,x})\Gamma_{xy}^t, \\ \Sigma_x^{st}(\mathcal{I} + \mathcal{J}_{t,x}) &= (\mathcal{I} + \mathcal{J}_{s,x})\Sigma_x^{st},\end{aligned}\tag{2.27}$$

for all $s, t \in \mathbf{R}$ and $x, y \in \mathbf{R}^d$.

Remark 2.18. We define the integrals in (2.25) and (2.26) as sums of the same integrals, but using the functions $K^{(n)}$ from the expansion (2.23). Since these integrals coincide with those from [Hai14] for the model (2.7), it follows from [Hai14, Lem. 5.19] that these sums converge absolutely, and hence the expressions in (2.25) and (2.26) are well defined.

Remark 2.19. The identities (2.27) should be viewed as defining $\Gamma_{xy}^t \mathcal{I} \tau$ and $\Sigma_x^{st} \mathcal{I} \tau$ in terms of $\Gamma_{xy}^t \tau$, $\Sigma_x^{st} \tau$, and (2.26).

With all these notations at hand we introduce the following operator acting on modelled distribution $H \in \mathcal{D}_T^{\gamma, \eta}(Z)$ with $\gamma + \beta > 0$:

$$(\mathcal{K}_\gamma H)_t(x) \stackrel{\text{def}}{=} \mathcal{I}H_t(x) + \mathcal{J}_{t,x}H_t(x) + (\mathcal{N}_\gamma H)_t(x).\tag{2.28}$$

Here, the last term is $\mathcal{T}_{\text{poly}}$ -valued and is given by

$$(\mathcal{N}_\gamma H)_t(x) \stackrel{\text{def}}{=} \sum_{|k|_s < \gamma + \beta} \frac{X^k}{k!} \int_{\mathbf{R}} \langle \mathcal{R}_s H_s - \Pi_x^s \Sigma_x^{st} H_t(x), D^k K_{t-s}(x - \cdot) \rangle ds,\tag{2.29}$$

where as before $k \in \mathbf{N}^{d+1}$ and the derivative D^k is in time-space (see Definition 2.12 for consistency of notation).

Remark 2.20. It follows from Remark 2.13 and the proof of [Hai14, Thm. 5.12], that the integral in (2.29) is well-defined, if we express it as a sum of the respective integrals with the functions $K^{(n)}$ in place of K . (See also the definition of the operator \mathbf{R}^+ in [Hai14, Sec. 7.1].)

The modelled distribution $\mathcal{K}_\gamma H$ represents the space-time convolution of H with K , and the following result shows that this action ‘‘improves’’ regularity by β .

Theorem 2.21. *Let $\mathcal{T} = (\mathcal{T}, \mathcal{G})$ be a regularity structure with the minimal homogeneity α , let \mathcal{I} be an abstract integration map of an integer order $\beta > 0$, let K be a singular function regularising by β , and let $Z = (\Pi, \Gamma, \Sigma)$ be a model, which realises K for \mathcal{I} . Furthermore, let $\gamma > 0$, $\eta < \gamma$, $\eta > -\mathfrak{s}_0$, $\gamma < \eta + \mathfrak{s}_0$, $\gamma + \beta \notin \mathbf{N}$, $\alpha + \beta > 0$ and $r > -\lfloor \alpha \rfloor$, $r > \gamma + \beta$ in Definition 2.16.*

Then \mathcal{K}_γ maps $\mathcal{D}_T^{\gamma, \eta}(Z)$ into $\mathcal{D}_T^{\bar{\gamma}, \bar{\eta}}(Z)$, where $\bar{\gamma} = \gamma + \beta$, $\bar{\eta} = \eta \wedge \alpha + \beta$, and for any $H \in \mathcal{D}_T^{\gamma, \eta}(Z)$ the following bound holds

$$\|\mathcal{K}_\gamma H\|_{\bar{\gamma}, \bar{\eta}; T} \lesssim \|H\|_{\gamma, \eta; T} \|\Pi\|_{\gamma; T} \|\Sigma\|_{\gamma; T} (1 + \|\Gamma\|_{\bar{\gamma}; T} + \|\Sigma\|_{\bar{\gamma}; T}).\tag{2.30}$$

Furthermore, for every $t \in (0, T]$, one has the identity

$$\mathcal{R}_t(\mathcal{K}_\gamma H)_t(x) = \int_0^t \langle \mathcal{R}_s H_s, K_{t-s}(x - \cdot) \rangle ds . \quad (2.31)$$

Let $\bar{Z} = (\bar{\Pi}, \bar{\Gamma}, \bar{\Sigma})$ be another model realising K for \mathcal{I} , which satisfies the same assumptions, and let $\bar{\mathcal{K}}_\gamma$ be defined by (2.28) for this model. Then one has

$$\|\mathcal{K}_\gamma H; \bar{\mathcal{K}}_\gamma \bar{H}\|_{\bar{\gamma}, \bar{\eta}; T} \lesssim \|H; \bar{H}\|_{\gamma, \eta; T} + \|Z; \bar{Z}\|_{\bar{\gamma}; T} , \quad (2.32)$$

for all $H \in \mathcal{D}_T^{\gamma, \eta}(Z)$ and $\bar{H} \in \mathcal{D}_T^{\bar{\gamma}, \bar{\eta}}(\bar{Z})$. Here, the proportionality constant depends on $\|H\|_{\gamma, \eta; T}$, $\|\bar{H}\|_{\bar{\gamma}, \bar{\eta}; T}$ and the norms on the models Z and \bar{Z} involved in the estimate (2.30).

Proof. In view of Remarks 2.7 and 2.13, the required bounds on the components of $(\mathcal{K}_\gamma H)_t(x)$ and $(\mathcal{K}_\gamma H)_t(x) - \Sigma_x^{ts}(\mathcal{K}_\gamma H)_s(x)$, as well as on the components of $(\mathcal{K}_\gamma H)_t(y) - \Gamma_{yx}^t(\mathcal{K}_\gamma H)_t(x)$ with non-integer homogeneities, can be obtained in exactly the same way as in [Hai14, Prop. 6.16]. (See the definition of the operator \mathbf{R}^+ in [Hai14, Sec. 7.1].)

In order to get the required bounds on the elements of $(\mathcal{K}_\gamma H)_t(x) - \Gamma_{xy}^t(\mathcal{K}_\gamma H)_t(y)$ with integer homogeneities, we need to modify the proof of [Hai14, Prop. 6.16]. The problem is that our definition of modelled distributions is slightly different than the one in [Hai14, Def. 6.2] (see Remark 2.9). That's why we have to consider only two regimes, $c2^{-n+1} \leq |x - y|$ and $c2^{-n+1} > |x - y|$, in the proof of [Hai14, Prop. 6.16], where c is from Definition 2.16. The only place in the proof, which requires a special treatment, is the derivation of the estimate

$$\left| \int_{\mathbf{R}} \langle \mathcal{R}_s H_s - \Pi_x^s H_s(x), D^k K_{t-s}^{(n)}(x - \cdot) \rangle ds \right| \lesssim 2^{(|k|_s - \gamma - \beta)n} |t|_0^{\eta - \gamma} ,$$

which in our case follows trivially from Theorem 2.11 and Definition 2.16. Here is the place where we need $\gamma - \eta < \mathfrak{s}_0$, in order to have an integrable singularity. Here, we use the same argument as in the proof of [Hai14, Thm. 7.1] to make sure that the time interval does not increase.

With respective modifications of the proof of [Hai14, Prop. 6.16] we can also show that (2.31) and (2.32) hold. \square

3 Solutions to parabolic stochastic PDEs

We consider a general parabolic stochastic PDE of the form

$$\partial_t u = Au + F(u, \xi) , \quad u(0, \cdot) = u_0(\cdot) , \quad (3.1)$$

on $\mathbf{R}_+ \times \mathbf{R}^d$, where u_0 is the initial data, ξ is a rough noise, F is a function in u and ξ , which depends in general on the space-time point z and which is affine in ξ , and A is a differential operator such that $\partial_t - A$ has a Green's function G , i.e. G is the distributional solution of $(\partial_t - A)G = \delta_0$. Then we require the following assumption to be satisfied.

Assumption 3.1. *The operator A is given by $Q(\nabla)$, for Q a homogeneous polynomial on \mathbf{R}^d of some even degree $\beta > 0$. Its Green's function $G : \mathbf{R}^{d+1} \setminus \{0\} \mapsto \mathbf{R}$ is smooth, non-anticipative, i.e. $G_t = 0$ for $t \leq 0$, and for $\lambda > 0$ satisfies the scaling relation*

$$\lambda^d G_{\lambda^\beta t}(\lambda x) = G_t(x) .$$

Remark 3.2. One can find in [Hör55] precise conditions on Q such that G satisfies Assumption 3.1.

In order to apply the abstract integration developed in the previous section, we would like the localised singular part of G to have the properties from Definition 2.16. The following result, following from [Hai14, Lem. 7.7], shows that this is indeed the case.

Lemma 3.3. *Let us consider functions u supported in $\mathbf{R}_+ \times \mathbf{R}^d$ and periodic in the spatial variable with some fixed period. If Assumption 3.1 is satisfied with some $\beta > 0$, then we can write $G = K + R$, in such a way that the identity*

$$(G \star u)(z) = (K \star u)(z) + (R \star u)(z) ,$$

holds for every such function u and every $z \in (-\infty, 1] \times \mathbf{R}^d$, where \star is the space-time convolution. Furthermore, K has the properties from Definition 2.16 with the parameters β and some arbitrary (but fixed) value r , and the scaling $\mathfrak{s} = (\beta, 1, \dots, 1)$. The function R is smooth, non-anticipative and compactly supported.

In particular, it follows from Lemma 3.3 that for any $\gamma > 0$ and any periodic $\zeta_t \in \mathcal{C}^\alpha(\mathbf{R}^d)$, with $t \in \mathbf{R}$ and with probably an integrable singularity at $t = 0$, we can define

$$(R_\gamma \zeta)_t(x) \stackrel{\text{def}}{=} \sum_{|k|_s < \gamma} \frac{X^k}{k!} \int_{\mathbf{R}} \langle \zeta_s, D^k R_{t-s}(x - \cdot) \rangle ds , \quad (3.2)$$

where $k \in \mathbf{N}^{d+1}$ and D^k is taken in time-space.

3.1 Regularity structures for locally subcritical stochastic PDEs

In this section we provide conditions on the equation (3.1), under which one can build a regularity structure for it. More precisely, we consider the mild form of equation (3.1):

$$u = G \star F(u, \xi) + S u_0 , \quad (3.3)$$

where \star is the space-time convolution, S is the semigroup generated by A and G is its fundamental solution. We will always assume that we are in a subcritical setting, as defined in [Hai14, Sec. 8].

It was shown in [Hai14, Sec. 8.1] that it is possible to build a regularity structure $\mathcal{T} = (\mathcal{T}, \mathcal{G})$ for a locally subcritical equation and to reformulate it as a fixed point

problem in an associated space of modelled distributions. We do not want to give a precise description of this regularity structure, see for example [Hai14, Hai15] for details in the case of Φ_3^4 . Let us just mention that we can recursively build two sets of symbols, \mathcal{F} and \mathcal{U} . The set \mathcal{F} contains $\Xi, \mathbf{1}, X_i$, as well as some of the symbols that can be built recursively from these basic building blocks by the operations

$$\tau \mapsto \mathcal{I}(\tau), \quad (\tau, \bar{\tau}) \mapsto \tau \bar{\tau}, \quad (3.4)$$

subject to the equivalences $\tau \bar{\tau} = \bar{\tau} \tau$, $\mathbf{1} \tau = \tau$, and $\mathcal{I}(X^k) = 0$. These symbols are involved in the description of the right hand side of (3.1). The set $\mathcal{U} \subset \mathcal{F}$ on the other hand contains only those symbols which are used in the description of the solution itself, which are either of the form X^k or of the form $\mathcal{I}(\tau)$ with $\tau \in \mathcal{F}$. The model space \mathcal{T} is then defined as $\text{span}\{\tau \in \mathcal{F} : |\tau| \leq r\}$ for a sufficiently large $r > 0$, the set of all (real) linear combinations of symbols in \mathcal{F} of homogeneity $|\tau| \leq r$, where $\tau \mapsto |\tau|$ is given by

$$|\mathbf{1}| = 0, \quad |X_i| = 1, \quad |\Xi| = \alpha, \quad |\mathcal{I}(\tau)| = |\tau| + \beta, \quad |\tau \bar{\tau}| = |\tau| + |\bar{\tau}|. \quad (3.5)$$

In the situation of interest, namely the Φ_3^4 model, one chooses $\beta = 2$ and $\alpha = -\frac{5}{2} - \kappa$ for some $\kappa > 0$ sufficiently small. Subcriticality then guarantees that \mathcal{T} is finite-dimensional. We will also write $\mathcal{T}_{\mathcal{U}}$ for the linear span of \mathcal{U} in \mathcal{T} .

One can also build a structure group \mathcal{G} acting on \mathcal{T} in such a way that the operation \mathcal{I} satisfies the assumptions of Definition 2.14 (corresponding to the convolution operation with the kernel K), and such that it acts on $\mathcal{T}_{\text{poly}}$ by translations as required.

Let now Z be a model realising K for \mathcal{I} , we denote by $\mathcal{R}, \mathcal{K}_{\bar{\gamma}}$ and R_{γ} the reconstruction operator, and the corresponding operators (2.28) and (3.2). We also use the notation $\mathcal{P} \stackrel{\text{def}}{=} \mathcal{K}_{\bar{\gamma}} + R_{\gamma} \mathcal{R}$ for the operator representing convolution with the heat kernel. With these notations at hand, it was shown in [Hai14] that one can associate to (3.3) the fixed point problem in $\mathcal{D}_T^{\gamma, \eta}(Z)$ given by

$$U = \mathcal{P}F(U) + Su_0, \quad (3.6)$$

for a suitable function (which we call again F) which ‘‘represents’’ the nonlinearity of the SPDE in the sense of [Hai14, Sec. 8] and which is such that $\mathcal{I}F(\tau) \in \mathcal{T}$ for every $\tau \in \mathcal{T}_{\mathcal{U}}$. In our running example, we would take

$$F(\tau) = -Q_{\leq 0}(\tau^3) + \Xi, \quad (3.7)$$

where $Q_{\leq 0}$ denotes the canonical projection onto $\mathcal{T}_{\leq 0}$ as before¹. The problem we encounter is that since we impose that our models are functions of time, there exists no model for which $\Pi_z \Xi = \xi$ with ξ a typical realisation of space-time white noise. We would like to replace (3.6) by an equivalent fixed point problem that circumvents this problem, and this is the content of the next two subsections.

¹The reason for adding this projection is to guarantee that $\mathcal{I}F$ maps $\mathcal{T}_{\mathcal{U}}$ into \mathcal{T} , since we truncated \mathcal{T} at homogeneity r .

3.2 Truncation of regularity structures

In general, as just discussed, we cannot always define a suitable inhomogeneous model for the regularity structure $\mathcal{T} = (\mathcal{T}, \mathcal{G})$, so we introduce the following truncation procedure, which amounts to simply removing the problematic symbols.

Definition 3.4. Consider a set of *generators* $\mathcal{F}^{\text{gen}} \subset \mathcal{F}$ such that $\mathcal{F}_{\text{poly}} \subset \mathcal{F}^{\text{gen}}$ and such that $\mathcal{T}^{\text{gen}} \stackrel{\text{def}}{=} \text{span}\{\tau \in \mathcal{F}^{\text{gen}} : |\tau| \leq r\} \subset \mathcal{T}$ is closed under the action of \mathcal{G} . We then define the corresponding *generating regularity structure* $\mathcal{T}^{\text{gen}} = (\mathcal{T}^{\text{gen}}, \mathcal{G})$.

Moreover, we define $\hat{\mathcal{F}}$ as the subset of \mathcal{F} generated by \mathcal{F}^{gen} via the two operations (3.4), and we assume that \mathcal{F}^{gen} was chosen in such a way that $\mathcal{U} \subset \hat{\mathcal{F}}$, with \mathcal{U} as in the previous section. Finally, we define the *truncated regularity structure* $\hat{\mathcal{T}} = (\hat{\mathcal{T}}, \mathcal{G})$ with $\hat{\mathcal{T}} \stackrel{\text{def}}{=} \text{span}\{\tau \in \hat{\mathcal{F}} : |\tau| \leq r\} \subset \mathcal{T}$.

Remark 3.5. Note that $\hat{\mathcal{T}}$ is indeed a regularity structure since $\hat{\mathcal{T}}$ is automatically closed under \mathcal{G} . This can easily be verified by induction using the definition of \mathcal{G} given in [Hai14].

A set \mathcal{F}^{gen} with these properties always exists, because one can take either $\mathcal{F}^{\text{gen}} = \mathcal{F}$ or $\mathcal{F}^{\text{gen}} = \{\Xi\} \cup \mathcal{F}_{\text{poly}}$. In both of these examples, one simply has $\hat{\mathcal{F}} = \mathcal{F}$, but in the case of (Φ_3^4) , it turns out to be convenient to make a choice for which this is not the case (see Section 7 below).

3.3 A general fixed point map

We now reformulate (3.1), with the operator A such that Assumption 3.1 is satisfied, using the regularity structure from the previous section, and show that the corresponding fixed point problem admits local solutions. For an initial condition u_0 in (3.1) with “sufficiently nice” behavior at infinity, we can define the function $S_t u_0 : \mathbf{R}^d \rightarrow \mathbf{R}$, which has a singularity at $t = 0$, where as before S_t is the semigroup generated by A . In particular, we have a precise description of its singularity, the proof of which is provided in [Hai14, Lem. 7.5]:

Lemma 3.6. *For some $\eta < 0$, let $u_0 \in \mathcal{C}^\eta(\mathbf{R}^d)$ be periodic. Then, for every $\gamma > 0$ and every $T > 0$, the map $(t, x) \mapsto S_t u_0(x)$ can be lifted to $\mathcal{D}_T^{\gamma, \eta}$ via its Taylor expansion. Furthermore, one has the bound*

$$\|S u_0\|_{\gamma, \eta; T} \lesssim \|u_0\|_{\mathcal{C}^\eta} . \quad (3.8)$$

Before reformulating (3.1), we make some assumptions on its nonlinear term F . For a regularity structure $\mathcal{T} = (\mathcal{T}, \mathcal{G})$, let $\hat{\mathcal{T}} = (\hat{\mathcal{T}}, \mathcal{G})$ be as in Definition 3.4 for a suitable set \mathcal{F}^{gen} . In what follows, we consider models on $\hat{\mathcal{T}}$ and denote by $\mathcal{D}_T^{\gamma, \eta}$ the respective spaces of modelled distributions. We also assume that we are given a function $F : \mathcal{T}_{\mathcal{U}} \rightarrow \mathcal{T}$ as above (for example (3.7)), and we make the following assumption on F .

For some fixed $\bar{\gamma} > 0$, $\eta \in \mathbf{R}$ we choose, for any model Z on $\hat{\mathcal{T}}$, elements $F_0(Z), I_0(Z) \in \mathcal{D}_T^{\bar{\gamma}, \eta}(Z)$ such that, for every z , $I_0(z) \in \hat{\mathcal{T}}$, $I_0(z) - \mathcal{I}F_0(z) \in \mathcal{T}_{\text{poly}}$ and such that, setting

$$\hat{F}(z, \tau) \stackrel{\text{def}}{=} F(z, \tau) - F_0(z) , \quad (3.9)$$

$\hat{F}(z, \cdot)$ maps $\hat{\mathcal{T}} \cap \mathcal{T}_{\mathcal{U}}$ into $\hat{\mathcal{T}}$. Here we suppressed the argument Z for conciseness by writing for example $I_0(z)$ instead of $I_0(Z)(z)$.

Remark 3.7. Since it is the *same* structure group \mathcal{G} acting on both \mathcal{T} and $\hat{\mathcal{T}}$, the condition $F_0 \in \mathcal{D}_T^{\bar{\gamma}, \eta}$ makes sense for a given model on $\hat{\mathcal{T}}$, even though $F_0(z)$ takes values in all of \mathcal{T} rather than just $\hat{\mathcal{T}}$.

Given such a choice of I_0 and F_0 and given $H : \mathbf{R}^{d+1} \rightarrow \hat{\mathcal{T}} \cap \mathcal{T}_{\mathcal{U}}$, we denote by $\hat{F}(H)$ the function

$$(\hat{F}(H))_t(x) \stackrel{\text{def}}{=} \hat{F}((t, x), H_t(x)) . \quad (3.10)$$

With this notation, we replace the problem (3.6) by the problem

$$U = \mathcal{P}\hat{F}(U) + Su_0 + I_0 . \quad (3.11)$$

This shows that one should really think of I_0 as being given by $I_0 = \mathcal{P}F_0$ since, at least formally, this would then turn (3.11) into (3.6). The advantage of (3.11) is that it makes sense for any model on $\hat{\mathcal{T}}$ and does not require a model on all of \mathcal{T} .

We then assume that \hat{F} , I_0 and F_0 satisfy the following conditions.

Assumption 3.8. *In the above context, we assume that there exists $\gamma \geq \bar{\gamma}$ such that, for every $B > 0$ there exists a constant $C > 0$ such that the bounds*

$$\begin{aligned} \|\hat{F}(H); \hat{F}(\bar{H})\|_{\bar{\gamma}, \bar{\eta}; T} &\leq C (\|H; \bar{H}\|_{\gamma, \eta; T} + \|Z; \bar{Z}\|_{\gamma; T}) , \\ \|I_0(Z); I_0(\bar{Z})\|_{\bar{\gamma}, \bar{\eta}; T} &\leq C \|Z; \bar{Z}\|_{\gamma; T} , \quad \|F_0(Z); F_0(\bar{Z})\|_{\bar{\gamma}, \bar{\eta}; T} \leq C \|Z; \bar{Z}\|_{\gamma; T} , \end{aligned} \quad (3.12)$$

hold for any two models Z, \bar{Z} with $\|Z\|_{\gamma; T} + \|\bar{Z}\|_{\gamma; T} \leq B$, and for $H \in \mathcal{D}_T^{\gamma, \eta}(Z)$, $\bar{H} \in \mathcal{D}_T^{\gamma, \eta}(\bar{Z})$ such that $\|H\|_{\gamma, \eta; T} + \|\bar{H}\|_{\gamma, \eta; T} \leq B$.

The following theorem provides the existence and uniqueness results of a local solution to this equation.

Theorem 3.9. *In the described context, let $\alpha \stackrel{\text{def}}{=} \min \hat{\mathcal{A}}$, and an abstract integration map \mathcal{I} be of order $\beta > -\alpha$. Furthermore, let the values $\gamma \geq \bar{\gamma} > 0$ and $\eta, \bar{\eta} \in \mathbf{R}$ from Assumption 3.8 satisfy $\eta < \bar{\eta} \wedge \alpha + \beta$, $\gamma < \bar{\gamma} + \beta$ and $\bar{\eta} > -\beta$.*

Then, for every model Z as above, and for every periodic $u_0 \in \mathcal{C}^\eta(\mathbf{R}^d)$, there exists a time $T_\star \in (0, +\infty]$ such that, for every $T < T_\star$ the equation (3.11) admits a unique solution $U \in \mathcal{D}_T^{\gamma, \eta}(Z)$. Furthermore, if $T_\star < \infty$, then

$$\lim_{T \rightarrow T_\star} \|\mathcal{R}_T \mathcal{S}_T(u_0, Z)_T\|_{\mathcal{C}^\eta} = \infty ,$$

where $\mathcal{S}_T : (u_0, Z) \mapsto U$ is the solution map. Finally, for every $T < T_\star$, the solution map \mathcal{S}_T is jointly Lipschitz continuous in a neighbourhood around (u_0, Z) in the sense that, for any $B > 0$ there is $C > 0$ such that, if $\bar{U} = \mathcal{S}_T(\bar{u}_0, \bar{Z})$ for some initial data (\bar{u}_0, \bar{Z}) , then one has the bound $\|U; \bar{U}\|_{\gamma, \eta; T} \leq C\delta$, provided $\|u_0 - \bar{u}_0\|_{\mathcal{C}^\eta} + \|Z; \bar{Z}\|_{\gamma; T} \leq \delta$, for any $\delta \in (0, B]$.

Proof. See [Hai14, Thm. 7.8], combined with [Hai14, Prop. 7.11]. Note that since we consider inhomogeneous models, we have no problems in evaluating $\mathcal{R}_t U_t$. \square

4 Discrete models and modelled distributions

In order to be able to consider discretisations of the equations whose solutions were provided in Section 3, we introduce the discrete counterparts of inhomogeneous models and modelled distributions. In this section we use the following notation: for $N \in \mathbf{N}$, we denote by $\varepsilon \stackrel{\text{def}}{=} 2^{-N}$ the mesh size of the grid $\Lambda_\varepsilon^d \stackrel{\text{def}}{=} (\varepsilon \mathbf{Z})^d$, and we fix some scaling $\mathfrak{s} = (\mathfrak{s}_0, 1, \dots, 1)$ of \mathbf{R}^{d+1} with an integer $\mathfrak{s}_0 > 0$.

4.1 Definitions and the reconstruction theorem

Now we define discrete analogues of the objects from Sections 2.1 and 2.2.

Definition 4.1. Given a regularity structure \mathcal{T} and $\varepsilon > 0$, a *discrete model* $(\Pi^\varepsilon, \Gamma^\varepsilon, \Sigma^\varepsilon)$ consists of the collections of maps

$$\Pi_x^{\varepsilon, t} : \mathcal{T} \rightarrow \mathbf{R}^{\Lambda_\varepsilon^d}, \quad \Gamma^{\varepsilon, t} : \Lambda_\varepsilon^d \times \Lambda_\varepsilon^d \rightarrow \mathcal{G}, \quad \Sigma_x^\varepsilon : \mathbf{R} \times \mathbf{R} \rightarrow \mathcal{G},$$

parametrised by $t \in \mathbf{R}$ and $x \in \Lambda_\varepsilon^d$, which have all the algebraic properties of their continuous counterparts in Definition 2.4, with the spatial variables restricted to the grid. Additionally, we require $(\Pi_x^{\varepsilon, t} \tau)(x) = 0$, for all $\tau \in \mathcal{T}_l$ with $l > 0$, and all $x \in \Lambda_\varepsilon^d$ and $t \in \mathbf{R}$.

We define the quantities $\|\Pi^\varepsilon\|_{\gamma; T}^{(\varepsilon)}$ and $\|\Gamma^\varepsilon\|_{\gamma; T}^{(\varepsilon)}$ to be the smallest constants C such that the bounds (2.4a) hold uniformly in $x, y \in \Lambda_\varepsilon^d$, $t \in \mathbf{R}$, $\lambda \in [\varepsilon, 1]$ and with the discrete pairing (1.7) in place of the standard one. The quantity $\|\Sigma^\varepsilon\|_{\gamma; T}^{(\varepsilon)}$ is defined as the smallest constant C such that the bounds

$$\|\Sigma_x^{\varepsilon, st} \tau\|_m \leq C \|\tau\| (|t - s|^{1/\mathfrak{s}_0} \vee \varepsilon)^{l-m}, \quad (4.1)$$

hold uniformly in $x \in \Lambda_\varepsilon^d$ and the other parameters as in (2.4b).

We measure the time regularity of Π^ε as in (2.5), by substituting the continuous objects by their discrete analogues, and by using $|t - s|^{1/\mathfrak{s}_0} \vee \varepsilon$ instead of $|t - s|^{1/\mathfrak{s}_0}$ on the right-hand side. All the other quantities $\|\cdot\|^{(\varepsilon)}$, $\|\!\|\cdot\|\!\|^{(\varepsilon)}$, etc. are defined by analogy with Remark 2.5.

Remark 4.2. The fact that $(\Pi_x^{\varepsilon, t} \tau)(x) = 0$ if $|\tau| > 0$ does not follow automatically from the discrete analogue of (2.4a) since these are only assumed to hold for test functions at scale $\lambda \geq \varepsilon$. We use this property in the proof of (4.35).

Remark 4.3. The weakening of the continuity property of $\Sigma_x^{\varepsilon, st}$ given by (4.1) will be used in the analysis of the ‘‘discrete abstract integration’’ in Section 4.2. It allows us to deal with the fact that the discrete heat kernel is discontinuous at $t = 0$, so we simply use uniform bounds on very small time scales (see [HMW14, Lem. 6.7] for a simple explanation in a related context).

For $\gamma, \eta \in \mathbf{R}$ and $T > 0$, for a discrete model $Z^\varepsilon = (\Pi^\varepsilon, \Gamma^\varepsilon, \Sigma^\varepsilon)$ on a regularity structure $\mathcal{T} = (\mathcal{T}, \mathcal{G})$, and for a function $H : (0, T] \times \Lambda_\varepsilon^d \rightarrow \mathcal{T}_{<\gamma}$, we define

$$\begin{aligned} \|H\|_{\gamma, \eta; T}^{(\varepsilon)} &\stackrel{\text{def}}{=} \sup_{t \in (0, T]} \sup_{x \in \Lambda_\varepsilon^d} \sup_{l < \gamma} |t|_\varepsilon^{(l-\eta) \vee 0} \|H_t(x)\|_l \\ &\quad + \sup_{t \in (0, T]} \sup_{\substack{x \neq y \in \Lambda_\varepsilon^d \\ |x-y| \leq 1}} \sup_{l < \gamma} \frac{\|H_t(x) - \Gamma_{xy}^{\varepsilon, t} H_t(y)\|_l}{|t|_\varepsilon^{\eta-\gamma} |x-y|^{\gamma-l}}, \end{aligned} \quad (4.2)$$

where $l \in \mathcal{A}$. Furthermore, we define the norm

$$\|H\|_{\gamma, \eta; T}^{(\varepsilon)} \stackrel{\text{def}}{=} \|H\|_{\gamma, \eta; T}^{(\varepsilon)} + \sup_{\substack{s \neq t \in (0, T] \\ |t-s| \leq |t, s|_0^{s_0}}} \sup_{x \in \Lambda_\varepsilon^d} \sup_{l < \gamma} \frac{\|H_t(x) - \Sigma_x^{\varepsilon, ts} H_s(x)\|_l}{|t, s|_\varepsilon^{\eta-\gamma} (|t-s|^{1/s_0} \vee \varepsilon)^{\gamma-l}}, \quad (4.3)$$

where the quantities $|t|_\varepsilon$ and $|t, s|_\varepsilon$ are defined in (1.8). We will call such functions H *discrete modeled distributions*.

Definition 4.4. Given a discrete model $Z^\varepsilon = (\Pi^\varepsilon, \Gamma^\varepsilon, \Sigma^\varepsilon)$ and a discrete modeled distribution H we define the *discrete reconstruction map* \mathcal{R}^ε by $\mathcal{R}_t^\varepsilon = 0$ for $t \leq 0$, and

$$(\mathcal{R}_t^\varepsilon H_t)(x) \stackrel{\text{def}}{=} (\Pi_x^{\varepsilon, t} H_t(x))(x), \quad (t, x) \in (0, T] \times \Lambda_\varepsilon^d. \quad (4.4)$$

Recalling the definition of the norms from (1.8), the following result is a discrete analogue of Theorem 2.11.

Theorem 4.5. Let \mathcal{T} be a regularity structure with $\alpha \stackrel{\text{def}}{=} \min \mathcal{A} < 0$ and $Z^\varepsilon = (\Pi^\varepsilon, \Gamma^\varepsilon, \Sigma^\varepsilon)$ be a discrete model. Then the bound

$$|\langle \mathcal{R}_t^\varepsilon H_t - \Pi_x^{\varepsilon, t} H_t(x), \varrho_x^\lambda \rangle_\varepsilon| \lesssim \lambda^\gamma |t|_\varepsilon^{\eta-\gamma} \|H\|_{\gamma, \eta; T}^{(\varepsilon)} \|\Pi^\varepsilon\|_{\gamma; T}^{(\varepsilon)},$$

holds for all discrete modeled distributions H , all $t \in (0, T]$, $x \in \Lambda_\varepsilon^d$, $\varrho \in \mathcal{B}_0^r(\mathbf{R}^d)$ with $r > -\lfloor \alpha \rfloor$, all $\lambda \in [\varepsilon, 1]$.

Let furthermore $\bar{Z}^\varepsilon = (\bar{\Pi}^\varepsilon, \bar{\Gamma}^\varepsilon, \bar{\Sigma}^\varepsilon)$ be another model for \mathcal{T} with the reconstruction operator $\bar{\mathcal{R}}_t^\varepsilon$, and let the maps Π^ε and $\bar{\Pi}^\varepsilon$ have time regularities $\delta > 0$. Then, for any two discrete modeled distributions H and \bar{H} , the maps $t \mapsto \mathcal{R}_t^\varepsilon H_t$ and $t \mapsto \bar{\mathcal{R}}_t^\varepsilon \bar{H}_t$ satisfy

$$\|\mathcal{R}^\varepsilon H\|_{\mathcal{C}_{\eta-\gamma, T}^{\delta, \alpha}}^{(\varepsilon)} \lesssim \|\Pi^\varepsilon\|_{\delta, \gamma; T}^{(\varepsilon)} (1 + \|\Sigma^\varepsilon\|_{\gamma; T}^{(\varepsilon)}) \|H\|_{\gamma, \eta; T}^{(\varepsilon)}, \quad (4.5a)$$

$$\|\mathcal{R}^\varepsilon H - \bar{\mathcal{R}}^\varepsilon \bar{H}\|_{\mathcal{C}_{\eta-\gamma, T}^{\delta, \alpha}}^{(\varepsilon)} \lesssim \|H; \bar{H}\|_{\gamma, \eta; T}^{(\varepsilon)} + \|Z^\varepsilon; \bar{Z}^\varepsilon\|_{\delta, \gamma; T}^{(\varepsilon)}, \quad (4.5b)$$

for any δ as in Theorem 2.11. Here, the norms of H and \bar{H} are defined via the models Z^ε and \bar{Z}^ε respectively, and the proportionality constants depend on ε only via $\|H\|_{\gamma, \eta; T}^{(\varepsilon)}$, $\|\bar{H}\|_{\gamma, \eta; T}^{(\varepsilon)}$, $\|Z^\varepsilon\|_{\delta, \gamma; T}^{(\varepsilon)}$ and $\|\bar{Z}^\varepsilon\|_{\delta, \gamma; T}^{(\varepsilon)}$.

Remark 4.6. To compare a discrete model $Z^\varepsilon = (\Pi^\varepsilon, \Gamma^\varepsilon, \Sigma^\varepsilon)$ to a continuous model $Z = (\Pi, \Gamma, \Sigma)$, we can define

$$\begin{aligned} \|\Pi; \Pi^\varepsilon\|_{\delta, \gamma; T}^{(\varepsilon)} &\stackrel{\text{def}}{=} \sup_{\varphi, x, \lambda, l, \tau} \sup_{t \in [-T, T]} \lambda^{-l} |\langle \Pi_x^t \tau, \varphi_x^\lambda \rangle - \langle \Pi_x^{\varepsilon, t} \tau, \varphi_x^\lambda \rangle_\varepsilon| \\ &+ \sup_{\varphi, x, \lambda, l, \tau} \sup_{\substack{s \neq t \in [-T, T] \\ |t-s| \leq 1}} \lambda^{-l+\delta} \frac{|\langle (\Pi_x^t - \Pi_x^s) \tau, \varphi_x^\lambda \rangle - \langle (\Pi_x^{\varepsilon, t} - \Pi_x^{\varepsilon, s}) \tau, \varphi_x^\lambda \rangle_\varepsilon|}{(|t-s|^{1/s_0} \vee \varepsilon)^\delta}, \end{aligned}$$

where the supremum is taken over $\varphi \in \mathcal{B}_0^r$, $x \in \Lambda_\varepsilon^d$, $\lambda \in [\varepsilon, 1]$, $l < \gamma$ and $\tau \in \mathcal{T}_l$ with $\|\tau\| = 1$. In order to compare discrete and continuous modelled distributions, we use the quantities as in (2.10), but with the respective modifications as in (4.3).

Then one can show similarly to (2.13) that for $H \in \mathcal{D}_T^{\gamma, \eta}(Z)$ and a discrete modeled distribution H^ε the maps $t \mapsto \mathcal{R}_t H_t$ and $t \mapsto \mathcal{R}_t^\varepsilon H_t^\varepsilon$ satisfy the estimate

$$\|\mathcal{R}H; \mathcal{R}^\varepsilon H^\varepsilon\|_{\mathcal{C}_{\eta-\gamma, T}^{\delta, \alpha}}^{(\varepsilon)} \lesssim \|H; H^\varepsilon\|_{\mathcal{C}_{\gamma, \eta; T}^{(\varepsilon)}} + \|Z; Z^\varepsilon\|_{\delta, \gamma; T}^{(\varepsilon)} + \varepsilon^\theta,$$

for $\tilde{\delta} > 0$ and $\theta > 0$ small enough. We will however not make use of this in the present article.

In order to prove Theorem 4.5, we need to introduce a multiresolution analysis and its discrete analogue.

4.1.1 Elements of multiresolution analysis

In this section we provide only the very basics of the multiresolution analysis, which are used in the sequel. For a more detailed introduction and for the proofs of the provided results we refer to [Dau92] and [Mey92].

One of the remarkable results of [Dau88] is that for every $r > 0$ there exists a compactly supported function $\varphi \in \mathcal{C}^r(\mathbf{R})$ (called *scaling function*) such that

$$\int_{\mathbf{R}} \varphi(x) dx = 1, \quad \int_{\mathbf{R}} \varphi(x) \varphi(x+k) dx = \delta_{0,k}, \quad k \in \mathbf{Z}, \quad (4.6)$$

where $\delta_{\cdot, \cdot}$ is the Kronecker's delta on \mathbf{Z} . Furthermore, if for $n \in \mathbf{N}$ we define the grid $\Lambda_n \stackrel{\text{def}}{=} \{2^{-n}k : k \in \mathbf{Z}\}$ and the family of functions

$$\varphi_x^n(\cdot) \stackrel{\text{def}}{=} 2^{n/2} \varphi(2^n(\cdot - x)), \quad x \in \Lambda_n, \quad (4.7)$$

then there is a finite collection of vectors $\mathcal{K} \subset \Lambda_1$ and a collection of structure constants $\{a_k : k \in \mathcal{K}\}$ such that the *refinement equation*

$$\varphi_x^n = \sum_{k \in \mathcal{K}} a_k \varphi_{x+2^{-n}k}^{n+1} \quad (4.8)$$

holds. Note that the multiplier in (4.7) preserves the L^2 -norm of the scaled functions rather than their L^1 -norm. It follows immediately from (4.6) and (4.8) that one has the identities

$$\sum_{k \in \mathcal{K}} a_k = 2^{d/2}, \quad \sum_{k \in \mathcal{K}} a_k a_{k+m} = \delta_{0,m}, \quad m \in \mathbf{Z}^d. \quad (4.9)$$

For a fixed scaling function φ , we denote by $V_n \subset L^2(\mathbf{R})$ the subspace spanned by $\{\varphi_x^n : x \in \Lambda_n\}$. Then the relation (4.8) ensures the inclusion $V_n \subset V_{n+1}$ for every n . It turns out that there is a compactly supported function $\psi \in \mathcal{C}^r(\mathbf{R})$ (called *wavelet function*) such that the space V_n^\perp , which is the orthogonal complement of V_n in V_{n+1} , is given by

$$V_n^\perp = \text{span}\{\psi_x^n : x \in \Lambda_n\},$$

where ψ_x^n is as in (4.8). Moreover, there are constants $\{b_k : k \in \mathcal{K}\}$, such that the *wavelet equation* holds:

$$\psi_x^n = \sum_{k \in \mathcal{K}} b_k \varphi_{x+2^{-n}k}^{n+1}. \quad (4.10)$$

One more useful property of the wavelet function is that it has vanishing moments, in the sense that the identity

$$\int_{\mathbf{R}} \psi(x) x^m dx = 0 \quad (4.11)$$

holds for all $m \in \mathbf{N}$ such that $m \leq r$.

There is a standard generalization of scaling and wavelet functions to \mathbf{R}^d , namely for $n \geq 0$ and $x = (x_1, \dots, x_d) \in \Lambda_n^d$ we define

$$\varphi_x^n(y) \stackrel{\text{def}}{=} \varphi_{x_1}^n(y_1) \cdots \varphi_{x_d}^n(y_d), \quad y = (y_1, \dots, y_d) \in \mathbf{R}^d.$$

For these scaling functions we also define V_n as the closed subspace in L^2 spanned by $\{\varphi_x^n : x \in \Lambda_n^d\}$. Then there is a finite set Ψ of functions on \mathbf{R}^d such that the space $V_n^\perp \stackrel{\text{def}}{=} V_{n+1} \setminus V_n$ is a span of $\{\psi_x^n : \psi \in \Psi, x \in \Lambda_n^d\}$, where we define the scaled function ψ_x^n by

$$\psi_x^n(y) \stackrel{\text{def}}{=} 2^{nd/2} \psi(2^n(y_1 - x_1), \dots, 2^n(y_d - x_d)).$$

All the results mentioned above can be literally translated from \mathbf{R} to \mathbf{R}^d , but of course with $\mathcal{K} \subset \Lambda_1^d$ and with different structure constants $\{a_k : k \in \mathcal{K}\}$ and $\{b_k : k \in \mathcal{K}\}$.

4.1.2 An analogue of the multiresolution analysis on the grid

In this section we will develop an analogue of the multiresolution analysis which will be useful for working with functions defined on a dyadic grid. Our construction

agrees with the standard discrete wavelets on gridpoints, but also extends off the grid. To this end, we use the notation of Section 4.1.1. We recall furthermore that we use $\varepsilon = 2^{-N}$ for a fixed $N \in \mathbf{N}$.

Let us fix a scaling function $\varphi \in \mathcal{C}_0^r(\mathbf{R})$, for some integer $r > 0$, as in Section 4.1.1. For integers $0 \leq n \leq N$ we define the functions

$$\varphi_x^{N,n}(\cdot) \stackrel{\text{def}}{=} 2^{Nd/2} \langle \varphi^N, \varphi_x^n \rangle, \quad x \in \Lambda_n^d. \quad (4.12)$$

One has that $\varphi_x^{N,n} \in \mathcal{C}^r(\mathbf{R}^d)$, it is supported in a ball of radius $\mathcal{O}(2^{-n})$ centered at x , it has the same scaling properties as φ_x^n , and it satisfies

$$\varphi_x^{N,N}(y) = 2^{Nd/2} \delta_{x,y}, \quad x, y \in \Lambda_N^d, \quad (4.13)$$

where $\delta_{\cdot,\cdot}$ is the Kronecker's delta on Λ_N^d . The last property follows from (4.6). Furthermore, it follows from (4.8) that for $n < N$ these functions satisfy the refinement identity

$$\varphi_x^{N,n} = \sum_{k \in \mathcal{K}} a_k \varphi_{x+2^{-n}k}^{N,n+1}, \quad (4.14)$$

with the same structure constants $\{a_k : k \in \mathcal{K}\}$ as for the functions φ_x^n . One more consequence of (4.6) is

$$2^{-Nd} \sum_{y \in \Lambda_N^d} \varphi_x^{N,n}(y) = 2^{-nd/2},$$

which obviously holds for $n = N$, and for $n < N$ it can be proved by induction, using (4.14) and (4.9).

The functions $\varphi_x^{N,n}$ inherit many of the crucial properties of the functions φ_x^n , which allows us to use them in the multiresolution analysis. In particular, for $n < N$ and $\psi \in \Psi$ (the set of wavelet functions, introduced in Section 4.1.1), we can define the functions

$$\psi_x^{N,n}(\cdot) \stackrel{\text{def}}{=} 2^{Nd/2} \langle \varphi^N, \psi_x^n \rangle, \quad x \in \Lambda_n^d,$$

whose properties are similar to those of ψ_x^n . For example, $\psi_x^{N,n} \in \mathcal{C}^r(\mathbf{R})$, and it has the same scaling and support properties as ψ_x^n . Furthermore, it follows from (4.10) that for $n < N$ the following identity holds

$$\psi_x^{N,n} = \sum_{k \in \mathcal{K}} b_k \varphi_{x+2^{-n}k}^{N,n+1}, \quad (4.15)$$

with the same constants $\{b_k : k \in \mathcal{K}\}$.

4.1.3 Proof of the discrete reconstruction theorem

With the help of the discrete analogue of the multiresolution analysis introduced in the previous section we are ready to prove Theorem 4.5.

Proof of Theorem 4.5. We take a compactly supported scaling function $\varphi \in \mathcal{C}^r(\mathbf{R}^d)$ of regularity $r > -\lfloor \alpha \rfloor$, where α is as in the statement of the theorem, and build the functions $\varphi_x^{N,n}$ as in (4.12). Furthermore, we define the discrete functions $\zeta_x^{\varepsilon,t} \stackrel{\text{def}}{=} \Pi_x^{\varepsilon,t} H_t(x)$ and $\zeta_{xy}^{\varepsilon,t} \stackrel{\text{def}}{=} \zeta_y^{\varepsilon,t} - \zeta_x^{\varepsilon,t}$. Then from Definition 4.1 we obtain

$$\begin{aligned} |\langle \zeta_{xy}^{\varepsilon,t}, \varphi_y^{N,n} \rangle_\varepsilon| &\lesssim \|\Pi^{\varepsilon,t}\|_{\gamma;T}^{(\varepsilon)} \sum_{l \in [\alpha, \gamma] \cap \mathcal{A}} 2^{-nd/2-ln} \|H_t(y) - \Gamma_{yx}^{\varepsilon,t} H_t(x)\|_l \\ &\lesssim \|\Pi^{\varepsilon,t}\|_{\gamma;T}^{(\varepsilon)} \|H\|_{\gamma,\eta;T}^{(\varepsilon)} |t|_\varepsilon^{\eta-\gamma} \sum_{l \in [\alpha, \gamma] \cap \mathcal{A}} 2^{-nd/2-ln} |y-x|^{\gamma-l} \\ &\lesssim \|\Pi^{\varepsilon,t}\|_{\gamma;T}^{(\varepsilon)} \|H\|_{\gamma,\eta;T}^{(\varepsilon)} |t|_\varepsilon^{\eta-\gamma} 2^{-nd/2-\alpha n} |y-x|^{\gamma-\alpha}, \end{aligned} \quad (4.16)$$

which holds as soon as $|x-y| \geq 2^{-n}$. Moreover, we define

$$\mathcal{R}_t^{\varepsilon,n} H_t \stackrel{\text{def}}{=} \sum_{y \in \Lambda_n^d} \langle \zeta_y^{\varepsilon,t}, \varphi_y^{N,n} \rangle_\varepsilon \varphi_y^{N,n}.$$

It follows from the property (4.13) that $\mathcal{R}_t^\varepsilon H_t = \mathcal{R}_t^{\varepsilon,N} H_t$ and $\Pi_x^{\varepsilon,t} H_t(x) = \mathcal{P}_{\varepsilon,N}(\zeta_x^{\varepsilon,t})$ (recall that $\varepsilon = 2^{-N}$), where the operator $\mathcal{P}_{\varepsilon,n}$ is defined by

$$\mathcal{P}_{\varepsilon,n}(\zeta) \stackrel{\text{def}}{=} \sum_{y \in \Lambda_n^d} \langle \zeta, \varphi_y^{N,n} \rangle_\varepsilon \varphi_y^{N,n}.$$

This allows us to choose $n_0 \geq 0$ to be the smallest integer such that $2^{-n_0} \leq \lambda$ and rewrite

$$\begin{aligned} \mathcal{R}_t^\varepsilon H_t - \Pi_x^{\varepsilon,t} H_t(x) &= (\mathcal{R}_t^{\varepsilon,n_0} H_t - \mathcal{P}_{\varepsilon,n_0}(\zeta_x^{\varepsilon,t})) \\ &\quad + \sum_{n=n_0}^{N-1} \left(\mathcal{R}_t^{\varepsilon,n+1} H_t - \mathcal{P}_{\varepsilon,n+1}(\zeta_x^{\varepsilon,t}) - \mathcal{R}_t^{\varepsilon,n} H_t + \mathcal{P}_{\varepsilon,n}(\zeta_x^{\varepsilon,t}) \right). \end{aligned} \quad (4.17)$$

The first term on the right hand side yields

$$\langle \mathcal{R}_t^{\varepsilon,n_0} H_t - \mathcal{P}_{\varepsilon,n_0}(\zeta_x^{\varepsilon,t}), \varrho_x^\lambda \rangle_\varepsilon = \sum_{y \in \Lambda_{n_0}^d} \langle \zeta_{xy}^{\varepsilon,t}, \varphi_y^{N,n_0} \rangle_\varepsilon \langle \varphi_y^{N,n_0}, \varrho_x^\lambda \rangle_\varepsilon. \quad (4.18)$$

Using (4.16) and the bound $|\langle \varphi_y^{N,n_0}, \varrho_x^\lambda \rangle_\varepsilon| \lesssim 2^{n_0 d/2}$, we obtain

$$|\langle \mathcal{R}_t^{\varepsilon,n_0} H_t - \mathcal{P}_{\varepsilon,n_0}(\zeta_x^{\varepsilon,t}), \varrho_x^\lambda \rangle_\varepsilon| \lesssim \|\Pi^{\varepsilon,t}\|_{\gamma;T}^{(\varepsilon)} \|H\|_{\gamma,\eta;T}^{(\varepsilon)} |t|_\varepsilon^{\eta-\gamma} 2^{-\gamma n_0}.$$

Here, we have also used $|x-y| \lesssim 2^{-n_0}$ in the sum in (4.18), and the fact that only a finite number of points $y \in \Lambda_{n_0}^d$ contribute to this sum.

Now we will bound each term in the sum in (4.17). Using (4.14) and (4.15), we can write

$$\mathcal{R}_t^{\varepsilon,n+1} H_t - \mathcal{P}_{\varepsilon,n+1}(\zeta_x^{\varepsilon,t}) - \mathcal{R}_t^{\varepsilon,n} H_t + \mathcal{P}_{\varepsilon,n}(\zeta_x^{\varepsilon,t}) = g_{t,n}^\varepsilon + h_{t,n}^\varepsilon,$$

where $g_{t,n}^\varepsilon$ is defined by

$$g_{t,n}^\varepsilon = \sum_{y \in \Lambda_n^d} \sum_{k \in \mathcal{K}} a_k \langle \zeta_{y,y+2^{-n}k}^{\varepsilon,t}, \varphi_{y+2^{-n}k}^{N,n+1} \rangle_\varepsilon \varphi_y^{N,n}$$

and the constants $\{a_k : k \in \mathcal{K}\}$ are from (4.14). For $h_{t,n}^\varepsilon$ we have the identity

$$h_{t,n}^\varepsilon = \sum_{y \in \Lambda_{n+1}^d} \sum_{k \in \mathcal{K}} \sum_{\psi \in \Psi} b_k \langle \zeta_{xy}^{\varepsilon,t}, \varphi_y^{N,n+1} \rangle_\varepsilon \psi_{y-2^{-n}k}^{N,n}. \quad (4.19)$$

Moreover, the following bounds, for $n \in [n_0, N]$, follow from the properties of the functions φ_x^n and ψ_x^n :

$$|\langle \varphi_y^{N,n}, \varrho_x^\lambda \rangle_\varepsilon| \lesssim 2^{n_0 d/2} 2^{-(n-n_0)d/2}, \quad |\langle \psi_y^{N,n}, \varrho_x^\lambda \rangle_\varepsilon| \lesssim 2^{n_0 d/2} 2^{-(n-n_0)(r+d/2)}.$$

Using them and (4.16), we obtain a bound on $g_{t,n}^\varepsilon$:

$$\begin{aligned} |\langle g_{t,n}^\varepsilon, \varrho_x^\lambda \rangle_\varepsilon| &\lesssim \sum_{y \in \Lambda_n^d} \sum_{k \in \mathcal{K}} |\langle \zeta_{y,y+2^{-n}k}^{\varepsilon,t}, \varphi_{y+2^{-n}k}^{N,n+1} \rangle_\varepsilon| |\langle \varphi_y^{N,n}, \varrho_x^\lambda \rangle_\varepsilon| \\ &\lesssim \|\Pi^\varepsilon\|_{\gamma;T}^{(\varepsilon)} \|H\|_{\gamma,\eta;T}^{(\varepsilon)} |t|_\varepsilon^{\eta-\gamma} 2^{-\gamma n}, \end{aligned}$$

where we have used $|x-y| \lesssim 2^{-n}$ in the sum. Summing these bounds over $n \in [n_0, N]$, we obtain a bound of the required order. Similarly, we obtain the following bound on (4.19):

$$|\langle h_{t,n}^\varepsilon, \varrho_x^\lambda \rangle_\varepsilon| \lesssim \|\Pi^\varepsilon\|_{\gamma;T}^{(\varepsilon)} \|H\|_{\gamma,\eta;T}^{(\varepsilon)} |t|_\varepsilon^{\eta-\gamma} 2^{-\gamma n_0} 2^{-(n-n_0)(r+\alpha)},$$

which gives the required bound after summing over $n \in [n_0, N]$. In this estimate we have used the fact that $|y-x| \approx 2^{-n_0}$ in the sum in (4.19).

The bounds (4.5) can be shown similarly to (2.12) and (2.13). \square

4.2 Convolutions with discrete kernels

In this section we describe on the abstract level convolutions with discrete kernels.

Definition 4.7. We say that a function $K^\varepsilon : \mathbf{R} \times \Lambda_\varepsilon^d \rightarrow \mathbf{R}$ is regularising of order $\beta > 0$, if one can find functions $K^{(\varepsilon,n)} : \mathbf{R}^{d+1} \rightarrow \mathbf{R}$ and $\hat{K}^\varepsilon : \mathbf{R} \times \Lambda_\varepsilon^d \rightarrow \mathbf{R}$ such that

$$K^\varepsilon = \sum_{n=0}^{N-1} K^{(\varepsilon,n)} + \hat{K}^\varepsilon \stackrel{\text{def}}{=} \bar{K}^\varepsilon + \check{K}^\varepsilon, \quad (4.20)$$

where the function $K^{(\varepsilon,n)}$ has the same support and bounds as the function $K^{(n)}$ in Definition 2.16, for some $c, r > 0$, and furthermore, for $k \in \mathbf{N}^{d+1}$ such that $|k|_s \leq r$, it satisfies

$$\int_{\mathbf{R} \times \Lambda_\varepsilon^d} z^k K^{(\varepsilon,n)}(z) dz = 0. \quad (4.21)$$

The function \hat{K}^ε is supported in $\{z \in \mathbf{R} \times \Lambda_\varepsilon^d : \|z\|_s \leq c\varepsilon\}$ and satisfies (4.21) with $k = 0$ and

$$\sup_{z \in \mathbf{R} \times \Lambda_\varepsilon^d} |\hat{K}^\varepsilon(z)| \leq C\varepsilon^{-|s|+\beta}. \quad (4.22)$$

Now, we will define how a discrete model acts on an abstract integration map.

Definition 4.8. Let \mathcal{I} be an abstract integration map of order β for a regularity structure $\mathcal{S} = (\mathcal{T}, \mathcal{G})$, let $Z^\varepsilon = (\Pi^\varepsilon, \Gamma^\varepsilon, \Sigma^\varepsilon)$ be a discrete model, and let K^ε be regularising of order β with $r > -\lfloor \min \mathcal{A} \rfloor$. Let furthermore \bar{K}^ε and \mathring{K}^ε be as in (4.20). We define $\bar{\mathcal{J}}^\varepsilon$ on the grid in the same way as its continuous analogue in (2.26), but using \bar{K}^ε instead of K and using the discrete objects instead of their continuous counterparts. Moreover, we define

$$\mathring{\mathcal{J}}_{t,x}^\varepsilon \tau \stackrel{\text{def}}{=} \mathbf{1} \int_{\mathbf{R}} \langle \Pi_x^{\varepsilon,s} \Sigma_x^{\varepsilon,st} \tau, \mathring{K}_{t-s}^\varepsilon(x - \cdot) \rangle_\varepsilon ds ,$$

and $\mathcal{J}_{t,x}^\varepsilon \stackrel{\text{def}}{=} \bar{\mathcal{J}}_{t,x}^\varepsilon + \mathring{\mathcal{J}}_{t,x}^\varepsilon$. We say that Z^ε realises K^ε for \mathcal{I} if the identities (2.25) and (2.27) hold for the corresponding discrete objects. As before, these two identities should be thought of as providing the definitions of $\Gamma_{xy}^{\varepsilon,t} \mathcal{I} \tau$ and $\Sigma_x^{\varepsilon,st} \mathcal{I} \tau$ via $\Gamma_{xy}^{\varepsilon,t}$ and $\Sigma_x^{\varepsilon,st} \tau$.

For a discrete modeled distribution H , we define $\bar{\mathcal{N}}_\gamma^\varepsilon H$ as in (2.29), but using the discrete objects instead of the continuous ones, and using the kernel \bar{K}^ε instead of K . Furthermore, we define the term containing \mathring{K}^ε by

$$(\mathring{\mathcal{N}}_\gamma^\varepsilon H)_t(x) \stackrel{\text{def}}{=} \mathbf{1} \int_{\mathbf{R}} \langle \mathcal{R}_s^\varepsilon H_s - \Pi_x^{\varepsilon,s} \Sigma_x^{\varepsilon,st} H_t(x), \mathring{K}_{t-s}^\varepsilon(x - \cdot) \rangle_\varepsilon ds , \quad (4.23)$$

and we set $\mathcal{N}_\gamma^\varepsilon H \stackrel{\text{def}}{=} \bar{\mathcal{N}}_\gamma^\varepsilon H + \mathring{\mathcal{N}}_\gamma^\varepsilon H$. Finally, we define the discrete analogue of (2.28) by

$$(\mathcal{K}_\gamma^\varepsilon H)_t(x) \stackrel{\text{def}}{=} \mathcal{I} H_t(x) + \mathcal{J}_{t,x}^\varepsilon H_t(x) + (\mathcal{N}_\gamma^\varepsilon H)_t(x) . \quad (4.24)$$

Our definition is consistent thanks to the following two lemmas.

Lemma 4.9. *In the setting of Definition 4.8, let $\min \mathcal{A} + \beta > 0$. Then all the algebraic relations of Definition 4.1 hold for the symbol $\mathcal{I} \tau$. Moreover, for $\delta > 0$ sufficiently small and for any $l \in \mathcal{A}$ and $\tau \in \mathcal{T}_l$ such that $l + \beta \notin \mathbf{N}$ and $\|\tau\| = 1$, one has the bounds*

$$|\langle \Pi_x^{\varepsilon,t} \mathcal{I} \tau, \varphi_x^\lambda \rangle_\varepsilon| \lesssim \lambda^{l+\beta} \|\Pi^\varepsilon\|_{l;T}^{(\varepsilon)} \|\Sigma^\varepsilon\|_{l;T}^{(\varepsilon)} (1 + \|\Gamma^\varepsilon\|_{l;T}^{(\varepsilon)}) , \quad (4.25)$$

$$\frac{|\langle (\Pi_x^{\varepsilon,t} - \Pi_x^{\varepsilon,s}) \mathcal{I} \tau, \varphi_x^\lambda \rangle_\varepsilon|}{(|t-s|^{1/s_0} \vee \varepsilon)^\delta} \lesssim \lambda^{l+\beta-\delta} \|\Pi^\varepsilon\|_{\delta;l;T}^{(\varepsilon)} \|\Sigma^\varepsilon\|_{l;T}^{(\varepsilon)} (1 + \|\Gamma^\varepsilon\|_{l;T}^{(\varepsilon)}) , \quad (4.26)$$

uniformly over $x \in \Lambda_\varepsilon^d$, $s, t \in [-T, T]$, $\lambda \in [\varepsilon, 1]$ and $\varphi \in \mathcal{B}_0^r(\mathbf{R}^d)$.

Proof. The algebraic properties of the models for the symbol $\mathcal{I} \tau$ follow easily from Definition 4.8. In order to prove (4.25), we will consider the terms in (2.25) containing \mathring{K}^ε separately from the others. To this end, we define

$$(\mathring{\Pi}_x^{\varepsilon,t} \mathcal{I} \tau)(y) \stackrel{\text{def}}{=} \int_{\mathbf{R}} \langle \Pi_x^{\varepsilon,s} \Sigma_x^{\varepsilon,st} \tau, \mathring{K}_{t-s}^\varepsilon(y - \cdot) - \mathring{K}_{t-s}^\varepsilon(x - \cdot) \rangle_\varepsilon ds , \quad (4.27)$$

$$(\bar{\Pi}_x^{\varepsilon,t} \mathcal{I}\tau)(y) \stackrel{\text{def}}{=} (\Pi_x^{\varepsilon,t} - \overset{\circ}{\Pi}_x^{\varepsilon,t})(\mathcal{I}\tau)(y).$$

Furthermore, for $x, y \in \Lambda_\varepsilon^d$ we use the assumption $0^0 \stackrel{\text{def}}{=} 1$ and set

$$T_{xy}^l K_t^{(\varepsilon,n)}(\cdot) \stackrel{\text{def}}{=} K_t^{(\varepsilon,n)}(y - \cdot) - \sum_{|k|_s < l + \beta} \frac{(0, y - x)^k}{k!} D^k K_t^{(\varepsilon,n)}(x - \cdot).$$

Using Definitions 4.1 and 4.7 and acting as in the proof of [Hai14, Lem. 5.19], we can obtain the following analogues of the bounds [Hai14, Eq. 5.33]:

$$\begin{aligned} |\langle \Pi_x^{\varepsilon,r} \Sigma_x^{\varepsilon,rt} \tau, T_{xy}^l K_{t-r}^{(\varepsilon,n)} \rangle_\varepsilon| &\lesssim \sum_{\zeta > 0} |y - x|^{l+\beta+\zeta} 2^{(s_0+\zeta)n} \mathbf{1}_{|t-r| \lesssim 2^{-s_0 n}}, \\ \left| \int_{\Lambda_\varepsilon^d} \langle \Pi_x^{\varepsilon,r} \Sigma_x^{\varepsilon,rt} \tau, T_{xy}^l K_{t-r}^{(\varepsilon,n)} \rangle_\varepsilon \varphi_x^\lambda(y) dy \right| &\lesssim \sum_{\zeta > 0} \lambda^{l+\beta-\zeta} 2^{(s_0-\zeta)n} \mathbf{1}_{|t-r| \lesssim 2^{-s_0 n}}, \end{aligned} \quad (4.28)$$

for $\varepsilon \leq |y - x| \leq 1$, $\lambda \in [\varepsilon, 1]$, with ζ taking a finite number of values and with the proportionality constants as in (4.25). Integrating these bounds in the time variable r and using the first bound in (4.28) in the case $|y - x| \leq 2^{-n}$ and the second bound in the case $2^{-n} \leq \lambda$, we obtain the required estimate on $\langle \bar{\Pi}_x^{\varepsilon,t} \mathcal{I}\tau, \varphi_x^\lambda \rangle_\varepsilon$.

In order to bound $(\bar{\Pi}_x^{\varepsilon,t} - \bar{\Pi}_x^{\varepsilon,s}) \mathcal{I}\tau$, we consider two cases $|t - s| \geq 2^{-s_0 n}$ and $|t - s| < 2^{-s_0 n}$. In the first case we estimate $\bar{\Pi}_x^{\varepsilon,t} \mathcal{I}\tau$ and $\bar{\Pi}_x^{\varepsilon,s} \mathcal{I}\tau$ separately using (4.28), and obtain the required bound, if $\delta > 0$ is sufficiently small. In the case $|t - s| < 2^{-s_0 n}$ we write

$$\begin{aligned} &\langle \Pi_x^{\varepsilon,r} \Sigma_x^{\varepsilon,rt} \tau, T_{xy}^l K_{t-r}^{(\varepsilon,n)} \rangle_\varepsilon - \langle \Pi_x^{\varepsilon,r} \Sigma_x^{\varepsilon,rs} \tau, T_{xy}^l K_{s-r}^{(\varepsilon,n)} \rangle_\varepsilon \\ &= \langle \Pi_x^{\varepsilon,r} \Sigma_x^{\varepsilon,rs} (\Sigma_x^{\varepsilon,st} - 1) \tau, T_{xy}^l K_{t-r}^{(\varepsilon,n)} \rangle_\varepsilon + \langle \Pi_x^{\varepsilon,r} \Sigma_x^{\varepsilon,rs} \tau, T_{xy}^l (K_{t-r}^{(\varepsilon,n)} - K_{s-r}^{(\varepsilon,n)}) \rangle_\varepsilon, \end{aligned}$$

and estimate each of these terms similarly to (4.28), which gives the required bound for sufficiently small $\delta > 0$.

It is only left to prove the required bounds for $\overset{\circ}{\Pi}_x^{\varepsilon,t}(\mathcal{I}\tau)$. It follows immediately from Definition 4.1 that $|(\Pi_x^{\varepsilon,t} a)(x)| \lesssim \|a\| \varepsilon^\zeta$, for $a \in \mathcal{T}_\zeta$. Hence, using the properties (2.2) and (2.3) we obtain

$$\begin{aligned} \int_{\mathbf{R}} |\langle \Pi_x^{\varepsilon,s} \Sigma_x^{\varepsilon,st} \tau, \overset{\circ}{K}_{t-s}^\varepsilon(y - \cdot) \rangle_\varepsilon| ds &= \int_{\mathbf{R}} |\langle \Pi_y^{\varepsilon,s} \Sigma_y^{\varepsilon,st} \Gamma_{yx}^{\varepsilon,t} \tau, \overset{\circ}{K}_{t-s}^\varepsilon(y - \cdot) \rangle_\varepsilon| ds \\ &\lesssim \sum_{\zeta \leq l} \varepsilon^{\zeta+\beta} |y - x|^{l-\zeta}, \end{aligned} \quad (4.29)$$

where $\zeta \in \mathcal{A}$. Similarly, the second term in (4.27) is bounded by $\varepsilon^{l+\beta}$, implying that if $\lambda \geq \varepsilon$ and $\min \mathcal{A} + \beta > 0$, then one has

$$|\langle \overset{\circ}{\Pi}_x^{\varepsilon,t} \mathcal{I}\tau, \varphi_x^\lambda \rangle_\varepsilon| \lesssim \sum_{\zeta \leq l} \varepsilon^{\zeta+\beta} \lambda^{l-\zeta} \lesssim \lambda^{l+\beta}, \quad (4.30)$$

which finishes the proof of (4.25). In order to complete the proof of (4.26), we use (4.29) and brutally bound

$$\begin{aligned} |\langle (\mathring{\Pi}_x^{\varepsilon,t} - \mathring{\Pi}_x^{\varepsilon,s})\mathcal{I}\tau, \varphi_x^\lambda \rangle_\varepsilon| &\leq |\langle \mathring{\Pi}_x^{\varepsilon,t}\mathcal{I}\tau, \varphi_x^\lambda \rangle_\varepsilon| + |\langle \mathring{\Pi}_x^{\varepsilon,s}\mathcal{I}\tau, \varphi_x^\lambda \rangle_\varepsilon| \\ &\lesssim \sum_{\zeta \leq l} \varepsilon^{\zeta+\beta} |y-x|^{l-\zeta} \lesssim (|t-s|^{1/s_0} \vee \varepsilon)^{\tilde{\delta}} \sum_{\zeta \leq l} \varepsilon^{\zeta+\beta-\tilde{\delta}} |y-x|^{l-\zeta}, \end{aligned}$$

from which we obtain the required bound in the same way as before, as soon as $\delta \in (0, \min \mathcal{A} + \beta)$. \square

The following lemma provides a relation between \mathcal{J}^ε and the operators $\Gamma^\varepsilon, \Sigma^\varepsilon$.

Lemma 4.10. *In the setting of Lemma 4.9, the operators*

$$\mathcal{J}_{xy}^{\varepsilon,t} \stackrel{\text{def}}{=} \mathcal{J}_{t,x}^\varepsilon \Gamma_{xy}^{\varepsilon,t} - \Gamma_{xy}^{\varepsilon,t} \mathcal{J}_{t,y}^\varepsilon, \quad \mathcal{J}_x^{\varepsilon,st} \stackrel{\text{def}}{=} \mathcal{J}_{s,x}^\varepsilon \Sigma_x^{\varepsilon,st} - \Sigma_x^{\varepsilon,st} \mathcal{J}_{t,x}^\varepsilon, \quad (4.31)$$

with $s, t \in \mathbf{R}$ and $x, y \in \Lambda_\varepsilon^d$, satisfy the following bounds:

$$\begin{aligned} |(\mathcal{J}_{xy}^{\varepsilon,t} \tau)_k| &\lesssim \|\Pi^\varepsilon\|_{l;T}^{(\varepsilon)} \|\Sigma^\varepsilon\|_{l;T}^{(\varepsilon)} (1 + \|\Gamma^\varepsilon\|_{l;T}^{(\varepsilon)}) |x-y|^{l+\beta-|k|_s}, \\ |(\mathcal{J}_x^{\varepsilon,st} \tau)_k| &\lesssim \|\Pi^\varepsilon\|_{l;T}^{(\varepsilon)} \|\Sigma^\varepsilon\|_{l;T}^{(\varepsilon)} (1 + \|\Gamma^\varepsilon\|_{l;T}^{(\varepsilon)}) (|t-s|^{1/s_0} \vee \varepsilon)^{l+\beta-|k|_s}, \end{aligned} \quad (4.32)$$

for τ as in Lemma 4.9, for any $k \in \mathbf{N}^{d+1}$ such that $|k|_s < l + \beta$, and for $(\cdot)_k$ being the multiplier of X^k . In particular, the required bounds on $\Gamma^\varepsilon \mathcal{I}\tau$ and $\Sigma^\varepsilon \mathcal{I}\tau$ from Definition 4.1 hold.

Proof. The bounds on the parts of $\mathcal{J}_{xy}^{\varepsilon,t} \tau$ and $\mathcal{J}_x^{\varepsilon,st} \tau$ not containing \mathring{K}^ε can be obtained as in [Hai14, Lem. 5.21], where the bound on the right-hand side of (4.32) comes from the fact that the scaling of the kernels $K^{(\varepsilon,n)}$ in (4.20) does not go below ε . The contributions to (4.31) from the kernel \mathring{K}^ε come via the terms $\mathring{\mathcal{J}}_{t,x}^\varepsilon \Gamma_{xy}^{\varepsilon,t}, \mathring{\mathcal{J}}_{t,y}^\varepsilon, \mathring{\mathcal{J}}_{s,x}^\varepsilon \Sigma_x^{\varepsilon,st}$ and $\mathring{\mathcal{J}}_{t,x}^\varepsilon$. We can bound all of them separately, similarly to (4.29), and use $|x-y| \geq \varepsilon$ and $|t-s|^{1/s_0} \vee \varepsilon \geq \varepsilon$ to estimate the powers of ε . Since all of these powers are positive by assumption, this yields the required bounds.

Now, we will prove the bound on $\Gamma^\varepsilon \mathcal{I}\tau$ required by Definition 4.1. For $m < l$ such that $m \notin \mathbf{N}$, (2.27) yields

$$\|\Gamma_{xy}^{\varepsilon,t} \mathcal{I}\tau\|_m = \|\mathcal{I}(\Gamma_{xy}^{\varepsilon,t} \tau)\|_m \leq \|\Gamma_{xy}^{\varepsilon,t} \tau\|_{m-\beta} \lesssim |y-x|^{l+\beta-m},$$

where we have used the properties of \mathcal{I} . Similarly, we can bound $\|\Sigma_x^{\varepsilon,st} \mathcal{I}\tau\|_m$. Furthermore, since the map \mathcal{I} does not produce elements of integer homogeneity, we have for $m \in \mathbf{N}$,

$$\|\Gamma_{xy}^{\varepsilon,t} \mathcal{I}\tau\|_m = \|\mathcal{J}_{xy}^{\varepsilon,t}\|_m \lesssim |y-x|^{l+\beta-m},$$

where the last bound we have proved above. In the same way we can obtain the required bound on $\|\Sigma_x^{\varepsilon,st} \mathcal{I}\tau\|_m$. \square

Remark 4.11. If $(\Pi^\varepsilon, \Gamma^\varepsilon, \Sigma^\varepsilon)$ is a discrete model on \mathcal{T}^{gen} , then there is a canonical way to extend it to a discrete model on $\hat{\mathcal{T}}$. Since the symbols from $\hat{\mathcal{T}}$ are “generated” by \mathcal{F}^{gen} , we only have to define the actions of $\Pi^\varepsilon, \Gamma^\varepsilon$ and Σ^ε on the symbols $\tau\bar{\tau}$ and $\mathcal{I}\tau \in \hat{\mathcal{T}} \setminus \mathcal{F}^{\text{gen}}$ with $\tau, \bar{\tau} \in \hat{\mathcal{T}}$, so that the extension of the model to $\hat{\mathcal{T}}$ will follow by induction. For the product $\tau\bar{\tau}$, we set

$$(\Pi_x^{\varepsilon,t}\tau\bar{\tau})(y) = (\Pi_x^{\varepsilon,t}\tau)(y) (\Pi_x^{\varepsilon,t}\bar{\tau})(y), \quad (4.33a)$$

$$\Sigma_x^{\varepsilon,st}\tau\bar{\tau} = (\Sigma_x^{\varepsilon,st}\tau) (\Sigma_x^{\varepsilon,st}\bar{\tau}), \quad \Gamma_{xy}^{\varepsilon,t}\tau\bar{\tau} = (\Gamma_{xy}^{\varepsilon,t}\tau) (\Gamma_{xy}^{\varepsilon,t}\bar{\tau}). \quad (4.33b)$$

For the symbol $\mathcal{I}\tau$ we define the actions of the maps $(\Pi^\varepsilon, \Gamma^\varepsilon, \Sigma^\varepsilon)$ by the identities (2.25) and (2.27). However, even if the family of models satisfy analytic bounds uniformly in ε on \mathcal{T}^{gen} , this is not necessarily true for its extension to $\hat{\mathcal{T}}$.

The structure of the canonical extension of a discrete model will be important for us. That is why we make the following definition.

Definition 4.12. We call a discrete model $Z^\varepsilon = (\Pi^\varepsilon, \Gamma^\varepsilon, \Sigma^\varepsilon)$ defined on $\hat{\mathcal{T}}$ *admissible*, if it satisfies the identities (4.33b) and furthermore realises K^ε for \mathcal{I} .

Remark 4.13. If $M \in \mathfrak{R}$ is a renormalisation map as in Section 3.1, such that $M\hat{\mathcal{T}} \subset \hat{\mathcal{T}}$, and if $Z^\varepsilon = (\Pi^\varepsilon, \Gamma^\varepsilon, \Sigma^\varepsilon)$ is an admissible model, then we can define a renormalised discrete model MZ^ε as in [Hai14, Sec. 8.3], which is also admissible.

The following result is a discrete analogue of Theorem 2.21.

Theorem 4.14. *For a regularity structure $\mathcal{T} = (\mathcal{T}, \mathcal{G})$ with the minimal homogeneity α , let $\beta, \gamma, \eta, \bar{\gamma}, \bar{\eta}$ and r be as in Theorem 2.21 and let $Z^\varepsilon = (\Pi^\varepsilon, \Gamma^\varepsilon, \Sigma^\varepsilon)$ be a discrete model which realises K^ε for \mathcal{I} . Then for any discrete modeled distribution H the following bound holds*

$$\|\mathcal{K}_\gamma^\varepsilon H\|_{\bar{\gamma}, \bar{\eta}; T}^{(\varepsilon)} \lesssim \|H\|_{\gamma, \eta; T}^{(\varepsilon)} \|\Pi^\varepsilon\|_{\gamma; T}^{(\varepsilon)} \|\Sigma^\varepsilon\|_{\gamma; T}^{(\varepsilon)} (1 + \|\Gamma^\varepsilon\|_{\bar{\gamma}; T}^{(\varepsilon)} + \|\Sigma^\varepsilon\|_{\bar{\gamma}; T}^{(\varepsilon)}), \quad (4.34)$$

and one has the identity

$$\mathcal{R}_t^\varepsilon(\mathcal{K}_\gamma^\varepsilon H)_t(x) = \int_0^t \langle \mathcal{R}_s^\varepsilon H_s, K_{t-s}^\varepsilon(x - \cdot) \rangle_\varepsilon ds. \quad (4.35)$$

Moreover, if $\bar{Z}^\varepsilon = (\bar{\Pi}^\varepsilon, \bar{\Gamma}^\varepsilon, \bar{\Sigma}^\varepsilon)$ is another discrete model realising K^ε for \mathcal{I} , and if $\bar{\mathcal{K}}_\gamma^\varepsilon$ is defined as in (4.24) for this model, then one has the bound

$$\|\mathcal{K}_\gamma^\varepsilon H; \bar{\mathcal{K}}_\gamma^\varepsilon \bar{H}\|_{\bar{\gamma}, \bar{\eta}; T}^{(\varepsilon)} \lesssim \|H; \bar{H}\|_{\gamma, \eta; T}^{(\varepsilon)} + \|Z^\varepsilon; \bar{Z}^\varepsilon\|_{\bar{\gamma}; T}^{(\varepsilon)}, \quad (4.36)$$

for all discrete modeled distributions H and \bar{H} , where the norms on H and \bar{H} are defined via the models Z^ε and \bar{Z}^ε respectively, and the proportionality constant depends on the same norms of the discrete objects as in (2.32).

Proof. The proof of the bound (4.34) for the components of $\mathcal{K}_\gamma^\varepsilon H$ not containing \mathring{K}^ε is almost identical to that of (2.30), and we only need to bound the terms $\mathring{J}^\varepsilon H$ and $\mathring{N}_\gamma^\varepsilon H$. The estimates on $\mathring{J}^\varepsilon H$ were obtained in the proof of Lemma 4.10. To bound $\mathring{N}_\gamma^\varepsilon H$, for $x, y \in \Lambda_\varepsilon^d$, we write

$$\begin{aligned} (\mathcal{R}_s^\varepsilon H_s - \Pi_x^{\varepsilon,s} \Sigma_x^{\varepsilon,st} H_t(x))(y) &= \Pi_y^{\varepsilon,s} (H_s(y) - \Gamma_{yx}^{\varepsilon,s} H_s(x))(y) \\ &\quad + \Pi_y^{\varepsilon,s} \Gamma_{yx}^{\varepsilon,s} (H_s(x) - \Sigma_x^{\varepsilon,st} H_t(x))(y), \end{aligned}$$

where we made use of Definitions 4.4 and 4.1. Estimating this expression similarly to (4.29), but using (4.3) this time, we obtain

$$\|(\mathring{N}_\gamma^\varepsilon H)_t(x)\|_0 \lesssim |t|_\varepsilon^{\eta-\gamma} \varepsilon^{\gamma+\beta} \lesssim |t|_\varepsilon^{\eta+\beta}, \quad (4.37)$$

where we have used $\gamma + \beta > 0$.

Furthermore, the operator $\Gamma_{yx}^{\varepsilon,t}$ leaves $\mathbf{1}$ invariant, and we have $\Gamma_{yx}^{\varepsilon,t}(\mathring{N}_\gamma^\varepsilon H)_t(x) = (\mathring{N}_\gamma^\varepsilon H)_t(x)$. Thus, estimating $(\mathring{N}_\gamma^\varepsilon H)_t(y)$ and $(\mathring{N}_\gamma^\varepsilon H)_t(x)$ separately by the intermediate bound in (4.37) and using $|x - y| \geq \varepsilon$, yields the required bound. In the same way we obtain the required estimate on $\Sigma_x^{\varepsilon,st}(\mathring{N}_\gamma^\varepsilon H)_t(x) - (\mathring{N}_\gamma^\varepsilon H)_s(x)$.

The bound (4.36) can be show similarly to (2.32), using the above approach. In order to show that the identity (4.35) holds, we notice that

$$(\mathcal{K}_\gamma^\varepsilon H)_t(x) \in \mathcal{T}_{\text{poly}} + \mathcal{T}_{\geq \alpha+\beta},$$

where $\mathcal{T}_{\text{poly}}$ contains only the abstract polynomials and $\alpha + \beta > 0$ by assumption. It hence follows from Definitions 4.1 and 4.4 that

$$\mathcal{R}_t^\varepsilon(\mathcal{K}_\gamma^\varepsilon H)_t(x) = \langle \mathbf{1}, (\mathcal{K}_\gamma^\varepsilon H)_t(x) \rangle,$$

which is equal to the right-hand side of (4.35). \square

5 Analysis of discrete stochastic PDEs

We consider the following spatial discretisation of equation (3.1) on $\mathbf{R}_+ \times \Lambda_\varepsilon^d$:

$$\partial_t u^\varepsilon = A^\varepsilon u^\varepsilon + F^\varepsilon(u^\varepsilon, \xi^\varepsilon), \quad u^\varepsilon(0, \cdot) = u_0^\varepsilon(\cdot), \quad (5.1)$$

where $u_0^\varepsilon \in \mathbf{R}^{\Lambda_\varepsilon^d}$, ξ^ε is a spatial discretisation of ξ , F^ε is a discrete approximation of F , and $A^\varepsilon : \ell^\infty(\Lambda_\varepsilon^d) \rightarrow \ell^\infty(\Lambda_\varepsilon^d)$ is a bounded linear operator satisfying the following assumption.

Assumption 5.1. *There exists an operator A given by a Fourier multiplier $a : \mathbf{R}^d \rightarrow \mathbf{R}$ satisfying Assumption 3.1 with an even integer parameter $\beta > 0$ and a measure μ on \mathbf{Z}^d with finite support such that*

$$(A^\varepsilon \varphi)(x) = \varepsilon^{-\beta} \int_{\mathbf{R}^d} \varphi(x - \varepsilon y) \mu(dy), \quad x \in \Lambda_\varepsilon^d, \quad (5.2)$$

for every $\varphi \in \mathcal{C}(\mathbf{R}^d)$, and such that the identity

$$\int_{\mathbf{R}^d} P(x-y) \mu(dy) = (AP)(x), \quad x \in \mathbf{R}^d, \quad (5.3)$$

holds for every polynomial P on \mathbf{R}^d with $\deg P \leq \beta$. Furthermore, the Fourier transform of μ only vanishes on \mathbf{Z}^d .

Example 5.2. A common example of the operator A is the Laplacian Δ , with its nearest neighbor discrete approximation Δ^ε , defined by (5.2) with the measure μ given by

$$\mu(\varphi) = \sum_{x \in \mathbf{Z}^d: \|x\|=1} (\varphi(x) - \varphi(0)), \quad (5.4)$$

for every $\varphi \in \ell^\infty(\mathbf{Z}^d)$, and where $\|x\|$ is the Euclidean norm. In this case, the Fourier multiplier of Δ is $a(\zeta) = -4\pi^2 \|\zeta\|^2$ and

$$(\mathcal{F}\mu)(\zeta) = -4 \sum_{i=1}^d \sin^2(\pi\zeta_i), \quad \zeta \in \mathbf{R}^d.$$

One can see that Assumption 5.1 is satisfied with $\beta = 2$.

The following section is devoted to the analysis of discrete operators.

5.1 Analysis of discrete operators

We assume that the operator $A^\varepsilon : \ell^\infty(\Lambda_\varepsilon^d) \rightarrow \ell^\infty(\Lambda_\varepsilon^d)$ satisfies Assumption 5.1 and we define the Green's function of $\partial_t - A^\varepsilon$ by

$$G_t^\varepsilon(x) \stackrel{\text{def}}{=} \varepsilon^{-d} \mathbf{1}_{t \geq 0} (e^{tA^\varepsilon} \delta_{0,\cdot})(x), \quad (t, x) \in \mathbf{R} \times \Lambda_\varepsilon^d, \quad (5.5)$$

where $\delta_{\cdot,\cdot}$ is the Kronecker's delta.

In order to build an extension of G^ε off the grid, we first choose a function $\varphi \in \mathcal{S}(\mathbf{R}^d)$ whose values coincide with $\delta_{0,\cdot}$ on \mathbf{Z}^d , and such that $(\mathcal{F}\varphi)(\zeta) = 0$ for $|\zeta|_\infty \geq 3/4$, say, where \mathcal{F} is the Fourier transform. To build such a function, write $\tilde{\varphi} \in \mathcal{C}^\infty(\mathbf{R}^d)$ for the Dirichlet kernel $\tilde{\varphi}(x) = \prod_{i=1}^d \frac{\sin(\pi x_i)}{\pi x_i}$, whose values coincide with $\delta_{0,x}$ for $x \in \mathbf{Z}^d$, and whose Fourier transform is supported in $\{\zeta : |\zeta|_\infty \leq \frac{1}{2}\}$. Choosing any function $\psi \in \mathcal{C}^\infty(\mathbf{R}^d)$ supported in the ball of radius $1/4$ around the origin and integrating to 1, it then suffices to set $\mathcal{F}\varphi = (\mathcal{F}\tilde{\varphi}) * \psi$.

Furthermore, we define the bounded operator $\tilde{A}^\varepsilon : \mathcal{C}_b(\mathbf{R}^d) \rightarrow \mathcal{C}_b(\mathbf{R}^d)$ by the right-hand side of (5.2). Then, denoting as usual by φ^ε the rescaled version of φ , we have for G^ε the representation

$$G_t^\varepsilon(x) = \mathbf{1}_{t \geq 0} (e^{t\tilde{A}^\varepsilon} \varphi^\varepsilon)(x), \quad (t, x) \in \mathbf{R} \times \Lambda_\varepsilon^d. \quad (5.6)$$

By setting $x \in \mathbf{R}^d$ in (5.6), we obtain an extension of G^ε to \mathbf{R}^{d+1} , which we again denote by G^ε .

Unfortunately, the function $G_t^\varepsilon(x)$ is discontinuous at $t = 0$, and our next aim is to modify it in such a way that it becomes differentiable at least for sufficiently large values of $|x|$. Since \tilde{A}^ε generates a strongly continuous semigroup, for every $m \in \mathbf{N}$ we have the uniform limit

$$\lim_{t \downarrow 0} \partial_t^m G_t^\varepsilon = (\tilde{A}^\varepsilon)^m \varphi^\varepsilon. \quad (5.7)$$

This gives us the terms which we have to subtract from G^ε to make it continuously differentiable at $t = 0$. For this, we take a function $\varrho : \mathbf{R} \rightarrow \mathbf{R}$ such that $\varrho(t) = 1$ for $t \in [0, \frac{1}{2}]$, $\varrho(t) = 0$ for $t \in (-\infty, 0) \cup [1, +\infty)$, and $\varrho(t)$ is smooth on $t > 0$. Then, for $r > 0$, we define

$$T^{\varepsilon, r}(t, x) \stackrel{\text{def}}{=} \varrho(t/\varepsilon^\beta) \sum_{m \leq r/\beta} \frac{t^m}{m!} (\tilde{A}^\varepsilon)^m \varphi^\varepsilon(x), \quad (t, x) \in \mathbf{R}^{d+1}. \quad (5.8)$$

The role of the function ϱ is to have $T^{\varepsilon, r}$ compactly supported in t . Then we have the following result.

Lemma 5.3. *In the described context, let Assumption 5.1 be satisfied. Then for every fixed value $r > 0$ there exists a constant $c > 0$ such that the bound*

$$|D^k(G^\varepsilon - T^{\varepsilon, r})(z)| \leq C \|z\|_{\mathfrak{s}}^{-d-|k|_{\mathfrak{s}}}, \quad (5.9)$$

holds uniformly over $z \in \mathbf{R}^{d+1}$ with $\|z\|_{\mathfrak{s}} \geq c\varepsilon$, for all $k \in \mathbf{N}^{d+1}$ with $|k|_{\mathfrak{s}} \leq r$, for D^k being a space-time derivative and for the space-time scaling $\mathfrak{s} = (\beta, 1, \dots, 1)$.

Moreover, for $|t|_\varepsilon \stackrel{\text{def}}{=} |t|^{1/\beta} \vee \varepsilon$, the function $\bar{G}_t^\varepsilon(x) \stackrel{\text{def}}{=} |t|_\varepsilon^d G_t^\varepsilon(|t|_\varepsilon x)$ is Schwartz in x , i.e. for every $m \in \mathbf{N}$ and $\bar{k} \in \mathbf{N}^d$ there is a constant \bar{C} such that the bound

$$|D_x^{\bar{k}} \bar{G}_t^\varepsilon(x)| \leq \bar{C} (1 + |x|)^{-m}, \quad (5.10)$$

holds uniformly over $(t, x) \in \mathbf{R}^{d+1}$.

Proof. The function $G^\varepsilon - T^{\varepsilon, r}$ is of class C_s^r on \mathbf{R}^{d+1} . Indeed, spatial regularity follows immediately from the regularity of φ and commutation of \tilde{A}^ε with the differential operator. Continuous differentiability at $t = 0$ follows from (5.7). Furthermore, since G^ε vanishes on $t \leq 0$, we only need to consider $t > 0$.

Next, we notice that the bound (5.9) follows from (5.10). Let $\hat{r} > 0$ be such that the measure μ in Assumption 5.1 is supported in the ball of radius \hat{r} . Then, for $k = (k_0, \bar{k}) \in \mathbf{N}^{d+1}$ with $k_0 \in \mathbf{N}$ and $|k|_{\mathfrak{s}} \leq r$ we use (5.6) and the identities (5.3), combined with the Taylor's formula, to get

$$|D^k G_t^\varepsilon(x)| = |(\tilde{A}^\varepsilon)^{k_0} D_x^{\bar{k}} G_t^\varepsilon(x)| \lesssim \sup_{y: |y-x| \leq k_0 \hat{r} \varepsilon} \sup_{l: |l| = \beta k_0} |D_y^{\bar{k}+l} G_t^\varepsilon(y)|, \quad (5.11)$$

where $y \in \mathbf{R}^d$, $l \in \mathbf{N}^d$. For $\|t, x\|_{\mathfrak{s}} \geq c\varepsilon$, in the case $|t|^{1/\beta} \geq |x|$, we bound the right-hand side of (5.11) using (5.10) with $m = 0$, what gives an estimate of order

$|t|^{-(d+|k|_s)/\beta}$. In the case $|t|^{1/\beta} < |x|$, we use (5.10) with $m = d + |k|_s$, and we get a bound of order $|x|^{-d-|k|_s}$, if we take $c \geq 2r\hat{r}/\beta$. Furthermore, the required bound on $T^{\varepsilon,r}$ follows easily from the properties of the functions φ and ϱ . Hence, we only need to prove the bound (5.10).

Denoting by \mathcal{F} the Fourier transform, we get from (5.6) and Assumption 5.1:

$$(\mathcal{F}\bar{G}_t^\varepsilon)(\zeta) = (\mathcal{F}\varphi)(\varepsilon|t|_\varepsilon^{-1}\zeta) e^{t|t|_\varepsilon^{-1}a(\zeta)f(\varepsilon|t|_\varepsilon^{-1}\zeta)}, \quad (5.12)$$

where we have used the scaling property $\lambda^\beta a(\zeta) = a(\lambda\zeta)$, and where $f \stackrel{\text{def}}{=} (\mathcal{F}\mu)/a$.

We start with considering the case $t \geq \varepsilon^\beta$. It follows from the last part of Assumption 5.1 that there exists $\bar{c} > 0$ such that $f(\zeta) \geq \bar{c}$ for $|\zeta|_\infty \leq 3/4$. Since $\varepsilon|t|_\varepsilon^{-1} \leq 1$, we conclude that

$$|D_\zeta^{\bar{k}} e^{a(\zeta)f(\varepsilon|t|_\varepsilon^{-1}\zeta)}| \lesssim |\zeta|^{|\bar{k}|} e^{a(\zeta)\bar{c}} \lesssim (1 + |\zeta|)^{-m},$$

for $|\zeta|_\infty < 3/(4\varepsilon|t|_\varepsilon^{-1})$, for every $m \geq 0$ and for a proportionality constant dependent on m and \bar{k} . Here, we have used $a(\zeta) < 0$ and polynomial growth of $|a(\zeta)|$. Since $(\mathcal{F}\varphi)(\varepsilon|t|_\varepsilon^{-1}\zeta)$ vanishes for $|\zeta|_\infty \geq 3/(4\varepsilon|t|_\varepsilon^{-1})$, we conclude that

$$|D_\zeta^{\bar{k}}(\mathcal{F}\bar{G}_t^\varepsilon)(\zeta)| \lesssim (1 + |\zeta|)^{-m},$$

uniformly in t and ε (provided that $t \geq \varepsilon^\beta$), and for every $m \in \mathbf{N}$ and $\bar{k} \in \mathbf{N}^d$.

In the case $t < \varepsilon^\beta$, we can bound the exponent in (5.12) by 1, and the polynomial decay comes from the factor $(\mathcal{F}\varphi)(\zeta)$, because $\varphi \in \mathcal{S}(\mathbf{R}^d)$. Since the Fourier transform is continuous on Schwartz space, this implies that \bar{G}_t^ε is a Schwartz function, with bounds uniform in ε and t , which is exactly the claim. \square

The following result is an analogue of Lemma 3.3 for G^ε .

Lemma 5.4. *Let Assumption 5.1 be satisfied. Then, the function G^ε defined in (5.6) can be written as $G^\varepsilon = K^\varepsilon + R^\varepsilon$ in such a way that the identity*

$$(G^\varepsilon \star_\varepsilon u)(z) = (K^\varepsilon \star_\varepsilon u)(z) + (R^\varepsilon \star_\varepsilon u)(z), \quad (5.13)$$

holds for all $z \in (-\infty, 1] \times \Lambda_\varepsilon^d$ and all functions u on $\mathbf{R}_+ \times \Lambda_\varepsilon^d$, periodic in the spatial variable with some fixed period. Furthermore, K^ε is regularising of order β in the sense of Definition 4.7, for arbitrary (but fixed) r and with the scaling $\mathfrak{s} = (\beta, 1, \dots, 1)$. The function R^ε is compactly supported, non-anticipative and the norm $\|R^\varepsilon\|_{C^r}$ is bounded uniformly in ε .

Proof. Let $M : \mathbf{R}^{d+1} \rightarrow \mathbf{R}_+$ be a smooth norm for the scaling \mathfrak{s} (see for example [Hai14, Rem. 2.13]). Furthermore, let $\bar{\varrho} : \mathbf{R}_+ \rightarrow [0, 1]$ be a smooth ‘‘cutoff function’’ such that $\bar{\varrho}(s) = 0$ if $s \notin [1/2, 2]$, and such that $\sum_{n \in \mathbf{Z}} \bar{\varrho}(2^n s) = 1$ for all $s > 0$ (see the construction of the partition of unity in [BCD11]). For integers $n \in [0, N)$ we set the functions $\bar{\varrho}_n(z) \stackrel{\text{def}}{=} \bar{\varrho}(2^n M(z))$, $\bar{\varrho}_{<0} \stackrel{\text{def}}{=} \sum_{n < 0} \bar{\varrho}_n$, $\bar{\varrho}_{\geq N} \stackrel{\text{def}}{=} \sum_{n \geq N} \bar{\varrho}_n$, and

$$\bar{K}^{\varepsilon,n}(z) = \bar{\varrho}_n(z)(G^\varepsilon - T^{\varepsilon,r})(z), \quad \bar{R}^\varepsilon(z) = \bar{\varrho}_{<0}(z)(G^\varepsilon - T^{\varepsilon,r})(z),$$

$$\tilde{K}^\varepsilon(z) = \bar{\varrho}_{\geq N}(z)(G^\varepsilon - T^{\varepsilon,r})(z) + T^{\varepsilon,r}(z). \quad (5.14)$$

Then it follows immediately from the properties of $\bar{\varrho}$ that

$$G^\varepsilon = \sum_{n=0}^{N-1} \bar{K}^{(\varepsilon,n)} + \tilde{K}^\varepsilon + \bar{R}^\varepsilon.$$

Since $\bar{\varrho}_{<0}$ is supported away from the origin, we use (5.9) and Assumption 5.1 to conclude that $\|\bar{R}^\varepsilon\|_{C^r}$ is bounded uniformly in ε . (Actually, its value and derivatives even decay faster than any power.)

Furthermore, the function $\bar{K}^{(\varepsilon,n)}$ is supported in the ball of radius $c2^{-n}$, for c as in Lemma 5.3, provided that the norm M was chosen such that $M(z) \geq 2c\|z\|_s$. By the same reason, the first term in (5.14) is supported in the ball of radius $c\varepsilon$. Moreover, the support property of the measure μ and the properties of the functions ϱ and φ^ε in (5.8) yield that the restriction of $T^{\varepsilon,r}$ to the grid Λ_ε^d in space is supported in the ball of radius $c\varepsilon$, as soon as $c \geq 2r\hat{r}/\beta$, where \hat{r} is from Assumption 5.1.

As a consequence of (5.2), (5.6) and (5.8), we get for $0 \leq n < N$ the scaling properties

$$\bar{K}^{(\varepsilon,n)}(z) = 2^{nd} \bar{K}^{(\varepsilon 2^n, 0)}(2^{sn} z), \quad \tilde{K}^\varepsilon(z) = \varepsilon^{-d} \tilde{K}^1(\varepsilon^{-sn} z),$$

and (2.24) and (4.22) follow immediately from (5.9) and (5.8).

It remains to modify these functions in such a way that they “kill” polynomials in the sense of (4.21). To this end, we take a smooth function $P^{(N)}$ on \mathbf{R}^{d+1} , whose support coincides with the support of \tilde{K}^ε , which satisfies $|P^{(N)}(z)| \lesssim \varepsilon^{-d}$, for every $z \in \mathbf{R}^{d+1}$, and such that one has

$$\int_{\mathbf{R} \times \Lambda_\varepsilon^d} (\tilde{K}^\varepsilon - P^{(N)})(z) dz = 0. \quad (5.15)$$

Then we define \hat{K}^ε to be the restriction of $\tilde{K}^\varepsilon - P^{(N)}$ to the grid Λ_ε^d in space. Apparently, the function \hat{K}^ε has the same scaling and support properties as \tilde{K}^ε , and it follows from (5.15) that it satisfies (4.21) with $k = 0$.

Moreover, we can recursively build a sequence of smooth functions $P^{(n)}$, for integers $n \in [0, N)$, such that $P^{(n)}$ is supported in the ball of radius 2^{-n} , the function $P^{(n)}$ satisfies the bounds in (2.24), and for every $k \in \mathbf{N}^{d+1}$ with $|k|_s \leq r$ one has

$$\int_{\mathbf{R} \times \Lambda_\varepsilon^d} z^k (\bar{K}^{(\varepsilon,n)} - P^{(n)} + P^{(n+1)})(z) dz = 0. \quad (5.16)$$

Then, for such values of n , we define

$$K^{(\varepsilon,n)} = \bar{K}^{(\varepsilon,n)} - P^{(n)} + P^{(n+1)}, \quad R^\varepsilon \stackrel{\text{def}}{=} \bar{R}^\varepsilon + P^{(0)}.$$

It follows from the properties of the functions $P^{(n)}$ that $K^{(\varepsilon,n)}$ has all the required properties. The function R^ε also has the required properties, and the decompositions (4.20) and (5.13) hold by construction. Finally, using (5.10), we can make the function R^ε compactly supported in the same way as in [Hai14, Lem. 7.7]. \square

By analogy with (3.2), we use the function R^ε from Lemma 5.4 to define for periodic $\zeta_t \in \mathbf{R}^{\Lambda_\varepsilon^d}$, $t \in \mathbf{R}$, the abstract polynomial

$$(R_\gamma^\varepsilon \zeta)_t(x) \stackrel{\text{def}}{=} \sum_{|k|_s < \gamma} \frac{X^k}{k!} \int_{\mathbf{R}} \langle \zeta_s, D^k R_{t-s}^\varepsilon(x - \cdot) \rangle_\varepsilon ds, \quad (5.17)$$

where as before $k \in \mathbf{N}^{d+1}$ and the mixed derivative D^k is in space-time.

5.2 Properties of the discrete equation

In this section we show that a discrete analogue of Theorem 3.9 holds for the solution map of the equation (5.1) with an operator A^ε satisfying Assumption 5.1.

Similarly to [Hai14, Lem. 7.5], but using the properties of G^ε proved in the previous section, we can show that for every periodic $u_0^\varepsilon \in \mathbf{R}^{\Lambda_\varepsilon^d}$, we have a discrete analogue of Lemma 3.6 for the map $(t, x) \mapsto S_t^\varepsilon u_0^\varepsilon(x)$, where S^ε is the semigroup generated by A^ε .

For the regularity structure \mathcal{F} from Section 3.1, we take a truncated regularity structure $\hat{\mathcal{F}} = (\hat{\mathcal{T}}, \mathcal{G})$ and make the following assumption on the nonlinearity F^ε .

Assumption 5.5. *For some $0 < \bar{\gamma} \leq \gamma$, $\eta \in \mathbf{R}$, every $\varepsilon > 0$ and every discrete model Z^ε on $\hat{\mathcal{F}}$, there exist discrete modeled distributions $F_0^\varepsilon(Z^\varepsilon)$ and $I_0^\varepsilon(Z^\varepsilon)$, with exactly the same properties as of F_0 and I_0 in Assumption 3.8 on the grid. Furthermore, we define \hat{F}^ε as in (3.9), but via F^ε and F_0^ε , and we define $\hat{F}^\varepsilon(H)$ for $H : \mathbf{R}_+ \times \Lambda_\varepsilon^d \rightarrow \mathcal{T}_{<\gamma}$ as in (3.10). Finally, we assume that the discrete analogue of the Lipschitz condition (3.12) holds for \hat{F}^ε , with the constant C independent of ε .*

Similarly to (3.11), but using the discrete operators (4.4), (5.17) and (4.24), we reformulate the equation (5.1) as

$$U^\varepsilon = \mathcal{P}^\varepsilon \hat{F}^\varepsilon(U^\varepsilon) + S^\varepsilon u_0^\varepsilon + I_0^\varepsilon, \quad (5.18)$$

where $\mathcal{P}^\varepsilon \stackrel{\text{def}}{=} \mathcal{K}_{\bar{\gamma}}^\varepsilon + R_\gamma^\varepsilon \mathcal{R}^\varepsilon$ and U^ε is a discrete modeled distribution.

Remark 5.6. If Z^ε is a canonical discrete model, then it follows from (4.35), (5.17), (4.4), Definition 4.1 and Assumption 5.5 that

$$u_t^\varepsilon(x) = (\mathcal{R}_t^\varepsilon U_t^\varepsilon)(x), \quad (t, x) \in \mathbf{R}_+ \times \Lambda_\varepsilon^d. \quad (5.19)$$

is a solution of the equation (5.1).

The following result can be proven in the same way as Theorem 3.9.

Theorem 5.7. *Let Z^ε be a sequence of models and let u_0^ε be a sequence of periodic functions on Λ_ε^d . Let furthermore the assumptions of Theorem 3.9 and Assumption 5.5 be satisfied. Then there exists $T_\star \in (0, +\infty]$ such that for every $T < T_\star$ the sequence of solution maps $\mathcal{S}_T^\varepsilon : (u_0^\varepsilon, Z^\varepsilon) \mapsto U^\varepsilon$ of the equation (5.18) is jointly Lipschitz continuous (uniformly in ε !) in the sense of Theorem 3.9, but for the discrete objects.*

Remark 5.8. Since we require uniformity in ε in Theorem 5.7, the solution of equation (5.18) is considered only up to some time point T_\star .

6 Inhomogeneous Gaussian models

In this section we analyse the discrete models which are built from a Gaussian noise. We will work as usual on the grid Λ_ε^d , with $\varepsilon = 2^{-N}$ and $N \in \mathbf{N}$, and with the time-space scaling $\mathfrak{s} = (\mathfrak{s}_0, 1, \dots, 1)$.

We assume that we are given a probability space $(\Omega, \mathcal{F}, \mathbf{P})$, together with a white noise ξ over the Hilbert space $H \stackrel{\text{def}}{=} L^2(D)$ (see [Nua06]), where $D \stackrel{\text{def}}{=} \mathbf{R} \times \mathbf{T}^d$ and $\mathbf{T} \stackrel{\text{def}}{=} \mathbf{R}/\mathbf{Z}$ is the unit circle. In the sequel, we will always identify ξ with its periodic extension to \mathbf{R}^{d+1} .

In order to build a spatial discretisation of ξ , we take a compactly supported function $\varrho : \mathbf{R}^d \rightarrow \mathbf{R}$, such that for every $y \in \mathbf{Z}^d$ one has

$$\int_{\mathbf{R}^d} \varrho(x) \varrho(x - y) dx = \delta_{0,y} ,$$

where $\delta_{\cdot, \cdot}$ is the Kronecker's function. Then, for $x \in \Lambda_\varepsilon^d$, we define the scaled function $\varrho_x^\varepsilon(y) \stackrel{\text{def}}{=} \varepsilon^{-d} \varrho((y - x)/\varepsilon)$ and

$$\xi^\varepsilon(t, x) \stackrel{\text{def}}{=} \xi(t, \varrho_x^\varepsilon) , \quad (t, x) \in \mathbf{R} \times \Lambda_\varepsilon^d . \quad (6.1)$$

One can see that ξ^ε is a white noise on the Hilbert space $H_\varepsilon \stackrel{\text{def}}{=} L^2(\mathbf{R}) \otimes \ell^2(\mathbf{T}_\varepsilon^d)$, where $\mathbf{T}_\varepsilon \stackrel{\text{def}}{=} (\varepsilon\mathbf{Z})/\mathbf{Z}$ and $\ell^2(\mathbf{T}_\varepsilon^d)$ is equipped with the inner product $\langle \cdot, \cdot \rangle_\varepsilon$.

In the setting of Section 3.2, we assume that $Z^\varepsilon = (\Pi^\varepsilon, \Gamma^\varepsilon, \Sigma^\varepsilon)$ is a discrete model on \mathcal{T}^{gen} such that, for each $\tau \in \mathcal{F}^{\text{gen}}$, the maps $\langle \Pi_x^{\varepsilon,t} \tau, \varphi \rangle_\varepsilon$, $\Gamma_{xy}^{\varepsilon,t} \tau$ and $\Sigma_x^{\varepsilon,st} \tau$ belong to the inhomogeneous Wiener chaos of order $\|\tau\|$ (the number of occurrences of Ξ in τ) with respect to ξ^ε . Moreover, we assume that the distributions of the functions $(t, x) \mapsto \langle \Pi_x^{\varepsilon,t} \tau, \varphi \rangle_\varepsilon$, $(t, x) \mapsto \Gamma_{x, x+h}^{\varepsilon,t} \tau$ and $(t, x) \mapsto \Sigma_x^{\varepsilon,t, t+h} \tau$ are stationary. It follows then from Remark 4.11 and Wick's lemma [Nua06, Prop. 1.1.3], that the canonical extension of Z^ε to $\hat{\mathcal{T}}$, in the sense of Remark 4.11, has the same property. In what follows, we will call the discrete models with these properties *stationary Gaussian discrete models*.

The following result provides a criterion for such a model to be bounded uniformly in ε . In its statement we use the following set:

$$\hat{\mathcal{F}}^- \stackrel{\text{def}}{=} \left(\{ \tau \in \hat{\mathcal{F}} : |\tau| < 0 \} \cup \mathcal{F}^{\text{gen}} \right) \setminus \mathcal{F}_{\text{poly}} . \quad (6.2)$$

Theorem 6.1. *Let $\hat{\mathcal{T}} = (\hat{\mathcal{T}}, \mathcal{G})$ be a truncated regularity structure and let Z^ε be an admissible stationary Gaussian discrete model on it. Let furthermore $\kappa > 0$ be such that the bounds*

$$\begin{aligned} \mathbf{E} \|\Gamma_{xy}^{\varepsilon,t} \tau\|_m^2 &\lesssim |x - y|^{2(|\tau| - m) + \kappa} , \\ \mathbf{E} \|\Sigma_x^{\varepsilon,st} \tau\|_m^2 &\lesssim (|t - s|^{1/\mathfrak{s}_0} \vee \varepsilon)^{2(|\tau| - m) + \kappa} , \end{aligned} \quad (6.3)$$

hold for $\tau \in \mathcal{F}^{\text{gen}} \setminus \mathcal{F}_{\text{poly}}$, for all $s, t \in [-T, T]$, all $x, y \in \Lambda_\varepsilon^d$, all $m \in \hat{\mathcal{A}}$ such that $m < |\tau|$, and for some $T \geq c$, where $c > 0$ is from Definition 4.7. Let finally

for some $\delta > 0$ and for each $\tau \in \hat{\mathcal{F}}^-$ the bounds

$$\begin{aligned} \mathbf{E} |\langle \Pi_x^{\varepsilon,t}, \varphi_x^\lambda \rangle_\varepsilon|^2 &\lesssim \lambda^{2|\tau|+\kappa}, \\ \mathbf{E} |\langle (\Pi_x^{\varepsilon,t} - \Pi_x^{\varepsilon,s}), \varphi_x^\lambda \rangle_\varepsilon|^2 &\lesssim \lambda^{2(|\tau|-\delta)+\kappa} (|t-s|^{1/50} \vee \varepsilon)^{2\delta}, \end{aligned} \quad (6.4)$$

hold uniformly in all $\lambda \in [\varepsilon, 1]$, all $s, t \in [-T, T]$, all $x \in \Lambda_\varepsilon^d$ and all $\varphi \in \mathcal{B}_0^r(\mathbf{R}^d)$ with $r > -[\min \hat{\mathcal{A}}]$. Then, for every $\gamma > 0$, $p \geq 1$ and $\bar{\delta} \in [0, \delta)$, one has the following bound:

$$\mathbf{E} \left[\left\| Z^\varepsilon \right\|_{\bar{\delta}, \gamma; T}^{(\varepsilon)} \right]^p \lesssim 1. \quad (6.5)$$

Finally, let \bar{Z}^ε be another admissible stationary Gaussian discrete model on $\hat{\mathcal{T}}$, such that for some $\theta > 0$ and some $\bar{\varepsilon} > 0$ the maps $\Gamma^\varepsilon - \bar{\Gamma}^\varepsilon$, $\Sigma^\varepsilon - \bar{\Sigma}^\varepsilon$ and $\Pi^\varepsilon - \bar{\Pi}^\varepsilon$ satisfy the bounds (6.3) and (6.4) respectively with proportionality constants of order $\bar{\varepsilon}^{2\theta}$. Then, for every $\gamma > 0$, $p \geq 1$ and $\bar{\delta} \in [0, \delta)$, one has the bound

$$\mathbf{E} \left[\left\| Z^\varepsilon; \bar{Z}^\varepsilon \right\|_{\bar{\delta}, \gamma; T}^{(\varepsilon)} \right]^p \lesssim \bar{\varepsilon}^{\theta p}. \quad (6.6)$$

Proof. The bound (6.5) on the regularity structure \mathcal{T}^{gen} follows from (6.3), (6.4) and equivalence of moments for elements from a fixed Wiener chaos. Going by induction from the elements of \mathcal{T}^{gen} to the elements of $\hat{\mathcal{T}}$ and using Lemmas 4.9 and 4.10, we can obtain (6.5) in the same way as in the proof of [Hai14, Thm. 10.7]. Similarly we can prove that the bound (6.6) holds. \square

The conditions (6.4) can be checked quite easily if the maps $\Pi^\varepsilon \tau$ have certain Wiener chaos expansions. More precisely, we assume that there exist kernels $\mathcal{W}^{(\varepsilon;k)}_\tau$ such that $(\mathcal{W}^{(\varepsilon;k)}_\tau)(z) \in H_\varepsilon^{\otimes k}$, for $z \in \mathbf{R} \times \Lambda_\varepsilon^d$, and

$$\langle \Pi_0^{\varepsilon,t} \tau, \varphi \rangle_\varepsilon = \sum_{k \leq \|\tau\|} I_k^\varepsilon \left(\int_{\Lambda_\varepsilon^d} \varphi(y) (\mathcal{W}^{(\varepsilon;k)}_\tau)(t, y) dy \right), \quad (6.7)$$

where I_k^ε is the k -th order Wiener integrals with respect to ξ^ε and the space H_ε is introduced above. Then we define

$$(\mathcal{K}^{(\varepsilon;k)}_\tau)(z_1, z_2) \stackrel{\text{def}}{=} \langle (\mathcal{W}^{(\varepsilon;k)}_\tau)(z_1), (\mathcal{W}^{(\varepsilon;k)}_\tau)(z_2) \rangle_{H_\varepsilon^{\otimes k}}, \quad (6.8)$$

for $z_1 \neq z_2 \in \mathbf{R} \times \Lambda_\varepsilon^d$, assuming that the expressions on the right-hand side are well-defined.

In the same way, we assume that the maps $\bar{\Pi}^\varepsilon \tau$ are given by (6.7) via the respective kernels $\bar{\mathcal{W}}^{(\varepsilon;k)}_\tau$. Moreover, we define the functions $\delta \mathcal{K}^{(\varepsilon;k)}_\tau$ as in (6.8), but via the kernels $\mathcal{W}^{(\varepsilon;k)}_\tau - \bar{\mathcal{W}}^{(\varepsilon;k)}_\tau$, and we assume that the functions $\mathcal{K}^{(\varepsilon;k)}_\tau$ and $\delta \mathcal{K}^{(\varepsilon;k)}_\tau$ depend on the time variables t_1 and t_2 only via $t_1 - t_2$, i.e.

$$(\mathcal{K}^{(k)}_\tau)_{t_1-t_2}(x_1, x_2) \stackrel{\text{def}}{=} (\mathcal{K}^{(k)}_\tau)(z_1, z_2), \quad (6.9)$$

where $z_i = (t_i, x_i)$, and similarly for $\delta \mathcal{K}^{(\varepsilon;k)}_\tau$.

The following result shows that the bounds (6.4) follows from corresponding bounds on these functions.

Proposition 6.2. *In the described context, we assume that for some $\tau \in \hat{\mathcal{F}}^-$ there are values $\alpha > |\tau| \vee (-d/2)$ and $\delta \in (0, \alpha + d/2)$ such that the bounds*

$$\begin{aligned} |(\mathcal{K}^{(\varepsilon;k)}\tau)_0(x_1, x_2)| &\lesssim \sum_{\zeta \geq 0} (|x_1| \vee |x_2|)^\zeta (|x_1 - x_2| \vee \varepsilon)^{2\alpha - \zeta}, \\ \frac{|\delta^{0,t}(\mathcal{K}^{(\varepsilon;k)}\tau)(x_1, x_2)|}{(|t|^{1/s_0} \vee \varepsilon)^{2\delta}} &\lesssim \sum_{\zeta \geq 0} (|x_1| \vee |x_2| \vee |t|^{1/s_0})^\zeta (|x_1 - x_2| \vee \varepsilon)^{2\alpha - 2\delta - \zeta}, \end{aligned} \quad (6.10)$$

hold for $x_1, x_2 \in \Lambda_\varepsilon^d$ and $k \leq \|\tau\|$, where the operator $\delta^{0,t}$ is defined in (1.4), and where the sums run over finitely many values of $\zeta \in [0, 2\alpha - 2\delta + d]$. Then the bounds (6.4) hold for τ with a sufficiently small κ .

Let furthermore (6.10) hold for the function $\delta\mathcal{K}^{(\varepsilon;k)}\tau$ with the proportionality constant of order $\varepsilon^{2\theta}$, for some $\theta > 0$. Then the required bounds on $(\Pi^\varepsilon - \bar{\Pi}^\varepsilon)\tau$ in Theorem 6.1 hold.

Proof. We note that due to our assumptions on stationarity of the models, it is sufficient to check the conditions (6.4) only for $\langle \Pi_0^{\varepsilon,0}\tau, \varphi_0^\lambda \rangle_\varepsilon$ and $\langle (\Pi_0^{\varepsilon,t} - \Pi_0^{\varepsilon,0})\tau, \varphi_0^\lambda \rangle_\varepsilon$, and respectively for the other model.

We start with the proof of the first statement of this proposition. We denote by $\Pi_0^{(\varepsilon,k),t}\tau$ the component of $\Pi_0^{\varepsilon,t}\tau$ belonging to the k -th homogeneous Wiener chaos. Furthermore, we will use the following property of the Wiener integral [Nua06]:

$$\mathbf{E}[I_k^\varepsilon(f)^2] \leq \|f\|_{H_\varepsilon^{\otimes k}}, \quad f \in H_\varepsilon^{\otimes k}. \quad (6.11)$$

Thus, from this property, (6.9) and the first bound in (6.10), we get

$$\begin{aligned} \mathbf{E}|\langle \Pi_0^{(\varepsilon,k),0}\tau, \varphi_0^\lambda \rangle_\varepsilon|^2 &\lesssim \int_{\Lambda_\varepsilon^d} \int_{\Lambda_\varepsilon^d} |\varphi_0^\lambda(x_1)| |\varphi_0^\lambda(x_2)| |(\mathcal{K}^{(\varepsilon,k)}\tau)_0(x_1, x_2)| dx_1 dx_2 \\ &\lesssim \lambda^{-2d} \sum_{\zeta \geq 0} \int_{\substack{|x_1| \leq \lambda \\ |x_2| \leq \lambda}} (|x_1| \vee |x_2|)^\zeta (|x_1 - x_2| \vee \varepsilon)^{2\alpha - \zeta} dx_1 dx_2 \\ &\lesssim \lambda^{-2d} \sum_{\zeta \geq 0} \lambda^{d+\zeta} \int_{|x| \leq 2\lambda} (|x| \vee \varepsilon)^{2\alpha - \zeta} dx \lesssim \lambda^{2\alpha}, \end{aligned} \quad (6.12)$$

for $\lambda \geq \varepsilon$. Here, to have the proportionality constant independent of ε , we need $2\alpha - \zeta > -d$. The first estimate in (6.4) obviously follows from these bounds for all k .

Now, we will investigate the time regularity of the discrete model. For $|t| \geq \lambda^{s_0}$ we can use (6.12) and brutally bound

$$\begin{aligned} \mathbf{E}|\langle (\Pi_0^{(\varepsilon,k),t} - \Pi_0^{(\varepsilon,k),0})\tau, \varphi_0^\lambda \rangle_\varepsilon|^2 &\lesssim \mathbf{E}|\langle \Pi_0^{(\varepsilon,k),t}\tau, \varphi_0^\lambda \rangle_\varepsilon|^2 + \mathbf{E}|\langle \Pi_0^{(\varepsilon,k),0}\tau, \varphi_0^\lambda \rangle_\varepsilon|^2 \\ &\lesssim \lambda^{2\alpha} \lesssim |t|^{2\delta/s_0} \lambda^{2\alpha - 2\delta}, \end{aligned}$$

for any $\delta \geq 0$, which is the required estimate. In the case $|t| < \lambda^{s_0}$, the second bound in (6.10) and (6.11) yield

$$\mathbf{E}|\langle \delta^{0,t}\Pi_0^{(\varepsilon,k)}\tau, \varphi_0^\lambda \rangle_\varepsilon|^2 \lesssim \int_{\Lambda_\varepsilon^d} \int_{\Lambda_\varepsilon^d} |\varphi_0^\lambda(x_1)| |\varphi_0^\lambda(x_2)| |\delta^{0,t}(\mathcal{K}^{(\varepsilon,k)}\tau)(x_1, x_2)| dx_1 dx_2$$

$$\begin{aligned}
& + \int_{\Lambda_\varepsilon^d} \int_{\Lambda_\varepsilon^d} |\varphi_0^\lambda(x_1)| |\varphi_0^\lambda(x_2)| |\delta^{-t,0}(\mathcal{K}^{(\varepsilon,k)}\tau)(x_1, x_2)| dx_1 dx_2 \\
& \lesssim (|t|^{1/s_0} \vee \varepsilon)^{2\delta} \lambda^{-2d} \\
& \quad \times \sum_{\zeta \geq 0} \int_{\substack{|x_1| \leq \lambda \\ |x_2| \leq \lambda}} (|x_1| \vee |x_2| \vee \lambda)^\zeta (|x_1 - x_2| \vee \varepsilon)^{2\alpha - 2\delta - \zeta} dx_1 dx_2 \\
& \lesssim (|t|^{1/s_0} \vee \varepsilon)^{2\delta} \lambda^{2\alpha - 2\delta},
\end{aligned}$$

where the integral is estimated as before. This completes the proof of the second bound in (6.4). The bounds on $(\Pi^\varepsilon - \bar{\Pi}^\varepsilon)\tau$ can be proved in the same way. \square

7 Convergence of the discrete dynamical Φ_3^4 model

In this section we use the theory developed above to prove convergence of the solutions of $(\Phi_{3,\varepsilon}^4)$, where Δ^ε is the nearest-neighbour approximation of Δ and the discrete noise ξ^ε is defined in (1.1) via a space-time white noise ξ .

Example 5.2 yields that Assumption 5.1 is satisfied, and moreover ξ^ε is a discrete noise as in (6.1). The time-space scaling for the equation (Φ_3^4) is $\mathfrak{s} = (2, 1, 1, 1)$ and the kernels K and K^ε are defined in Lemma 5.4 with the parameters $\beta = 2$ and $r > 2$, for the operators Δ and Δ^ε respectively.

The regularity structure $\mathcal{F} = (\mathcal{T}, \mathcal{G})$ for the equation (Φ_3^4) , introduced in Section 3.1, has the model space $\mathcal{T} = \text{span}\{\mathcal{F}\}$, where

$$\begin{aligned}
\mathcal{F} = \{ & \mathbf{1}, \Xi, \Psi, \Psi^2, \Psi^3, \Psi^2 X_i, \mathcal{I}(\Psi^3)\Psi, \mathcal{I}(\Psi^3)\Psi^2, \mathcal{I}(\Psi^2)\Psi^2, \mathcal{I}(\Psi^2), \\
& \mathcal{I}(\Psi)\Psi, \mathcal{I}(\Psi)\Psi^2, X_i \},
\end{aligned}$$

$\Psi \stackrel{\text{def}}{=} \mathcal{I}(\Xi)$, $|\Xi| = \alpha \in (-\frac{18}{7}, -\frac{5}{2})$ and the index i corresponds to any of the three spatial dimensions, see [Hai14, Sec. 9.2]. The bound $\alpha > -\frac{18}{7}$ is required, in order to have a fixed regularity structure. If $\alpha \leq -\frac{18}{7}$, then we have to add the symbol $\Psi^2 \mathcal{I}(\Psi^2 \mathcal{I}(\Psi^3))$ into the set \mathcal{F} .

A two-parameter renormalisation subgroup $\mathfrak{R}^0 \subset \mathfrak{R}$ for this problem consists of the linear maps M on \mathcal{T} , defined by

$$\begin{aligned}
M\Psi^2 &= \Psi^2 - C_1 \mathbf{1}, \\
M(\Psi^2 X_i) &= \Psi^2 X_i - 3C_1 X_i, \\
M\Psi^3 &= \Psi^3 - C_1 \Psi, \\
M(\mathcal{I}(\Psi^2)\Psi^2) &= \mathcal{I}(\Psi^2)(\Psi^2 - C_1 \mathbf{1}) - C_2 \mathbf{1}, \\
M(\mathcal{I}(\Psi^3)\Psi) &= (\mathcal{I}(\Psi^3) - 3C_1 \mathcal{I}(\Psi))\Psi, \\
M(\mathcal{I}(\Psi^3)\Psi^2) &= (\mathcal{I}(\Psi^3) - 3C_1 \mathcal{I}(\Psi))(\Psi^2 - C_1 \mathbf{1}) - 3C_2 \Psi, \\
M(\mathcal{I}(\Psi)\Psi^2) &= \mathcal{I}(\Psi)(\Psi^2 - C_1 \mathbf{1}),
\end{aligned} \tag{7.1}$$

as well as $M\tau = \tau$ for the remaining elements $\tau \in \mathcal{F}$, and where C_1 and C_2 are the two parameter constants (see [Hai14, Sec. 9.2] for the proof that $M \in \mathfrak{R}$).

As one can see from Proposition 6.2, one can expect a concrete realisation of an abstract symbol τ to be a function in time, only if $|\tau| > -\frac{3}{2}$. In our case, the symbols Ξ and Ψ^3 don't satisfy this condition, having homogeneities $\alpha < -\frac{5}{2}$ and $3(\alpha + \beta) < -\frac{3}{2}$ respectively. To resolve this problem, we define $\bar{\Psi} \stackrel{\text{def}}{=} \mathcal{I}(\Psi^3)$ and the sets $\mathcal{F}^{\text{gen}} \stackrel{\text{def}}{=} \{\Psi, \bar{\Psi}\} \cup \mathcal{F}_{\text{poly}}$ and

$$\hat{\mathcal{F}} = \{\mathbf{1}, \Psi, \Psi^2, \Psi^2 X_i, \Psi \bar{\Psi}, \Psi^2 \bar{\Psi}, \mathcal{I}(\Psi^2) \Psi^2, \mathcal{I}(\Psi^2), \mathcal{I}(\Psi) \Psi, \mathcal{I}(\Psi) \Psi^2, X_i\}.$$

Then the model spaces of the regularity structures \mathcal{T}^{gen} and $\hat{\mathcal{T}}$ from Definition 3.4 are the linear spans of \mathcal{F}^{gen} and $\hat{\mathcal{F}}$ respectively. The set $\hat{\mathcal{F}}^-$ from (6.2) is given by

$$\hat{\mathcal{F}}^- = \{\Psi, \bar{\Psi}, \Psi^2, \Psi^2 X_i, \Psi \bar{\Psi}, \mathcal{I}(\Psi^2) \Psi^2, \Psi^2 \bar{\Psi}\}.$$

In the following lemma we show that our nonlinearities satisfy the required assumptions, provided that the appearance of the renormalisation constant is being dealt with at the level of the corresponding models.

Lemma 7.1. *Let $\hat{\alpha} \stackrel{\text{def}}{=} \min \hat{\mathcal{A}}$. Then, for any $\gamma > |2\hat{\alpha}|$ and any $\eta \leq \hat{\alpha}$, the maps $F(u, \xi) = F^\varepsilon(u, \xi) = \xi - u^3$ satisfy Assumptions 3.8 and 5.5 with*

$$F_0(\Xi) = F_0^\varepsilon(\Xi) = \Xi + \Psi^3, \quad I_0(\Xi) = I_0^\varepsilon(\Xi) = \Psi + \bar{\Psi},$$

and $\bar{\gamma} = \gamma + 2\hat{\alpha}$, $\bar{\eta} = 3\eta$.

Proof. Since the set $\mathcal{U} \subset \mathcal{F}$, introduced in Section 3.1, is given by

$$\mathcal{U} = \{\mathbf{1}, \Psi, \mathcal{I}(\Psi^2), X_i\},$$

the property (3.9) of the functions F and F^ε is obvious. The bounds (3.12) follow from [Hai14, Prop. 6.12]. \square

Our following aim is to define a discrete model $Z^\varepsilon = (\Pi^\varepsilon, \Gamma^\varepsilon, \Sigma^\varepsilon)$ on \mathcal{T}^{gen} , and to extend it in the canonical way to $\hat{\mathcal{T}}$ as in Remark 4.11. To this end, we postulate, for $s, t \in \mathbf{R}$ and $x, y \in \Lambda_\varepsilon^3$,

$$(\Pi_x^{\varepsilon, t} \Psi)(y) = (K^\varepsilon \star_\varepsilon \xi^\varepsilon)(t, y), \quad \Gamma_{xy}^{\varepsilon, t} \Psi = \Psi, \quad \Sigma_x^{\varepsilon, st} \Psi = \Psi.$$

Furthermore, we denote the function $\bar{\psi}^\varepsilon(t, x) \stackrel{\text{def}}{=} (K^\varepsilon \star_\varepsilon (\Pi_x^{\varepsilon, t} \Psi)^3)(t, x)$ and set

$$\begin{aligned} (\Pi_x^{\varepsilon, t} \bar{\Psi})(y) &= \bar{\psi}^\varepsilon(t, y) - \bar{\psi}^\varepsilon(t, x), \quad \Gamma_{xy}^{\varepsilon, t} \bar{\Psi} = \bar{\Psi} - (\bar{\psi}^\varepsilon(t, y) - \bar{\psi}^\varepsilon(t, x)) \mathbf{1}, \\ \Sigma_x^{\varepsilon, st} \bar{\Psi} &= \bar{\Psi} - (\bar{\psi}^\varepsilon(t, x) - \bar{\psi}^\varepsilon(s, x)) \mathbf{1}. \end{aligned}$$

Postulating the actions of these maps on the abstract polynomials in the standard way, we canonically extend Z^ε to the whole $\hat{\mathcal{T}}$.

Furthermore, we define the renormalisation constants²

$$C_1^{(\varepsilon)} \stackrel{\text{def}}{=} \int_{\mathbf{R} \times \Lambda_\varepsilon^3} (K^\varepsilon(z))^2 dz, \quad C_2^{(\varepsilon)} \stackrel{\text{def}}{=} 2 \int_{\mathbf{R} \times \Lambda_\varepsilon^3} (K^\varepsilon \star_\varepsilon K^\varepsilon)(z)^2 K^\varepsilon(z) dz, \quad (7.2)$$

²One can show that $C_1^{(\varepsilon)} \sim \varepsilon^{-1}$ and $C_2^{(\varepsilon)} \sim \log \varepsilon$.

and use them to define the renormalisation map M^ε as in (7.1). Finally, we define the renormalised model $\hat{Z}^\varepsilon \stackrel{\text{def}}{=} M^\varepsilon Z^\varepsilon$ as in Remark 4.13. Using the model \hat{Z}^ε in (5.19) we obtain a solution to $(\Phi_{3,\varepsilon}^4)$ with $C^{(\varepsilon)} \stackrel{\text{def}}{=} 3C_1^{(\varepsilon)} - 9C_2^{(\varepsilon)}$.

Before giving a proof of Theorem 1.1 we provide some technical results.

7.1 Discrete functions with prescribed singularities

It follows from Proposition 6.2 that the ‘‘strength’’ of singularity of a kernel determines the regularity of the respective distribution. In this section we provide some properties of singular discrete functions. As usual we fix a scaling $\mathfrak{s} = (\mathfrak{s}_0, 1, \dots, 1)$ of \mathbf{R}^{d+1} with $\mathfrak{s}_0 \geq 1$.

For a function K defined on $\mathbf{R} \times \Lambda_\varepsilon^d$ and supported in a ball centered at the origin, we denote by $D_{i,\varepsilon}$ the finite difference derivative, i.e.

$$D_{i,\varepsilon} K(t, x) \stackrel{\text{def}}{=} \varepsilon^{-1} (K(t, x + \varepsilon e_i) - K(t, x)),$$

where $\{e_i\}_{i=1\dots d}$ is the canonical basis of \mathbf{R}^d , and for $k = (k_0, \dots, k_d) \in \mathbf{N}^{d+1}$ we define $D_\varepsilon^k \stackrel{\text{def}}{=} D_t^{k_0} D_{1,\varepsilon}^{k_1} \dots D_{d,\varepsilon}^{k_d}$. Then for $\zeta \in \mathbf{R}$ and $m \geq 0$ we define

$$\|K\|_{\zeta; m, c}^{(\varepsilon)} \stackrel{\text{def}}{=} \max_{|k|_{\mathfrak{s}} \leq m} \sup_{\|z\|_{\mathfrak{s}} \geq c\varepsilon} \|z\|_{\mathfrak{s}}^{|k|_{\mathfrak{s}} - \zeta} |D_\varepsilon^k K(z)| + \sup_{\|z\|_{\mathfrak{s}} < c\varepsilon} \varepsilon^{-\zeta} |K(z)|,$$

where $z \in \mathbf{R} \times \Lambda_\varepsilon^d$ and $k \in \mathbf{N}^{d+1}$. If this quantity is bounded, we will say that K is of order ζ .

In what follows we provide properties of such functions analogous to [Hai14, Sec. 10.3]. The following result establishes how product and convolution change orders of singularities.

Lemma 7.2. *Let K_1 and K_2 be of orders ζ_1 and ζ_2 respectively. Then $K_1 K_2$ is of order $\zeta \stackrel{\text{def}}{=} \zeta_1 + \zeta_2$ and for every $m \geq 0$ and $c > 0$ one has*

$$\|K_1 K_2\|_{\zeta; m, c+m}^{(\varepsilon)} \lesssim \|K_1\|_{\zeta_1; m, c}^{(\varepsilon)} \|K_2\|_{\zeta_2; m, c}^{(\varepsilon)}. \quad (7.3)$$

If $\zeta_1 \wedge \zeta_2 > -|\mathfrak{s}|$ and $\bar{\zeta} \stackrel{\text{def}}{=} \zeta_1 + \zeta_2 + |\mathfrak{s}| < 0$, then one has the bound

$$\|K_1 \star_\varepsilon K_2\|_{\bar{\zeta}; m, \bar{c}}^{(\varepsilon)} \lesssim \|K_1\|_{\zeta_1; m, c}^{(\varepsilon)} \|K_2\|_{\zeta_2; m, c}^{(\varepsilon)}, \quad (7.4)$$

for some $\bar{c} > c$ depending on m . In the case $\bar{\zeta} \in \mathbf{R}_+ \setminus \mathbf{N}$, the function

$$\bar{K}(z) \stackrel{\text{def}}{=} (K_1 \star_\varepsilon K_2)(z) - \sum_{|k|_{\mathfrak{s}} < \bar{\zeta}} \frac{(z)_{k,\varepsilon}}{k!} D_\varepsilon^k (K_1 \star_\varepsilon K_2)(0),$$

where $(t, x)_{k,\varepsilon} \stackrel{\text{def}}{=} t^{k_0} \prod_{i \neq 0} \prod_{0 \leq j < k_i} (x_i - \varepsilon j)$, satisfies the bound

$$\|\bar{K}\|_{\bar{\zeta}; m, \bar{c}}^{(\varepsilon)} \lesssim \|K_1\|_{\zeta_1; \bar{m}, c}^{(\varepsilon)} \|K_2\|_{\zeta_2; \bar{m}, c}^{(\varepsilon)}, \quad (7.5)$$

with $\bar{m} = m \vee (\lfloor \bar{\zeta} \rfloor + \mathfrak{s}_0)$. In all these estimates the proportionality constants depend only on the support of the functions K_i .

Proof. The bound (7.3) follows from the Leibniz rule for the discrete derivative:

$$D_\varepsilon^k(K_1 K_2)(z) = \sum_{l=(l_0, \bar{l}) \leq k} \binom{k}{l} D_\varepsilon^l K_1(z) D_\varepsilon^{k-l} K_2(z + (0, \varepsilon \bar{l})), \quad (7.6)$$

where $l_0 \in \mathbf{N}$ and $\bar{l} \in \mathbf{N}^d$, and from the simple estimate $\|z + (0, \varepsilon \bar{l})\|_{\mathfrak{s}} \geq \varepsilon c$ as soon as $\|z\|_{\mathfrak{s}} \geq (c + m)\varepsilon$. The bounds (7.4) and (7.5) can be shown similarly to [Hai14, Lem. 10.14], but using the Leibniz rule (7.6), summation by parts for the discrete derivative and the fact that the products $(z)_{k, \varepsilon}$ play the role of polynomials for the discrete derivative. \square

Sometimes we need to bound an increment of a singular function. The following lemma provides a relevant result.

Lemma 7.3. *Let K be of order $\zeta \leq 0$. Then, for every $\kappa \in [0, 1]$, one has*

$$|K(t, x_1) - K(t, x_2)| \lesssim |x_1 - x_2|^\kappa (\|t, x_1\|_{\mathfrak{s}, \varepsilon}^{\zeta - \kappa} + \|t, x_2\|_{\mathfrak{s}, \varepsilon}^{\zeta - \kappa}) \|K\|_{\zeta; 1, c}^{(\varepsilon)},$$

for $t \in \mathbf{R}$, $x_1, x_2 \in \Lambda_\varepsilon^d$ and $\|z\|_{\mathfrak{s}, \varepsilon} \stackrel{\text{def}}{=} \|z\|_{\mathfrak{s}} \vee \varepsilon$.

Proof. The proof is almost identical to that of [Hai14, Lem. 10.18]. \square

For a discrete singular function K , we define the function $\mathcal{R}_\varepsilon K$ by

$$(\mathcal{R}_\varepsilon K)(\varphi) \stackrel{\text{def}}{=} \int_{\mathbf{R} \times \Lambda_\varepsilon^d} K(z) (\varphi(z) - \varphi(0)) dz,$$

for every compactly supported test function φ on \mathbf{R}^{d+1} . The following result can be proved similarly to [Hai14, Lem. 10.16].

Lemma 7.4. *Let K_1 and K_2 be of orders ζ_1 and ζ_2 with $\zeta_1 \in (-|\mathfrak{s}| - 1, -|\mathfrak{s}|]$ and $\zeta_2 \in (-2|\mathfrak{s}| - \zeta_1, 0]$. Then the function $(\mathcal{R}_\varepsilon K_1) \star_\varepsilon K_2$ is of order $\bar{\zeta} \stackrel{\text{def}}{=} \zeta_1 + \zeta_2 + |\mathfrak{s}|$ and, for any $m \geq 0$, one has*

$$\|(\mathcal{R}_\varepsilon K_1) \star_\varepsilon K_2\|_{\bar{\zeta}; m, \bar{c}}^{(\varepsilon)} \lesssim \|K_1\|_{\zeta_1; m, c}^{(\varepsilon)} \|K_2\|_{\zeta_2; m + \mathfrak{s}_0, c}^{(\varepsilon)},$$

where $\bar{c} > c$ depends on m .

The following result shows how certain convolutions change singular functions. Its proof is similar to [Hai14, Lem. 10.17].

Lemma 7.5. *Let, for some $\bar{\varepsilon} \in [\varepsilon, 1]$, the function $\psi^{\bar{\varepsilon}, \varepsilon} : \mathbf{R} \times \Lambda_\varepsilon^d \rightarrow \mathbf{R}$ be smooth in time, supported in the ball $B(0, R\bar{\varepsilon}) \subset \mathbf{R}^{d+1}$, for some $R \geq 1$, and satisfies*

$$\int_{\mathbf{R} \times \Lambda_\varepsilon^d} \psi^{\bar{\varepsilon}, \varepsilon}(z) dz = 1, \quad |D_\varepsilon^k \psi^{\bar{\varepsilon}, \varepsilon}(z)| \lesssim \bar{\varepsilon}^{-|\mathfrak{s}| - |k|_{\mathfrak{s}}},$$

for all $z \in \mathbf{R} \times \Lambda_\varepsilon^d$ and all $k \in \mathbf{N}^{d+1}$ such that $|k|_{\mathfrak{s}} \leq m + \mathfrak{s}_0$ with some $m \geq 0$. If K is of order $\zeta \in (-|\mathfrak{s}|, 0)$, then for all $\kappa \in (0, 1]$ one has the bound

$$\|K - K \star_\varepsilon \psi\|_{\zeta - \kappa; m, \bar{c}}^{(\varepsilon)} \lesssim \bar{\varepsilon}^\kappa \|K\|_{\zeta; m + \mathfrak{s}_0, c}^{(\varepsilon)},$$

where $\bar{c} > c$ depends on m and R .

7.2 Proof of the convergence result

Using the results from the previous section, we are ready to prove Theorem 1.1.

Proof of Theorem 1.1. In order to prove the claim, we proceed as in [Par75] and introduce intermediate equations driven by a smooth noise. Precisely, we take a function $\psi : \mathbf{R}^4 \rightarrow \mathbf{R}$ which is smooth, compactly supported and integrates to 1, and for some $\bar{\varepsilon} \in [\varepsilon, 1]$ we define $\psi^{\bar{\varepsilon}}(t, x) \stackrel{\text{def}}{=} \bar{\varepsilon}^{-|s|} \psi(\bar{\varepsilon}^{-2}t, \bar{\varepsilon}^{-1}x)$ and the mollified noise $\xi^{\bar{\varepsilon},0} \stackrel{\text{def}}{=} \xi \star \psi^{\bar{\varepsilon}}$. Then we denote by $\Phi^{\bar{\varepsilon},0}$ the global solution of

$$\partial_t \Phi^{\bar{\varepsilon},0} = \Delta \Phi^{\bar{\varepsilon},0} - ((\Phi^{\bar{\varepsilon},0})^2 - C^{(\bar{\varepsilon},0)}) \Phi^{\bar{\varepsilon},0} + \xi^{\bar{\varepsilon},0}, \quad \Phi^{\bar{\varepsilon},0}(0, \cdot) = \Phi_0(\cdot),$$

where $C^{(\bar{\varepsilon},0)}$ is as in [Hai14, Thm. 10.22 and Eq. 9.21].

For $z \in \mathbf{R}^4$, we define the map $\Pi_z^{(\bar{\varepsilon})} : \mathcal{T} \rightarrow \mathcal{C}^\infty(\mathbf{R}^4)$ as in the proof of [Hai14, Thm. 10.22] via the noise $\xi^{\bar{\varepsilon},0}$. It follows from the bounds obtained in the proof of [Hai14, Thm. 10.22] that we can define an inhomogeneous model $Z^{\bar{\varepsilon},0}$ by setting

$$\langle \Pi_x^{\bar{\varepsilon},t} \tau, \varphi_x^\lambda \rangle \stackrel{\text{def}}{=} \lim_{\nu \rightarrow 0} \langle \Pi_{(t,x)}^{(\bar{\varepsilon})} \tau, \psi_t^\nu \varphi_x^\lambda \rangle,$$

where $\psi_t^\nu(s) \stackrel{\text{def}}{=} \nu^{-1} \psi((s-t)\nu^{-1})$ and $\psi : \mathbf{R} \rightarrow \mathbf{R}$ is smooth, supported in the unit ball and $\|\psi\|_{C^r} \leq 1$. In the same way we can define the corresponding limiting model Z on $\hat{\mathcal{F}}$. Furthermore, similarly to how it was done in Proposition 6.2 and Theorem 6.1 we can show that

$$\mathbf{E} \|Z\|_{\delta,\gamma;T}^p \lesssim 1, \quad \mathbf{E} \|Z; Z^{\bar{\varepsilon},0}\|_{\delta,\gamma;T}^p \lesssim \bar{\varepsilon}^{\theta p}, \quad (7.7)$$

for any $p \geq 1$, $\delta > 0$ and $\theta > 0$ sufficiently small.

In order to discretise the noise $\xi^{\bar{\varepsilon},0}$, we define the function

$$\psi^{\bar{\varepsilon},\varepsilon}(t, x) \stackrel{\text{def}}{=} \varepsilon^{-d} \int_{\mathbf{R}^d} \psi^{\bar{\varepsilon}}(t, y) \mathbf{1}_{|y-x| \leq \varepsilon/2} dy, \quad (t, x) \in \mathbf{R} \times \Lambda_\varepsilon^d,$$

and the discrete noise $\xi^{\bar{\varepsilon},\varepsilon} \stackrel{\text{def}}{=} \psi^{\bar{\varepsilon},\varepsilon} \star_\varepsilon \xi^\varepsilon$, where ξ^ε is given in (1.1). We define the discrete model $Z^{\bar{\varepsilon},\varepsilon}$ by substituting each occurrence of ξ^ε , $C_1^{(\varepsilon)}$ and $C_2^{(\varepsilon)}$ in the definition of Z^ε by $\xi^{\bar{\varepsilon},\varepsilon}$, $C_1^{(\bar{\varepsilon},\varepsilon)}$ and $C_2^{(\bar{\varepsilon},\varepsilon)}$ respectively, where $C_1^{(\bar{\varepsilon},\varepsilon)}$ is defined as in (7.2), but via the kernel $K^{\bar{\varepsilon},\varepsilon} \stackrel{\text{def}}{=} K^\varepsilon \star_\varepsilon \psi^{\bar{\varepsilon},\varepsilon}$, and $C_2^{(\bar{\varepsilon},\varepsilon)}$ is defined by replacing $K^\varepsilon \star_\varepsilon K^\varepsilon$ by $K^{\bar{\varepsilon},\varepsilon} \star_\varepsilon K^{\bar{\varepsilon},\varepsilon}$ in the second expression in (7.2). Furthermore, using $\|K^\varepsilon\|_{-3;r}^{(\varepsilon)} \leq C$, which follows from Lemma 5.4, and proceeding exactly as in the proof of [Hai14, Thm. 10.22], but exploiting Proposition 6.2 and the results from Section 7.1 instead of their continuous counterparts, we obtain the bounds (6.3) for each $\tau \in \mathcal{F}^{\text{gen}} \setminus \mathcal{F}_{\text{poly}}$, and (6.4) for each $\tau \in \hat{\mathcal{F}}^-$, uniformly in $\varepsilon \leq \bar{\varepsilon}$ and for $\delta > 0$ small enough. We also obtain the respective bounds on the differences $Z^{\bar{\varepsilon},\varepsilon} - Z^\varepsilon$, with the proportionality constants of orders $\bar{\varepsilon}^{2\theta}$ with $\theta > 0$ sufficiently small. For this, we can use Lemma 7.5, because $\psi^{\bar{\varepsilon},\varepsilon}$ satisfies the required conditions, which follows from the properties of ψ . Thus, Theorem 6.1 yields

$$\mathbf{E} \left[\|Z^\varepsilon\|_{\delta,\gamma;T}^{(\varepsilon)} \right]^p \lesssim 1, \quad \mathbf{E} \left[\|Z^{\bar{\varepsilon},\varepsilon}; Z^\varepsilon\|_{\delta,\gamma;T}^{(\varepsilon)} \right]^p \lesssim \bar{\varepsilon}^{\theta p}, \quad (7.8)$$

uniformly in $\varepsilon \leq \bar{\varepsilon}$, for any $T > 0$ and $p \geq 1$.

Let $\Phi^{\bar{\varepsilon}, \varepsilon}$ be the solution of $(\Phi_{3, \varepsilon}^4)$, driven by the noise $\xi^{\bar{\varepsilon}, \varepsilon}$, with the renormalisation constant $C^{(\bar{\varepsilon}, \varepsilon)} \stackrel{\text{def}}{=} 3C_1^{(\bar{\varepsilon}, \varepsilon)} - 9C_2^{(\bar{\varepsilon}, \varepsilon)}$. Then we can bound, for some stopping time T_ε ,

$$\begin{aligned} \mathbf{E} \|\Phi; \Phi^\varepsilon\|_{C_{\bar{\eta}, T_\varepsilon}^{\delta, \alpha+2}}^{(\varepsilon)} &\leq \mathbf{E} \|\Phi - \Phi^{\bar{\varepsilon}, 0}\|_{C_{\bar{\eta}, T_\varepsilon}^{\delta, \alpha+2}} + \mathbf{E} \|\Phi^{\bar{\varepsilon}, 0}; \Phi^{\bar{\varepsilon}, \varepsilon}\|_{C_{\bar{\eta}, T_\varepsilon}^{\delta, \alpha+2}}^{(\varepsilon)} \\ &\quad + \mathbf{E} \|\Phi^{\bar{\varepsilon}, \varepsilon} - \Phi^\varepsilon\|_{C_{\bar{\eta}, T_\varepsilon}^{\delta, \alpha+2}}^{(\varepsilon)}, \end{aligned} \quad (7.9)$$

and to take first the limit $\varepsilon \rightarrow 0$ and after that $\bar{\varepsilon} \rightarrow 0$. Convergence of the first term in (7.9) follows from Theorems 2.11, 3.9 and (7.7), for a suitable T_ε . We can use Theorems 5.7, 4.5 and (7.8) to obtain convergence of the last term in (7.9).

Now, we turn to the second term in (7.9). It follows from our definitions that we have $\xi^{\bar{\varepsilon}, \varepsilon} = \varrho^{\bar{\varepsilon}, \varepsilon} \star \xi$, where

$$\varrho^{\bar{\varepsilon}, \varepsilon}(t, x) \stackrel{\text{def}}{=} \varepsilon^{-d} \int_{\Lambda_\varepsilon^d} \psi^{\bar{\varepsilon}, \varepsilon}(t, y) \mathbf{1}_{|y-x| \leq \varepsilon/2} dy.$$

Moreover, for $z = (t, x) \in \mathbf{R} \times \Lambda_\varepsilon^d$ one has the identity

$$(\psi^{\bar{\varepsilon}} - \varrho^{\bar{\varepsilon}, \varepsilon})(z) = \varepsilon^{-2d} \int_{\Lambda_\varepsilon^d} \int_{\mathbf{R}^d} (\psi^{\bar{\varepsilon}}(t, x) - \psi^{\bar{\varepsilon}}(t, u)) \mathbf{1}_{|u-y| \leq \varepsilon/2} \mathbf{1}_{|y-x| \leq \varepsilon/2} du dy,$$

from which we immediately obtain the bound

$$\sup_{z \in \mathbf{R} \times \Lambda_\varepsilon^d} |D_t^k (\psi^{\bar{\varepsilon}} - \varrho^{\bar{\varepsilon}, \varepsilon})(z)| \lesssim \varepsilon \bar{\varepsilon}^{-|s| - ks_0 - 1},$$

for every $k \in \mathbf{N}$. Hence, we can conclude from [BS08] that the second term in (7.9) vanishes as $\varepsilon \rightarrow 0$, as soon as $\bar{\varepsilon}$ is fixed, which finishes the proof. \square

Proof of Corollary 1.2. Let ξ be space-time white noise on some probability space $(\Omega, \mathcal{F}, \mathbf{P})$, and let its discretisation ξ^ε be given by (1.1). Let furthermore Φ_0^ε be a random variable on the same probability space which is independent of ξ and such that the solution to $(\Phi_{3, \varepsilon}^4)$ with the nearest neighbours approximate Laplacian Δ^ε and driven by ξ^ε is stationary. We denote by μ_ε its stationary distribution, which we view as a measure on \mathcal{C}^α with α as in (1.2), by extending it in a piecewise constant fashion. It then follows from [BFS83] (by combining Eq. 8.2 and Thm. 6.1 in that article) that the sequence μ_ε is tight in \mathcal{C}^α as $\varepsilon \rightarrow 0$ with uniformly bounded moments of all orders, so we can choose a subsequence (which we also denote by μ_ε), weakly converging to an accumulation point μ . Actually, combining this with [Par75] shows that μ is unique and coincides with the Φ_3^4 measure constructed in [Fel74]. In particular, if we view Φ_0^ε as an element of \mathcal{C}^α by piecewise constant extension, we can and will assume by Skorokhod's representation theorem that Φ_0^ε converges almost surely as $\varepsilon \rightarrow 0$ to a limit $\Phi_0 \in \mathcal{C}^\alpha$.

Before we proceed, we introduce the space $\bar{\mathcal{C}} \stackrel{\text{def}}{=} \mathcal{C}_{\bar{\eta}}^{\delta, \alpha}([0, 1], \mathbf{T}^3) \cup \{\infty\}$ (the latter Hölder space is a subspace of $\mathcal{C}_{\bar{\eta}}^{\delta, \alpha}([0, 1], \mathbf{R}^3)$, containing the spatially periodic distributions), for δ , α and $\bar{\eta}$ as in (1.2), and equipped with the metric such that

$$\begin{aligned} d(\zeta, \infty) &\stackrel{\text{def}}{=} d(\infty, \zeta) \stackrel{\text{def}}{=} (1 + \|\zeta\|_{\mathcal{C}_{\bar{\eta}}^{\delta, \alpha}})^{-1}, \quad \zeta \neq \infty, \\ d(\zeta_1, \zeta_2) &\stackrel{\text{def}}{=} \min\{\|\zeta_1 - \zeta_2\|_{\mathcal{C}_{\bar{\eta}}^{\delta, \alpha}}, d(\zeta_1, \infty) + d(\zeta_2, \infty)\}, \quad \zeta_i \neq \infty. \end{aligned}$$

Denote now by Φ^ε the solution to $(\Phi_{3, \varepsilon}^4)$ with initial condition Φ_0^ε and by Φ the solution to (Φ_3^4) with initial condition Φ_0 . We can view these as $\bar{\mathcal{C}}$ -valued random variables by postulating that $\Phi = \infty$ if its lifetime is smaller than 1. (The lifetime of Φ^ε is always infinite for fixed ε .)

Since the assumptions of Theorem 1.1 are fulfilled, the convergence (1.2) holds and, since solutions blow up at time T^* , this implies that $d(\Phi^\varepsilon, \Phi) \rightarrow 0$ in probability, as $\varepsilon \rightarrow 0$. In order to conclude, it remains to show that $\mathbf{P}(\Phi = \infty) = 0$. In particular, since the only point of discontinuity of the evaluation maps $\Phi \mapsto \Phi(t, \cdot)$ on $\bar{\mathcal{C}}$ is ∞ , this would then immediately show not only that solutions Φ live up to time 1 (and therefore any time) almost surely, but also that μ is invariant for Φ . To show that $\Phi \neq \infty$ a.s., it suffices to prove tightness of Φ^ε in $\mathcal{C}_{\bar{\eta}}^{\delta, \alpha}([0, 1], \mathbf{T}^3)$.

To prove this tightness, we take any $\delta' > \delta$ and $\alpha' > \alpha$ such that the convergence (1.2) still holds in $\mathcal{C}_{\bar{\eta}}^{\delta', \alpha'}([0, T_\varepsilon], \mathbf{T}^3)$, for a sequence of stopping times T_ε . Our aim is to show that for every $\bar{\varepsilon} > 0$ there exists a constant $C_{\bar{\varepsilon}} > 0$ such that

$$\mathbf{P}\left(\|\Phi^\varepsilon\|_{\mathcal{C}_{\bar{\eta}, 1}^{\delta', \alpha'}} \leq C_{\bar{\varepsilon}}\right) \geq 1 - \bar{\varepsilon}. \quad (7.10)$$

Since $\mathcal{C}_{\bar{\eta}}^{\delta', \alpha'}([0, 1], \mathbf{T}^3)$ is compactly embedded into $\mathcal{C}_{\bar{\eta}}^{\delta, \alpha}([0, 1], \mathbf{T}^3)$, this bound implies the required tightness of the stationary solutions.

We fix $\bar{\varepsilon} > 0$ in what follows and work with a generic constant $C_{\bar{\varepsilon}} > 0$, whose value will be chosen later. For integers $K \geq 2$ and $i \in \{0, \dots, K-2\}$, we denote

$$Q_{K, i}^\varepsilon \stackrel{\text{def}}{=} \|\Phi^\varepsilon\|_{\mathcal{C}_{\bar{\eta}, [i/K, (i+2)/K]}^{\delta', \alpha'}},$$

where the norm $\|\cdot\|_{\mathcal{C}_{\bar{\eta}, [T_1, T_2]}^{\delta', \alpha'}}$ is defined as in (1.5), but on the time interval $[T_1, T_2]$ and with a blow-up at T_1 . Splitting the time interval $(0, 1]$ in (1.5) into subintervals of length $1/K$, and deriving estimates on each subinterval, one gets

$$\|\Phi^\varepsilon\|_{\mathcal{C}_{\bar{\eta}, 1}^{\delta', \alpha'}} \leq Q_{K, 0}^\varepsilon + \sum_{i=1}^{K-1} (i+1)^{-\eta/2} Q_{K, i-1}^\varepsilon \leq \tilde{C} K^{-\eta/2} \sum_{i=0}^{K-2} Q_{K, i}^\varepsilon,$$

if $\eta \leq 0$, and for some \tilde{C} independent of K and ε . Since, by stationarity, the random variables $Q_{K, i}^\varepsilon$ all have the same law, it follows that

$$\mathbf{P}\left(\|\Phi^\varepsilon\|_{\mathcal{C}_{\bar{\eta}, 1}^{\delta', \alpha'}} \geq C_{\bar{\varepsilon}}\right) \leq \mathbf{P}\left(\tilde{C} K^{-\eta/2} \sum_{i=0}^{K-2} Q_{K, i}^\varepsilon \geq C_{\bar{\varepsilon}}\right)$$

$$\leq K\mathbf{P}\left(\|\Phi^\varepsilon\|_{C_{\bar{\eta}, 2/K}^{\delta', \alpha'}} \geq \tilde{C}^{-1}K^{\eta/2}C_{\bar{\varepsilon}}\right), \quad (7.11)$$

To make the notation concise, we write $\tilde{C}_{K, \bar{\varepsilon}} \stackrel{\text{def}}{=} \tilde{C}^{-1}K^{\eta/2}C_{\bar{\varepsilon}}$. Furthermore, in order to have a uniform bound on the initial data and the model, we use the following estimate

$$\begin{aligned} \mathbf{P}\left(\|\Phi^\varepsilon\|_{C_{\bar{\eta}, 2/K}^{\delta', \alpha'}} \geq \tilde{C}_{K, \bar{\varepsilon}}\right) &\leq \mathbf{P}\left(\|\Phi^\varepsilon\|_{C_{\bar{\eta}, 2/K}^{\delta', \alpha'}} \geq \tilde{C}_{K, \bar{\varepsilon}} \mid \|\Phi_0^\varepsilon\|_{C^\eta} \leq L, \|Z^\varepsilon\|_{\delta', \gamma; 1}^{(\varepsilon)} \leq L\right) \\ &+ \mathbf{P}\left(\|\Phi_0^\varepsilon\|_{C^\eta} > L\right) + \mathbf{P}\left(\|Z^\varepsilon\|_{\delta', \gamma; 1}^{(\varepsilon)} > L\right), \end{aligned} \quad (7.12)$$

valid for every L , where $\gamma > 0$ is as in the proof of Theorem 1.1.

Recalling that [BFS83, Sec. 8] yields uniform bounds on all moments of μ_ε , and using the first bound in (7.8), Markov's inequality implies that

$$\mathbf{P}\left(\|\Phi_0^\varepsilon\|_{C^\eta} > L\right) \leq B_1L^{-q}, \quad \mathbf{P}\left(\|Z^\varepsilon\|_{\delta', \gamma; 1}^{(\varepsilon)} > L\right) \leq B_2L^{-q}, \quad (7.13)$$

for any $q \geq 1$, and for constant B_1 and B_2 independent of ε and L .

Turning to the first term in (7.12), it follows from the fixed point argument in the proof of Theorem 5.7 and the bound (4.5a), that there exists $\tilde{p} \geq 1$ such that one has the bound

$$\|\Phi^\varepsilon\|_{C_{\bar{\eta}, 2/K}^{\delta', \alpha'}}^p \leq B_3L^{3p},$$

with B_3 being independent of ε and L , as soon as $\|\Phi_0^\varepsilon\|_{C^\eta} \leq L$, $\|Z^\varepsilon\|_{\delta', \gamma; 1}^{(\varepsilon)} \leq L$, $K \geq L^{\tilde{p}}$ and $L \geq 2$. In particular, the first term vanishes if we can ensure that

$$\tilde{C}_{K, \bar{\varepsilon}} \geq B_3L^{3p} \quad (7.14)$$

Choosing first L large enough so that the contribution of the two terms in (7.13) is smaller than $\bar{\varepsilon}/3$, then K large enough so that $K \geq L^{\tilde{p}}$, and finally $C_{\bar{\varepsilon}}$ large enough so that (7.14) holds, the claim follows. \square

References

- [AR91] S. ALBEVERIO and M. RÖCKNER. Stochastic differential equations in infinite dimensions: solutions via Dirichlet forms. *Probab. Theory Related Fields* **89**, no. 3, (1991), 347–386. doi:10.1007/BF01198791.
- [BCD11] H. BAHOURI, J.-Y. CHEMIN, and R. DANCHIN. *Fourier analysis and nonlinear partial differential equations*, vol. 343 of *Grundlehren der Mathematischen Wissenschaften [Fundamental Principles of Mathematical Sciences]*. Springer, Heidelberg, 2011. doi:10.1007/978-3-642-16830-7.
- [BFS83] D. C. BRYDGES, J. FRÖHLICH, and A. D. SOKAL. A new proof of the existence and nontriviality of the continuum φ_2^4 and φ_3^4 quantum field theories. *Comm. Math. Phys.* **91**, no. 2, (1983), 141–186.
- [BG97] L. BERTINI and G. GIACOMIN. Stochastic Burgers and KPZ equations from particle systems. *Comm. Math. Phys.* **183**, no. 3, (1997), 571–607. doi:10.1007/s002200050044.

- [Bou94] J. BOURGAIN. Periodic nonlinear Schrödinger equation and invariant measures. *Comm. Math. Phys.* **166**, no. 1, (1994), 1–26. doi:[10.1007/BF02099299](https://doi.org/10.1007/BF02099299).
- [BS08] S. C. BRENNER and L. R. SCOTT. *The mathematical theory of finite element methods*, vol. 15 of *Texts in Applied Mathematics*. Springer, New York, third ed., 2008. doi:[10.1007/978-0-387-75934-0](https://doi.org/10.1007/978-0-387-75934-0).
- [CC13] R. CATELLIER and K. CHOUK. Paracontrolled distributions and the 3-dimensional stochastic quantization equation (2013). arXiv:[1310.6869](https://arxiv.org/abs/1310.6869).
- [Dau88] I. DAUBECHIES. Orthonormal bases of compactly supported wavelets. *Comm. Pure Appl. Math.* **41**, no. 7, (1988), 909–996. doi:[10.1002/cpa.3160410705](https://doi.org/10.1002/cpa.3160410705).
- [Dau92] I. DAUBECHIES. *Ten lectures on wavelets*, vol. 61 of *CBMS-NSF Regional Conference Series in Applied Mathematics*. Society for Industrial and Applied Mathematics (SIAM), Philadelphia, PA, 1992. doi:[10.1137/1.9781611970104](https://doi.org/10.1137/1.9781611970104).
- [DPD03] G. DA PRATO and A. DEBUSSCHE. Strong solutions to the stochastic quantization equations. *Ann. Probab.* **31**, no. 4, (2003), 1900–1916. doi:[10.1214/aop/1068646370](https://doi.org/10.1214/aop/1068646370).
- [DPZ14] G. DA PRATO and J. ZABCZYK. *Stochastic equations in infinite dimensions*, vol. 152 of *Encyclopedia of Mathematics and its Applications*. Cambridge University Press, Cambridge, second ed., 2014. doi:[10.1017/CB09781107295513](https://doi.org/10.1017/CB09781107295513).
- [Fel74] J. FELDMAN. The $\lambda\phi_3^4$ field theory in a finite volume. *Comm. Math. Phys.* **37**, (1974), 93–120.
- [GIP15] M. GUBINELLI, P. IMKELLER, and N. PERKOWSKI. Paracontrolled distributions and singular PDEs. *Forum of Mathematics, Pi* **3**, no. 6, (2015), 1–75. arXiv:[1210.2684](https://arxiv.org/abs/1210.2684). doi:[10.1017/fmp.2015.2](https://doi.org/10.1017/fmp.2015.2).
- [GLP99] G. GIACOMIN, J. L. LEBOWITZ, and E. PRESUTTI. Deterministic and stochastic hydrodynamic equations arising from simple microscopic model systems. In *Stochastic partial differential equations: six perspectives*, vol. 64 of *Math. Surveys Monogr.*, 107–152. Amer. Math. Soc., Providence, RI, 1999. doi:[10.1090/surv/064/03](https://doi.org/10.1090/surv/064/03).
- [GP15] M. GUBINELLI and N. PERKOWSKI. KPZ reloaded. *ArXiv e-prints* (2015). arXiv:[1508.03877](https://arxiv.org/abs/1508.03877).
- [GRS75] F. GUERRA, L. ROSEN, and B. SIMON. The $P(\phi)_2$ Euclidean quantum field theory as classical statistical mechanics. I, II. *Ann. of Math. (2)* **101**, (1975), 111–189; *ibid.* (2) **101** (1975), 191–259.
- [Gub04] M. GUBINELLI. Controlling rough paths. *J. Funct. Anal.* **216**, no. 1, (2004), 86–140. arXiv:[math/0306433](https://arxiv.org/abs/math/0306433). doi:[10.1016/j.jfa.2004.01.002](https://doi.org/10.1016/j.jfa.2004.01.002).
- [Hai11] M. HAIRER. Rough stochastic PDEs. *Comm. Pure Appl. Math.* **64**, no. 11, (2011), 1547–1585. arXiv:[1008.1708](https://arxiv.org/abs/1008.1708). doi:[10.1002/cpa.20383](https://doi.org/10.1002/cpa.20383).
- [Hai13] M. HAIRER. Solving the KPZ equation. *Ann. of Math. (2)* **178**, no. 2, (2013), 559–664. arXiv:[1109.6811](https://arxiv.org/abs/1109.6811). doi:[10.4007/annals.2013.178.2.4](https://doi.org/10.4007/annals.2013.178.2.4).
- [Hai14] M. HAIRER. A theory of regularity structures. *Invent. Math.* **198**, no. 2, (2014), 269–504. arXiv:[1303.5113](https://arxiv.org/abs/1303.5113). doi:[10.1007/s00222-014-0505-4](https://doi.org/10.1007/s00222-014-0505-4).
- [Hai15] M. HAIRER. Regularity structures and the dynamical ϕ_3^4 model. *ArXiv e-prints* (2015). arXiv:[1508.05261](https://arxiv.org/abs/1508.05261).

- [Has80] R. Z. HAS'MINSKIĬ. *Stochastic stability of differential equations*, vol. 7 of *Monographs and Textbooks on Mechanics of Solids and Fluids: Mechanics and Analysis*. Sijthoff & Noordhoff, Alphen aan den Rijn—Germantown, Md., 1980. Translated from the Russian by D. Louvish.
- [HL15] M. HAIRER and C. LABBÉ. Multiplicative stochastic heat equations on the whole space (2015). [arXiv:1504.07162](https://arxiv.org/abs/1504.07162).
- [HM14] M. HAIRER and K. MATETSKI. Optimal rate of convergence for stochastic Burgers-type equations (2014). [arXiv:1504.05134](https://arxiv.org/abs/1504.05134).
- [HMW14] M. HAIRER, J. MAAS, and H. WEBER. Approximating rough stochastic PDEs. *Comm. Pure Appl. Math.* **67**, no. 5, (2014), 776–870. [arXiv:1202.3094](https://arxiv.org/abs/1202.3094). [doi:10.1002/cpa.21495](https://doi.org/10.1002/cpa.21495).
- [Hör55] L. HÖRMANDER. On the theory of general partial differential operators. *Acta Math.* **94**, (1955), 161–248.
- [IW89] N. IKEDA and S. WATANABE. *Stochastic differential equations and diffusion processes*, vol. 24 of *North-Holland Mathematical Library*. North-Holland Publishing Co., Amsterdam; Kodansha, Ltd., Tokyo, second ed., 1989.
- [JLM85] G. JONA-LASINIO and P. K. MITTER. On the stochastic quantization of field theory. *Comm. Math. Phys.* **101**, no. 3, (1985), 409–436.
- [KPZ86] M. KARDAR, G. PARISI, and Y.-C. ZHANG. Dynamic scaling of growing interfaces. *Phys. Rev. Lett.* **56**, no. 9, (1986), 889–892.
- [Kup15] A. KUPIAINEN. Renormalization group and stochastic PDEs. *Annales Henri Poincaré* 1–39. [arXiv:1410.3094](https://arxiv.org/abs/1410.3094). [doi:10.1007/s00023-015-0408-y](https://doi.org/10.1007/s00023-015-0408-y).
- [Lyo98] T. J. LYONS. Differential equations driven by rough signals. *Rev. Mat. Iberoamericana* **14**, no. 2, (1998), 215–310. [doi:10.4171/RMI/240](https://doi.org/10.4171/RMI/240).
- [Mey92] Y. MEYER. *Wavelets and operators*, vol. 37 of *Cambridge Studies in Advanced Mathematics*. Cambridge University Press, Cambridge, 1992. Translated from the 1990 French original by D. H. Salinger.
- [MW14] J.-C. MOURRAT and H. WEBER. Convergence of the two-dimensional dynamic Ising-Kac model to ϕ_2^4 (2014). [arXiv:1410.1179](https://arxiv.org/abs/1410.1179).
- [Nua06] D. NUALART. *The Malliavin calculus and related topics*. Probability and its Applications (New York). Springer-Verlag, Berlin, second ed., 2006.
- [Par75] Y. M. PARK. Lattice approximation of the $(\lambda\phi^4 - \mu\phi)_3$ field theory in a finite volume. *Journal of Mathematical Physics* **16**, no. 5, (1975), 1065–1075. [doi:10.1063/1.522661](https://doi.org/10.1063/1.522661).
- [PW81] G. PARISI and Y. S. WU. Perturbation theory without gauge fixing. *Sci. Sinica* **24**, no. 4, (1981), 483–496.
- [SG73] B. SIMON and R. B. GRIFFITHS. The $(\phi^4)_2$ field theory as a classical Ising model. *Comm. Math. Phys.* **33**, (1973), 145–164.
- [ZZ15] R. ZHU and X. ZHU. Lattice approximation to the dynamical ϕ_3^4 model. *ArXiv e-prints* (2015). [arXiv:1508.05613](https://arxiv.org/abs/1508.05613).