

Random homogenisation of a highly oscillatory singular potential

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Abstract

In this article, we consider the problem of homogenising the linear heat equation perturbed by a rapidly oscillating random potential. We consider the situation where the space-time scaling of the potential's oscillations is *not* given by the diffusion scaling that leaves the heat equation invariant. Instead, we treat the case where spatial oscillations are much faster than temporal oscillations. Under suitable scaling of the amplitude of the potential, we prove convergence to a deterministic heat equation with constant potential, thus completing the results previously obtained in [PP12].

1 Introduction

We consider the parabolic PDE with space-time random potential given by

$$\begin{aligned}\partial_t u^\varepsilon(x, t) &= \partial_x^2 u^\varepsilon(x, t) + \varepsilon^{-\beta} V\left(\frac{x}{\varepsilon}, \frac{t}{\varepsilon^\alpha}\right) u^\varepsilon(x, t), \\ u^\varepsilon(x, 0) &= u_0(x),\end{aligned}\tag{1.1}$$

where $x \in \mathbf{R}$, $t \geq 0$ and V is a stationary centred random field. The homogenisation theory of equations of this type has been studied by a number of authors. The case when V is time-independent was considered in [IPP08, Bal10]. The articles [CKP01, DIPP06] considered a situation where V is a stationary process as a function of time, but periodic in space. Purely periodic / quasiperiodic operators with large potential were also studied in [BLP78, Koz83].

For $\alpha \geq 2$ and $\beta = \frac{\alpha}{2}$, (1.1) was studied in [PP12], where it was shown that its solutions converge as $\varepsilon \rightarrow 0$ to the solutions to

$$\partial_t u(x, t) = \partial_x^2 u(x, t) + \bar{V} u(x, t), \quad u(x, 0) = u_0(x),\tag{1.2}$$

where the constant \bar{V} is given by

$$\bar{V} = \int_0^\infty \Phi(0, t) dt, \quad (1.3)$$

in the case $\alpha > 2$ and

$$\bar{V} = \int_0^\infty \int_{-\infty}^\infty \frac{e^{-\frac{x^2}{4t}}}{2\sqrt{\pi t}} \Phi(x, t) dx dt, \quad (1.4)$$

in the case $\alpha = 2$. Here, $\Phi(x, t) = \mathbf{E}V(0, 0)V(x, t)$ is the correlation function of V which is assumed to decay sufficiently fast.

In the case $0 < \alpha < 2$, it was conjectured in [PP12] that the correct scaling to use in order to obtain a non-trivial limit is $\beta = 1/2 + \alpha/4$, but the corresponding value of \bar{V} was not obtained. Furthermore, the techniques used there seem to break down in this case. The main result of the present article is that the conjecture does indeed hold true and that the solutions to (1.1) do again converge to those of (1.2) as $\varepsilon \rightarrow 0$. This time, the limiting constant \bar{V} is given by

$$\bar{V} = \frac{1}{2\sqrt{\pi}} \int_0^\infty \frac{\bar{\Phi}(t)}{\sqrt{t}} dt, \quad (1.5)$$

where we have set $\bar{\Phi}(s) := \int_{\mathbf{R}} \Phi(x, s) dx$.

Remark 1.1 One can “guess” both (1.3) and (1.5) if we admit that (1.4) holds. Indeed, (1.3) is obtained from (1.4) by replacing $\Phi(x, t)$ by $\Phi(\delta x, t)$ and taking the limit $\delta \rightarrow 0$. This reflects the fact that this corresponds to a situation in which, at the diffusive scale, the temporal oscillations of the potential are faster than the spatial oscillations. Similarly, (1.5) is obtained by replacing $\Phi(x, t)$ with $\delta^{-1}\Phi(\delta^{-1}x, t)$ and then taking the limit $\delta \rightarrow 0$, reflecting the fact that we are in the reverse situation where spatial oscillations are faster. These arguments also allow to guess the correct exponent β in both regimes.

The techniques employed in the present article are very different from [PP12]: instead of relying on probabilistic techniques, we adapt the analytical techniques from [Hai13a].

From now on, we will rewrite (1.1) as

$$\partial_t u^\varepsilon(x, t) = \partial_x^2 u^\varepsilon(x, t) + V_\varepsilon(x, t)u^\varepsilon(x, t), \quad u^\varepsilon(x, 0) = u_0(x),$$

where V_ε is the rescaled potential given by

$$V_\varepsilon(x, t) = \varepsilon^{-(1/2+\alpha/4)} V\left(\frac{x}{\varepsilon}, \frac{t}{\varepsilon^\alpha}\right).$$

Before we proceed, we give a more precise description of our assumptions on the random potential V .

1.1 Assumptions on the potential

Besides some regularity and integrability assumptions, our main assumption will be a sufficiently fast decay of maximal correlations for V . Recall that the “maximal correlation coefficient” of V , subsequently denoted by ϱ , is given by the following definition where, for any given compact set $K \subset \mathbf{R}^2$, we denote by \mathcal{F}_K the σ -algebra generated by $\{V(x, t) : (x, t) \in K\}$.

Definition 1.2 For any $r > 0$, $\varrho(r)$ is the smallest value such that the bound

$$\mathbf{E}(\varphi_1(V)\varphi_2(V)) \leq \varrho(r) \sqrt{\mathbf{E}\varphi_1^2(V) \mathbf{E}\varphi_2^2(V)},$$

holds for any two compact sets K_1, K_2 such that

$$d(K_1, K_2) \stackrel{\text{def}}{=} \inf_{(x_1, t_1) \in K_1} \inf_{(x_2, t_2) \in K_2} (|x_1 - x_2| + |t_1 - t_2|) \geq r,$$

and any two r.v.’s $\varphi_i(V)$ such that $\varphi_i(V)$ is \mathcal{F}_{K_i} -measurable and $\mathbf{E}\varphi_i(V) = 0$.

Note that ϱ is a decreasing function. With this notation at hand, we then make the following assumption:

Assumption 1.3 *The field V is stationary, centred, continuous, and \mathcal{C}^1 in the x -variable. Furthermore,*

$$\mathbf{E}(|V(x, t)|^p + |\partial_x V(x, t)|^p) < \infty$$

for every $p > 0$.

For most of our results, we will furthermore require that the correlations of V decay sufficiently fast in the following sense:

Assumption 1.4 *The maximal correlation function ϱ from Definition 1.2 satisfies $\varrho(R) \lesssim (1 + R)^{-q}$ for every $q > 0$.*

Remark 1.5 Retracing the steps of our proof, one can see that in order to obtain our main result, Theorem 1.8, we actually only need this bound for some sufficiently large q . Similarly, the assumption on the x -differentiability of V is not absolutely necessary, but simplifies some of our arguments.

Let us first give a few examples of random fields satisfying our assumptions.

Example 1.6 Take a measure space (\mathcal{M}, ν) with some finite measure ν and a function $\psi: \mathcal{M} \times \mathbf{R}^2 \rightarrow \mathbf{R}$ such that

$$\sup_{m \in \mathcal{M}} \sup_{x, t} \frac{|\psi(m, x, t)| + |\partial_x \psi(m, x, t)|}{1 + |x|^q + |t|^q} < \infty ,$$

for all $q > 0$. Assume furthermore that ψ satisfies the centering condition

$$\int_{\mathbf{R}} \int_{\mathbf{R}} \int_{\mathcal{M}} \psi(m, y, s) \nu(dm) dy ds = 0 .$$

Consider now a realisation μ of the Poisson point process on $\mathcal{M} \times \mathbf{R}^2$ with intensity measure $\nu(dm) dy ds$ and set

$$V(x, t) = \int_{\mathcal{M}} \int_{\mathbf{R}} \int_{\mathbf{R}} \psi(m, y - x, s - t) \mu(dm, dy, ds) .$$

Then V satisfies Assumptions 1.3 and 1.4.

Example 1.7 Take for V a centred Gaussian field with covariance Φ such that

$$\sup_{x, t} \frac{|\Phi(x, t)| + |\partial_x^2 \Phi(x, t)|}{1 + |x|^q + |t|^q} < \infty ,$$

for all $q > 0$. Then V does not quite satisfy Assumptions 1.3 and 1.4 because V and $\partial_x V$ are not necessarily continuous. However, it is easy to check that our proofs still work in this case.

The advantage of Definition 1.2 is that it is invariant under the composition by measurable functions. In particular, given a finite number of independent random fields $\{V_1, \dots, V_k\}$ of the type of Examples 1.6 and 1.7 (or, more generally, any mutually independent fields satisfying Assumptions 1.3 and 1.4) and a function $F: \mathbf{R}^k \rightarrow \mathbf{R}$ such that

1. $\mathbf{E}F(V_1(x, t), \dots, V_k(x, t)) = 0$,
2. F , together with its first partial derivatives, grows no faster than polynomially at infinity.

Then, our results hold with $V(x, t) = F(V_1(x, t), \dots, V_k(x, t))$.

1.2 Statement of the result

Consider the solution to the heat equation with constant potential

$$\begin{aligned}\partial_t u(x, t) &= \partial_x^2 u(x, t) + \bar{V}u(x, t), & t \geq 0, x \in \mathbf{R}; \\ u(x, 0) &= u_0(x),\end{aligned}\tag{1.6}$$

where \bar{V} is defined by (1.5). Then, the main result of this article is the following convergence result:

Theorem 1.8 *Let V be a random potential satisfying Assumptions 1.3 and 1.4, and let $u_0 \in \mathcal{C}^{3/2}(\mathbf{R})$ be of no more than exponential growth. Then, as $\varepsilon \rightarrow 0$, one has $u^\varepsilon(t, x) \rightarrow u(t, x)$ in probability, locally uniformly in $x \in \mathbf{R}$ and $t \geq 0$.*

Remark 1.9 The precise assumption on u_0 is that it belongs to the space $\mathcal{C}_{e^\ell}^{3/2}$ for some $\ell \in \mathbf{R}$, see Section 2.1 below for the definition of this space.

Remark 1.10 The fact that $\mathbf{E}V = 0$ is of course not essential, since one can easily subtract the mean by performing a suitable rescaling of the solution.

To prove Theorem 1.8, we use the standard “trick” to introduce a corrector that “kills” the large potential V_ε to highest order. The less usual feature of this problem is that, in order to obtain the required convergence, it turns out to be advantageous to use *two* correctors, which ensures that the remaining terms can be brought under control. These correctors, which we denote by Y^ε and Z^ε , are given by the solutions to the following inhomogeneous heat equations:

$$\begin{aligned}\partial_t Y^\varepsilon(x, t) &= \partial_x^2 Y^\varepsilon(x, t) + V_\varepsilon(x, t), \\ \partial_t Z^\varepsilon(x, t) &= \partial_x^2 Z^\varepsilon(x, t) + |\partial_x Y^\varepsilon(x, t)|^2 - \bar{V}_\varepsilon(t),\end{aligned}\tag{1.7}$$

where we have set $\bar{V}_\varepsilon(t) = \mathbf{E}|\partial_x Y^\varepsilon(x, t)|^2$. In both cases, we start with the flat (zero) initial condition at $t = 0$. Writing

$$v^\varepsilon(x, t) = u^\varepsilon(x, t) \exp[-(Y^\varepsilon(x, t) + Z^\varepsilon(x, t))],$$

Theorem 1.8 is then a consequence of the following two claims:

1. Both Y^ε and Z^ε converge locally uniformly to 0.
2. The process v^ε converges locally uniformly to the solution u of (1.6).

It is straightforward to verify that v^ε solves the equation

$$\partial_t v^\varepsilon = \partial_x^2 v^\varepsilon + \bar{V}_\varepsilon v^\varepsilon + 2(\partial_x Y^\varepsilon + \partial_x Z^\varepsilon) \partial_x v^\varepsilon + (|\partial_x Z^\varepsilon|^2 + 2\partial_x Z^\varepsilon \partial_x Y^\varepsilon) v^\varepsilon, \quad (1.8)$$

with initial condition u_0 . The second claim will then essentially follow from the first (except that, due to the appearance of nonlinear terms involving the derivatives of the correctors, we need somewhat tighter control than just locally uniform convergence), combined with the fact that the function $\bar{V}_\varepsilon(t)$ converges locally uniformly to the constant \bar{V} .

Remark 1.11 One way of “guessing” the correct forms for the correctors Y^ε and Z^ε is to note the analogy of the problem with that of building solutions to the KPZ equation. Indeed, performing the Cole-Hopf transform $h^\varepsilon = \log u^\varepsilon$, one obtains for h^ε the equation

$$\partial_t h^\varepsilon = \partial_x^2 h^\varepsilon + (\partial_x h^\varepsilon)^2 + V_\varepsilon,$$

which, in the case where V_ε is replaced by space-time white noise, was recently analysed in detail in [Hai13a]. The correctors Y^ε and Z^ε then arise naturally in this analysis as the first terms in the Wild expansion of the KPZ equation.

This also suggests that it would be possible to find a diverging sequence of constants C_ε such that the solutions to

$$\partial_t u^\varepsilon(x, t) = \partial_x^2 u^\varepsilon(x, t) + \varepsilon^{-\frac{1+\alpha}{2}} V\left(\frac{x}{\varepsilon}, \frac{t}{\varepsilon^\alpha}\right) u^\varepsilon(x, t) - C_\varepsilon u^\varepsilon(x, t),$$

converge in law to the solutions to the multiplicative stochastic heat equation driven by space-time white noise. In the non-Gaussian case, this does still seem out of reach at the moment, although some recent progress can be found in [Hai13b].

The proof of Theorem 1.8 now goes as follows. In a first step, which is rather long and technical and constitutes Section 2 below, we obtain sharp a priori bounds for Y^ε and Z^ε in various norms. In a second step, which is performed in Section 3, we then combine these estimates in order to show that the only terms in (1.8) that matter are indeed the first two terms on the right hand side.

Remark 1.12 Throughout this article, the notation $X \lesssim Y$ will be equivalent to the notation $X \leq CY$ for some constant C independent of ε .

2 Estimates of Y^ε and Z^ε

In this section, we shall prove that both Y^ε and Z^ε tend to zero as $\varepsilon \rightarrow 0$, and establish further estimates on those sequences of functions which will be needed for taking the limit of the sequence v^ε . But before doing so, let us first introduce some technical tools which will be needed both in this section and in the last one.

2.1 Weighted Hölder continuous spaces of functions and the heat semigroup

First of all, we define the notion of an admissible weight w as a function $w: \mathbf{R} \rightarrow \mathbf{R}_+$ such that there exists a constant $C \geq 1$ with

$$C^{-1} \leq \frac{w(x)}{w(y)} \leq C, \quad (2.1)$$

for all pairs (x, y) with $|x - y| \leq 1$. Given such an admissible weight w , we then define the space \mathcal{C}_w as the closure of \mathcal{C}_0^∞ under the norm

$$\|f\|_w = \|f\|_{0,w} = \sup_{x \in \mathbf{R}} \frac{|f(x)|}{w(x)}.$$

We also define \mathcal{C}_w^α for $\alpha \in (0, 1)$ as the closure of \mathcal{C}_0^∞ under the norm

$$\|f\|_{\alpha,w} = \|f\|_w + \sup_{|x-y| \leq 1} \frac{|f(x) - f(y)|}{w(x)|x-y|^\alpha}.$$

Similarly, for $\alpha \geq 1$, we define \mathcal{C}_w^α recursively as the closure of \mathcal{C}_0^∞ under the norm

$$\|f\|_{\alpha,w} = \|f\|_w + \|f'\|_{\alpha-1,w}.$$

It is clear that, if w_1 and w_2 are two admissible weights, then so is $w = w_1 w_2$. Furthermore, it is a straightforward exercise to use the Leibniz rule to verify that there exists a constant C such that the bound

$$\|f_1 f_2\|_{\alpha,w} \leq C \|f_1\|_{\alpha_1,w_1} \|f_2\|_{\alpha_2,w_2}, \quad (2.2)$$

holds for every $f_i \in \mathcal{C}_{w_i}^{\alpha_i}$, provided that $\alpha \leq \alpha_1 \wedge \alpha_2$.

We now show that a similar inequality still holds if one of the two Hölder exponents is *negative*. For $\alpha \in (-1, 0)$, we can indeed define weighted spaces of negative ‘‘Hölder regularity’’ by postulating that \mathcal{C}_w^α is the closure of \mathcal{C}_0^∞ under the norm

$$\|f\|_{\alpha,w} = \sup_{|x-y| \leq 1} \frac{|\int_x^y f(z) dz|}{w(x)|x-y|^{\alpha+1}}.$$

In other words, we essentially want the antiderivative of f to belong to $\mathcal{C}_w^{\alpha+1}$, except that we do not worry about its growth.

With these notations at hand, we then have the bound:

Proposition 2.1 *Let w_1 and w_2 be two admissible weights and let $\alpha_1 < 0 < \alpha_2$ be such that $\alpha_2 > |\alpha_1|$. Then, the bound (2.2) holds with $\alpha = \alpha_1$.*

Proof. We only need to show the bound for smooth and compactly supported elements f_1 and f_2 , the general case then follows by density. Denote now by F_1 an antiderivative for f_1 , so that

$$\int_x^y f_1(z)f_2(z) dz = \int_x^y f_2(z) dF_1(z) ,$$

where the right hand side is a Riemann-Stieltjes integral. For any interval $I \subset \mathbf{R}$, we now write

$$\|f\|_{\alpha,I} = \sup_{\{x,y\} \subset I} \frac{|f(x) - f(y)|}{|x - y|^\alpha} .$$

It then follows from Young's inequality [You36] that there exists a constant C depending only on the precise values of the α_i and on the constants appearing in the definition (2.1) of admissibility for the weights w_i , such that

$$\begin{aligned} \left| \int_x^y f_2(z) dF_1(z) \right| &\leq |f_2(x)| |F_1(y) - F_1(x)| + C \|f_2\|_{\alpha_2,[x,y]} \|F_1\|_{\alpha_1+1,[x,y]} |x - y|^{\alpha_1+\alpha_2+1} \\ &\leq w(x) |x - y|^{\alpha_1+1} (\|f_2\|_{0,w_2} \|f_1\|_{\alpha_1,w_1} + C \|f_2\|_{\alpha_2,w_2} \|f_1\|_{\alpha_1,w_1}) , \end{aligned}$$

which is precisely the requested bound. \square

There are two types of admissible weights that will play a crucial role in the sequel:

$$e_\ell(x) \stackrel{\text{def}}{=} \exp(-\ell|x|) , \quad p_\kappa(x) \stackrel{\text{def}}{=} 1 + |x|^\kappa ,$$

where the exponent κ will always be positive, but ℓ could have any sign. One has of course the identity

$$e_\ell \cdot e_m = e_{\ell+m} . \tag{2.3}$$

Furthermore, it is straightforward to verify that there exists a constant C such that the bound

$$p_\kappa(x)e_\ell(x) \leq C\ell^{-\kappa} , \tag{2.4}$$

holds uniformly in $x \in \mathbf{R}$, $\kappa \in (0, 1]$, and $\ell \in (0, 1]$.

Finally, we have the following regularising property of the heat semigroup:

Proposition 2.2 *Let $\alpha \in (-1, \infty)$, let $\beta > \alpha$, and let $\ell, \kappa \in \mathbf{R}$. Then, for every $t > 0$, the operator P_t extends to a bounded operator from $\mathcal{C}_{e_\ell}^\alpha$ to $\mathcal{C}_{e_\ell}^\beta$ and from $\mathcal{C}_{p_\kappa}^\alpha$ to $\mathcal{C}_{p_\kappa}^\beta$. Furthermore, for every $\ell_0 > 0$ and $\kappa_0 > 0$, there exists a constant C such that the bounds*

$$\|P_t f\|_{\beta, e_\ell} \leq C t^{-\frac{\beta-\alpha}{2}} \|f\|_{\alpha, e_\ell}, \quad \|P_t g\|_{\beta, p_\kappa} \leq C t^{-\frac{\beta-\alpha}{2}} \|g\|_{\alpha, p_\kappa},$$

hold for every $f \in \mathcal{C}_{e_\ell}^\alpha$, every $g \in \mathcal{C}_{p_\kappa}^\alpha$, every $t \in (0, 1]$, every $|\ell| \leq \ell_0$, and every $|\kappa| \leq \kappa_0$.

Proof. The proof is standard: one first verifies that the semigroup preserves these norms, so that the case $\beta = \alpha$ is covered. The case of integer values of β can easily be verified by an explicit calculation. The remaining values then follow by interpolation. \square

2.2 Bounds and convergence of Y^ε and Z^ε

For any integer $k \geq 2$, define the k -point correlation function $\Phi^{(k)}$ for $x, t \in \mathbf{R}^k$ by

$$\Phi^{(k)}(x, t) = \mathbf{E}(V(x_1, t_1) \dots V(x_k, t_k)).$$

(In particular, $\Phi^{(2)}(x_1, t_1, x_2, t_2) = \Phi(x_1 - x_2, t_1 - t_2)$, where Φ is the correlation function of V defined above.) With these notations at hand, we have the following bound which will prove to be useful:

Lemma 2.3 *The function $\Psi^{(4)}$ given by*

$$\Psi^{(4)}(x, t) = \Phi^{(4)}(x, t) - \Phi(x_1 - x_2, t_1 - t_2)\Phi(x_3 - x_4, t_3 - t_4),$$

satisfies the bound

$$\begin{aligned} |\Psi^{(4)}(x, t)| &\leq \eta(|x_1 - x_3| + |t_1 - t_3|)\eta(|x_2 - x_4| + |t_2 - t_4|) \\ &\quad + \eta(|x_1 - x_4| + |t_1 - t_4|)\eta(|x_2 - x_3| + |t_2 - t_3|), \end{aligned} \tag{2.5}$$

where the function $\eta: \mathbf{R}_+ \rightarrow \mathbf{R}_+$ is defined by

$$\eta(r) = \sqrt{K \varrho(r/3)}, \quad \text{with } K = 4(\|V(x, t)\|_2 \|V^3(x, t)\|_2 + \|V^2(x, t)\|_2^2),$$

where we write $\|\cdot\|_2$ for the $L^2(\Omega)$ norm of a real-valued random variable.

Remark 2.4 In the Gaussian case, one has the identity

$$\Psi^{(4)}(x, t) = \Phi(x_1 - x_3, t_1 - t_3)\Phi(x_2 - x_4, t_2 - t_4) + \Phi(x_1 - x_4, t_1 - t_4)\Phi(x_2 - x_3, t_2 - t_3),$$

so that the bound (2.5) follows from the fact that ϱ dominates the decay of the correlation function Φ .

Proof. For the sake of brevity denote $\xi_j = (x_j, t_j)$. We set

$$R_1 = \max_{1 \leq i \leq 4} \text{dist}\left(\xi_i, \bigcup_{j \neq i} \{\xi_j\}\right), \quad R_2 = \max \text{dist}\left(\{\xi_{i_1}, \xi_{i_2}\}, \{\xi_{i_3}, \xi_{i_4}\}\right),$$

where the second maximum is taken over all permutations $\{i_1, i_2, i_3, i_4\}$ of $\{1, 2, 3, 4\}$.

Consider first the case $R_1 \geq R_2$. Without loss of generality we can assume that $R_1 = \text{dist}(\xi_1, \bigcup_{j \neq 1} \{\xi_j\})$. It is easily seen that, in the case under consideration,

$$\text{dist}((\xi_i, \xi_j)) \leq 3R_1, \quad i, j = 1, 2, 3, 4. \quad (2.6)$$

Then the functions $\Phi^{(4)}$ and $\Phi(\xi_1 - \xi_2)\Phi(\xi_3 - \xi_4)$ admit the following upper bounds:

$$\begin{aligned} |\Phi^{(4)}(\xi_1, \xi_2, \xi_3, \xi_4)| &= |\mathbf{E}(V(\xi_1)V(\xi_2)V(\xi_3)V(\xi_4))| \\ &\leq \varrho(R_1)\|V(\xi_1)\|_2\|V(\xi_2)V(\xi_3)V(\xi_4)\|_2 \\ &\leq \varrho(R_1)\|V(\xi)\|_2\|(V(\xi))^3\|_2, \end{aligned}$$

and

$$\Phi(\xi_1 - \xi_2)\Phi(\xi_3 - \xi_4) \leq \varrho(R_1)\|V\|_2^2\|V\|_2^2$$

Therefore,

$$|\Psi^{(4)}(x, t)| \leq \varrho(\mathbf{R}_1)(\|V(\xi)\|_2\|(V(\xi))^3\|_2 + \|V\|_2^4)$$

From (2.6) and the fact that ϱ is a decreasing function we derive

$$K\varrho(R_1) = \eta(3R_1)\eta(3R_1) \leq \eta(|\xi_1 - \xi_3|)\eta(|\xi_2 - \xi_4|).$$

This yields the desired inequality.

Assume now that $R_1 < R_2$ and $\text{dist}(\{\xi_1, \xi_2\}, \{\xi_3, \xi_4\}) = R_2$. In this case

$$\text{dist}(\xi_1, \xi_2) < R_2 \quad \text{and} \quad \text{dist}(\xi_3, \xi_4) < R_2. \quad (2.7)$$

Indeed, if we assume that $\text{dist}(\xi_1, \xi_2) \geq R_2$, then $\text{dist}(\xi_1, \{\xi_2, \xi_3, \xi_4\}) \geq R_2$ and, thus, $R_1 \geq R_2$ which contradicts our assumption. We have

$$|\Psi^{(4)}(\xi_1, \xi_2, \xi_3, \xi_4)| = |\Phi^{(4)}(\xi_1, \xi_2, \xi_3, \xi_4) - \Phi(\xi_1 - \xi_2)\Phi(\xi_3 - \xi_4)| \quad (2.8)$$

$$\begin{aligned}
&= |\mathbf{E}([V(\xi_1)V(\xi_2) - \mathbf{E}(V(\xi_1)V(\xi_2))][V(\xi_3)V(\xi_4) - \mathbf{E}(V(\xi_3)V(\xi_4))])| \\
&\leq \varrho(R_2)\|(V(\xi))^2\|_2^2.
\end{aligned}$$

In view of (2.7), $\text{dist}(\xi_1, \xi_3) \leq 3R_2$ and $\text{dist}(\xi_2, \xi_4) \leq 3R_2$. Therefore,

$$K\varrho(R_2) \leq \eta(|\xi_1 - \xi_3|)\eta(|\xi_2 - \xi_4|),$$

and the desired inequality follows.

It remains to consider the case $R_1 < R_2$ and $\text{dist}(\{\xi_1, \xi_3\}, \{\xi_2, \xi_4\}) = R_2$; the case $\text{dist}(\{\xi_1, \xi_4\}, \{\xi_2, \xi_3\}) = R_2$ can be addressed in the same way. In this case

$$\text{dist}(\xi_1, \xi_2) \geq R_2, \quad \text{dist}(\xi_1, \xi_4) \geq R_2, \quad \text{dist}(\xi_1, \xi_3) < R_2.$$

Therefore, $\text{dist}(\xi_1, \{\xi_2, \xi_3, \xi_4\}) = \text{dist}(\xi_1, \xi_3)$, and we have

$$\begin{aligned}
|\Phi^{(4)}(\xi_1, \xi_2, \xi_3, \xi_4)| &\leq \varrho(|\xi_1 - \xi_3|)\|V(\xi)\|_2\|(V(\xi))^3\|_2 \\
|\Phi(\xi_1 - \xi_2)\Phi(\xi_3 - \xi_4)| &\leq \varrho(R_2)\|V\|_2^4 \leq \varrho(|\xi_1 - \xi_3|)\|V\|_2^4.
\end{aligned}$$

This yields

$$|\Psi^{(4)}(\xi_1, \xi_2, \xi_3, \xi_4)| \leq \varrho(|\xi_1 - \xi_3|)(\|V(\xi)\|_2\|(V(\xi))^3\|_2 + \|V\|_2^4)$$

In the same way one gets

$$|\Psi^{(4)}(\xi_1, \xi_2, \xi_3, \xi_4)| \leq \varrho(|\xi_2 - \xi_4|)(\|V(\xi)\|_2\|(V(\xi))^3\|_2 + \|V\|_2^4)$$

From the last two estimates we obtain

$$\begin{aligned}
|\Psi^{(4)}(\xi_1, \xi_2, \xi_3, \xi_4)| &\leq \sqrt{\varrho(|\xi_1 - \xi_3|)}\sqrt{\varrho(|\xi_2 - \xi_4|)}(\|V(\xi)\|_2\|(V(\xi))^3\|_2 + \|V\|_2^4) \\
&\leq \eta(|\xi_1 - \xi_3|)\eta(|\xi_2 - \xi_4|).
\end{aligned}$$

This implies the desired inequality and completes the proof of Lemma 2.3. \square

In order to prove our next result, we will need the following small lemma:

Lemma 2.5 *Let $F: \mathbf{R}_+ \rightarrow \mathbf{R}_+$ be an increasing function with $F(r) \leq r^q$. Then, $\int_0^\infty (1+r)^{-p}dF(r) < \infty$ as soon as $p > q$.*

Proof. We have $\int_0^\infty (1+r)^{-p}dF(r) \leq 1 + \int_1^\infty r^{-p}dF(r)$, so we only need to bound the latter. We write

$$\int_1^\infty r^{-p}dF(r) \leq \sum_{k \geq 0} \int_{2^k}^{2^{k+1}} r^{-p}dF(r) \leq \sum_{k \geq 0} 2^{-pk} \int_{2^k}^{2^{k+1}} dF(r) \leq \sum_{k \geq 0} 2^{-pk} 2^{q(k+1)}.$$

This expression is summable as soon as $p > q$, thus yielding the claim. \square

Lemma 2.6 Fix $t > 0$ and let $\varphi : \mathbf{R} \times \mathbf{R}_+ \rightarrow \mathbf{R}_+$ be a smooth function with compact support. Define $\varphi_\delta(x, t) = \delta^{-3} \varphi\left(\frac{x}{\delta}, \frac{t}{\delta^2}\right)$. Then, for all $p \geq 1$, $\varepsilon, \delta > 0$, one has the bound

$$\left[\mathbf{E} \left(\int_0^t \int_{\mathbf{R}} \varphi_\delta(x-y, t-s) V_\varepsilon(y, s) dy ds \right)^p \right]^{1/p} \leq C_\varphi (\varepsilon^{-1/2-\alpha/4} \wedge \delta^{-1/2} \varepsilon^{-\alpha/4} \wedge \delta^{-3/2} \varepsilon^{\alpha/4}),$$

where C_φ depends on p , on the supremum and the support of φ , and on the bound of Assumption 1.3.

Proof. We consider separately the cases $\delta > \max(\varepsilon, \varepsilon^\alpha)$, $\delta < \min(\varepsilon, \varepsilon^\alpha)$, as well as $\min(\varepsilon, \varepsilon^\alpha) \leq \delta \leq \max(\varepsilon, \varepsilon^\alpha)$.

Assume first that $\delta > \max(\varepsilon, \varepsilon^\alpha)$. Without loss of generality we also assume that p is even, that is $p = 2k$ with $k \in \mathbb{N}$. Then

$$\begin{aligned} \mathcal{J}_p^{\varepsilon, \delta} &:= \mathbf{E} \left(\int_0^t \int_{\mathbf{R}} \varphi_\delta(x-y, t-s) V_\varepsilon(y, s) dy ds \right)^p \\ &= \int_0^t \dots \int_0^t \int_{\mathbf{R}} \dots \int_{\mathbf{R}} \prod_{i=1}^{2k} \varphi_\delta(x-y_i, t-s_i) \mathbf{E} \left(\prod_{i=1}^{2k} V_\varepsilon(y_i, s_i) \right) d\vec{y} d\vec{s}, \end{aligned}$$

where $d\vec{y} = dy_1 \dots dy_{2k}$ and $d\vec{s} = ds_1 \dots ds_{2k}$. Changing the variables $\tilde{y}_i = \varepsilon^{-1} y_i$ and $\tilde{s}_i = \varepsilon^{-\alpha} s_i$, and considering the definition of φ_δ and V_ε , we obtain

$$\mathcal{J}_p^{\varepsilon, \delta} = \delta^{-6k} \varepsilon^{-k - \frac{\alpha k}{2}} \varepsilon^{2k+2\alpha k} \int_{[0, t/\varepsilon^\alpha]^{2k}} \int_{\mathbb{R}^{2k}} \prod_{i=1}^{2k} \varphi \left(\frac{x - \varepsilon \tilde{y}_i}{\delta}, \frac{t - \varepsilon^\alpha \tilde{s}_i}{\delta^2} \right) \mathbf{E} \left(\prod_{i=1}^{2k} V(\tilde{y}_i, \tilde{s}_i) \right) d\vec{y} d\vec{s}.$$

The support of the function $\prod_{i=1}^{2k} \varphi \left(\frac{x - \varepsilon \tilde{y}_i}{\delta}, \frac{t - \varepsilon^\alpha \tilde{s}_i}{\delta^2} \right)$ belongs to the rectangle $(x - k \frac{\delta}{\varepsilon} s_\varphi, x + k \frac{\delta}{\varepsilon} s_\varphi)^{2k} \times (t - k \frac{\delta^2}{\varepsilon^\alpha} s_\varphi, t + k \frac{\delta^2}{\varepsilon^\alpha} s_\varphi)^{2k}$, where s_φ is the diameter of support of $\varphi = \varphi(y, s)$. Denote $\Pi_{\delta, \varepsilon}^1 = (0, 2k \frac{\delta}{\varepsilon} s_\varphi)^{2k}$ and $\Pi_{\delta, \varepsilon}^2 = (0, 2k \frac{\delta^2}{\varepsilon^\alpha} s_\varphi)^{2k}$. Since $V(y, s)$ is stationary, we have

$$\mathcal{J}_p^{\varepsilon, \delta} \leq \delta^{-6k} \varepsilon^{-k - \frac{\alpha k}{2}} \varepsilon^{2k+2\alpha k} \|\varphi\|_C^{2k} \int_{(0, 2k \frac{\delta}{\varepsilon} s_\varphi)^{2k}} \int_{(0, 2k \frac{\delta^2}{\varepsilon^\alpha} s_\varphi)^{2k}} \left| \mathbf{E} \left(\prod_{i=1}^{2k} V(\tilde{y}_i, \tilde{s}_i) \right) \right| d\vec{y} d\vec{s}. \quad (2.9)$$

For any $R \geq 0$ we introduce a subset of \mathbb{R}^{4k}

$$\mathcal{V}_{\delta, \varepsilon}(R) = \left\{ (\vec{y}, \vec{s}) \in \Pi_{\delta, \varepsilon}^1 \times \Pi_{\delta, \varepsilon}^2 : \max_{1 \leq j \leq 2k} \text{dist}(\tilde{y}_j, \bigcup_{i \neq j} \tilde{y}_i) \leq R, \max_{1 \leq j \leq 2k} \text{dist}(\tilde{s}_j, \bigcup_{i \neq j} \tilde{s}_i) \leq R \right\},$$

and denote by $|\mathcal{V}_{\delta,\varepsilon}|(R)$ the Lebesgue measure of this set. It is easy to check that the set $\mathcal{V}_{\delta,\varepsilon}(0)$ is the union of sets of the form

$$\{(\tilde{y}, \tilde{s}) \in \Pi_{\delta,\varepsilon}^1 \times \Pi_{\delta,\varepsilon}^1 : \tilde{y}_{i_1} = \tilde{y}_{i_2}, \dots, \tilde{y}_{i_{2k-1}} = \tilde{y}_{i_{2k}}, \tilde{s}_{j_1} = \tilde{s}_{j_2}, \dots, \tilde{s}_{j_{2k-1}} = \tilde{s}_{j_{2k}}\}$$

with $i_l \neq i_m$ and $j_l \neq j_m$ if $l \neq m$, that is, $\mathcal{V}_{\delta,\varepsilon}(0)$ is the union of a finite number of subsets of $2k$ -dimensional planes in R^{4k} . The $2k$ -dimensional measure of this set satisfies the following upper bound

$$|\mathcal{V}_{\delta,\varepsilon}(0)|_{2k} \leq C(k) \left(\frac{\delta}{\varepsilon}\right)^k \left(\frac{\delta^2}{\varepsilon^\alpha}\right)^k,$$

Therefore,

$$|\mathcal{V}_{\delta,\varepsilon}|(R) \lesssim \left(\frac{\delta}{\varepsilon}\right)^k \left(\frac{\delta^2}{\varepsilon^\alpha}\right)^k R^{2k}, \quad (2.10)$$

For each $(\tilde{y}, \tilde{s}) \in \mathcal{V}_{\delta,\varepsilon}(R)$ we have

$$\left| \mathbf{E} \left(\prod_{i=1}^{2k} V(\tilde{y}_i, \tilde{s}_i) \right) \right| \leq \varrho(R) C_1(k) \|V\|_{L^2(\Omega)} \|V^{2k-1}\|_{L^2(\Omega)}. \quad (2.11)$$

Combining (2.9), (2.10) and (2.11) yields

$$\mathcal{J}_p^{\varepsilon,\delta} \lesssim \delta^{-6k} \varepsilon^{-k - \frac{\alpha k}{2}} \varepsilon^{2k+2\alpha k} \int_0^\infty \varrho(R) d|\mathcal{V}_{\delta,\varepsilon}|(R) \lesssim \delta^{-3k} \varepsilon^{\frac{\alpha k}{2}}.$$

Here, the last inequality holds due to Assumption 1.4, combined with (2.10) and Lemma 2.5. Therefore, recalling that $p = 2k$, we have the bound

$$(\mathcal{J}_p^{\varepsilon,\delta})^{1/p} \lesssim \delta^{-3/2} \varepsilon^{\alpha/4}. \quad (2.12)$$

In the case $\delta < \min(\varepsilon, \varepsilon^\alpha)$ we have

$$\begin{aligned} \mathcal{J}_p^{\varepsilon,\delta} &= \int_0^t \dots \int_0^t \int_{\mathbf{R}} \dots \int_{\mathbf{R}} \prod_{i=1}^{2k} \varphi_\delta(x - y_i, t - s_i) \mathbf{E} \left(\prod_{i=1}^{2k} V_\varepsilon(y_i, s_i) \right) d\vec{y} d\vec{s} \\ &\leq \int_0^t \dots \int_0^t \int_{\mathbf{R}} \dots \int_{\mathbf{R}} \prod_{i=1}^{2k} |\varphi_\delta(x - y_i, t - s_i)| \left| \mathbf{E} \left(\prod_{i=1}^{2k} V_\varepsilon(y_i, s_i) \right) \right| d\vec{y} d\vec{s} \\ &\leq \mathbf{E} \left((V_\varepsilon(y_1, s_1))^{2k} \right) \int_0^t \dots \int_0^t \int_{\mathbf{R}} \dots \int_{\mathbf{R}} \prod_{i=1}^{2k} |\varphi_\delta(x - y_i, t - s_i)| d\vec{y} d\vec{s} \end{aligned}$$

$$\lesssim \varepsilon^{-k-\frac{\alpha k}{2}} \|\varphi\|_{L^1}^{2k},$$

so that

$$(\mathcal{J}_p^{\varepsilon,\delta})^{1/p} \lesssim \varepsilon^{-1/2-\alpha/4}. \quad (2.13)$$

Finally, if we are in the regime $\varepsilon < \delta < \varepsilon^{\alpha/2}$, then

$$\begin{aligned} \mathcal{J}_p^{\varepsilon,\delta} &= \delta^{-6k} \varepsilon^{k+\frac{3\alpha}{2}k} \int_{[0,t/\varepsilon^\alpha]^{2k}} \int_{\mathbb{R}^{2k}} \prod_{i=1}^{2k} \varphi\left(\frac{x-\varepsilon\tilde{y}_i}{\delta}, \frac{t-\varepsilon^\alpha\tilde{s}_i}{\delta^2}\right) \mathbf{E}\left(\prod_{i=1}^{2k} V(\tilde{y}_i, \tilde{s}_i)\right) d\vec{\tilde{y}} d\vec{\tilde{s}} \\ &\leq \delta^{-6k} \varepsilon^{k+3\alpha k/2} \|\varphi\|_{L^\infty}^{2k} \int_{(0,2k\frac{\delta^2}{\varepsilon^\alpha}s_\varphi)^{2k}} \int_{(0,2k\frac{\delta}{\varepsilon}s_\varphi)^{2k}} \left| \mathbf{E}\left(\prod_{i=1}^{2k} V(\tilde{y}_i, \tilde{s}_i)\right) \right| d\vec{\tilde{y}} d\vec{\tilde{s}} \\ &\lesssim \delta^{-6k} \varepsilon^{k+3\alpha k/2} \|\varphi\|_{L^\infty}^{2k} \int_{(0,2k\frac{\delta^2}{\varepsilon^\alpha}s_\varphi)^{2k}} \|V\|_{L^2(\Omega)} \|V^{2k-1}\|_{L^2(\Omega)} \left(\frac{\delta}{\varepsilon}\right)^k \int_0^\infty \varrho(R) R^{k-1} dR \\ &\lesssim \delta^{-k} \varepsilon^{-\alpha k/2}. \end{aligned}$$

Hence,

$$(\mathcal{J}_p^{\varepsilon,\delta})^{1/p} \lesssim \delta^{-1/2} \varepsilon^{-\alpha/4} \quad (2.14)$$

so that, combining (2.12), (2.13) and (2.14), the desired estimate holds. \square

Lemma 2.7 Fix $t > 0$ and let $\varphi : \mathbf{R} \times \mathbf{R}_+ \rightarrow \mathbf{R}_+$ be a function which is uniformly bounded and decays exponentially in x , uniformly over $s \in [0, t]$.

Then, for all $p \geq 1$, $\varepsilon > 0$, one has the bound

$$\left[\mathbf{E} \left(\int_0^t \int_{\mathbf{R}} \varphi(x-y, t-s) V_\varepsilon(y, s) dy ds \right)^p \right]^{1/p} \leq C_\varphi (\varepsilon^{-1/2-\alpha/4} \wedge \varepsilon^{-\alpha/4} \wedge \varepsilon^{\alpha/4}).$$

Here, the proportionality constant depends on p , on t , on the bounds on φ , and on the bounds of Assumption 1.3.

Proof. The proof of this lemma is similar (with some simplifications) to that of the previous statement. We leave it to the reader. \square

Lemma 2.8 For each $p \geq 1$, there exists a constant C_p such that for all $\varepsilon > 0$, $t \geq 0$, $x \in \mathbf{R}$,

$$[\mathbf{E}(|Y^\varepsilon(x, t)|^p)]^{1/p} \leq C_p (1 + \sqrt{t}) \varepsilon^{\alpha/4} \quad (2.15)$$

$$[\mathbf{E}(|\partial_x Y^\varepsilon(x, t)|^p)]^{1/p} \leq C_p \quad (2.16)$$

$$[\mathbf{E}(|\partial_x^2 Y^\varepsilon(x, t)|^p)]^{1/p} \leq C_p \varepsilon^{-1}. \quad (2.17)$$

Proof. Our main ingredient is the existence of a function $\psi: \mathbf{R}_+ \rightarrow [0, 1]$ which is smooth, compactly supported in the interval $[1/2, 2]$, and such that

$$\sum_{n \in \mathbf{Z}} \psi(2^{-n}r) = 1 ,$$

for all $r > 0$.

As a consequence, we can rewrite the heat kernel as

$$p_t(x) = \sum_{n \in \mathbf{Z}} 2^{-2n} \varphi_n(x, t) , \quad (2.18)$$

where

$$\varphi_n(x, t) = 2^{3n} \varphi(2^n x, 2^{2n} t) , \quad \varphi(x, t) = p_t(x) \psi(\sqrt{x^2 + t}) . \quad (2.19)$$

The advantage of this formulation is that the function φ is smooth and compactly supported. The reason why we scale φ_n in this way, at the expense of still having a prefactor 2^{-2n} in (2.18) is that this is the scaling used in Lemma 2.6 (setting $\delta = 2^{-n}$).

We use this decomposition to define Y_n^ε by

$$Y_n^\varepsilon(x, t) = 2^{-2n} \int_0^t \int_{\mathbf{R}} \varphi_n(x - y, t - s) V_\varepsilon(y, s) dy ds , \quad (2.20)$$

so that, by (2.18), one has $Y^\varepsilon = \sum_n Y_n^\varepsilon$. Setting $\tilde{\varphi}(x, t) = \partial_x \varphi(x, t)$ and defining $\tilde{\varphi}_n(x, t) = 2^{3n} \tilde{\varphi}(2^n x, 2^{2n} t)$ as in (2.19), the derivative of Y^ε can be decomposed in the same way:

$$\partial_x Y_n^\varepsilon(x, t) = 2^{-n} \int_0^t \int_{\mathbf{R}} \tilde{\varphi}_n(x - y, t - s) V_\varepsilon(y, s) dy ds . \quad (2.21)$$

We first bound the derivative of Y^ε . Since $\tilde{\varphi}$ is smooth and compactly supported, the constants appearing in Lemma 2.6 do not depend on t and we have

$$(\mathbf{E} |\partial_x Y_n^\varepsilon(x, t)|^p)^{1/p} \lesssim 2^{n/2} \varepsilon^{\alpha/4} \wedge 2^{-n/2} \varepsilon^{-\alpha/4} = 2^{-|\frac{n}{2} + \frac{\alpha}{4} \log_2 \varepsilon|} .$$

Since the sum (over n) of this quantity is bounded independently of ε , (2.16) now follows by the triangle inequality.

Note that (2.17) follows from the same argument, if we integrate by parts (hence differentiate V_ε).

In order to finally establish (2.15), we bound Y^ε in a similar way. This time however, we combine all the terms with $n < 0$ into one single term, setting

$$p_t^-(x) = \sum_{n \leq 0} 2^{-2n} \varphi_n(x, t), \quad Y_-^\varepsilon(x, t) = \int_0^t \int_{\mathbf{R}} p_{t-s}^-(x-y) V_\varepsilon(y, s) dy ds,$$

so that $Y^\varepsilon = \sum_{n > 0} Y_n^\varepsilon + Y_-^\varepsilon$. Similarly to before, we obtain

$$(\mathbf{E}|Y_n^\varepsilon(x, t)|^p)^{1/p} \lesssim 2^{-n/2} \varepsilon^{\alpha/4}. \quad (2.22)$$

In order to bound Y_-^ε , we apply Lemma 2.6 with $\delta = 1$ and $\varphi = p^-$. It is immediate that $c_1(t) \lesssim \sqrt{t}$ and $c_3(t) \lesssim 1$, so that

$$(\mathbf{E}|Y_-^\varepsilon(x, t)|^p)^{1/p} \lesssim \sqrt{t} \varepsilon^{\alpha/4}.$$

Combining this with (2.22), summed over $n > 0$, yields the desired bound. \square

We deduce from Lemma 2.8 and equation 1.7

Corollary 2.9 *As $\varepsilon \rightarrow 0$, $Y^\varepsilon(x, t) \rightarrow 0$ in probability, locally uniformly with respect to x and t .*

Proof. It follows from Lemma 2.8 and equation 1.7 that for some $a, b > 0$ and all $p \geq 1$, all bounded subsets $D \subset \mathbf{R}_+ \times \mathbf{R}$,

$$\sup_{(x,t) \in D} \mathbf{E}[|Y^\varepsilon(x, t)|^p] \lesssim \varepsilon^{pa}, \quad (2.23)$$

$$\sup_{(x,t) \in D} \mathbf{E}[|\partial_x Y^\varepsilon(x, t)|^p] \lesssim \varepsilon^{-pb}, \quad \sup_{(x,t) \in D} \mathbf{E}[|\partial_t Y^\varepsilon(x, t)|^p] \lesssim \varepsilon^{-pb}. \quad (2.24)$$

We deduce from (2.23) that for all $(x, t), (y, s) \in D$, $p \geq 1$,

$$\mathbf{E}[|Y^\varepsilon(x, t) - Y^\varepsilon(y, s)|^p] \lesssim \varepsilon^{pa},$$

and from (2.24), writing $Y^\varepsilon(x, t) - Y^\varepsilon(y, s)$ as the sum of an integral of $\partial_x Y^\varepsilon$ and an integral of $\partial_t Y^\varepsilon$, we get

$$\mathbf{E}[|Y^\varepsilon(x, t) - Y^\varepsilon(y, s)|^p] \lesssim (|x - y| + |t - s|)^p \varepsilon^{-pb}.$$

Hence from Hölder's inequality

$$\mathbf{E}[|Y^\varepsilon(x, t) - Y^\varepsilon(y, s)|^{\alpha+\beta}] \leq (|x - y| + |t - s|)^\beta \varepsilon^{\alpha a - \beta b}.$$

Provided $\beta > 2$ and $\alpha > \beta b/a$, we obtain an estimate which allows us to deduce the result from a combination of (2.23) and Kolmogorov's Lemma. \square

We will also need

Lemma 2.10 *The function $t \rightarrow \bar{V}_\varepsilon(t)$ is continuous, and, for each $\varepsilon > 0$, there exists a positive constant \bar{V}_ε^0 such that*

$$\bar{V}_\varepsilon(t) \rightarrow \bar{V}_\varepsilon^0, \quad \text{as } t \rightarrow \infty.$$

Furthermore,

$$\lim_{\varepsilon \rightarrow 0} \bar{V}_\varepsilon^0 = \bar{V} := \int_0^\infty \int_{\mathbb{R}} \frac{\Phi(y, t)}{2\sqrt{\pi t}} dy dt,$$

and $\bar{V}_\varepsilon(t) \rightarrow \bar{V}$ as $\varepsilon \rightarrow 0$, uniformly in $t \in [1, +\infty]$.

Proof. Writing Φ_ε for the correlation function of V_ε and using the definition of $\bar{V}_\varepsilon(t)$, we have

$$\begin{aligned} \bar{V}_\varepsilon(t) &= \mathbf{E} \left[\left(\frac{\partial}{\partial x} \int_0^t \int_{\mathbb{R}} p_{t-s}(x-y) V_\varepsilon(y, s) dy ds \right)^2 \right] \\ &= \mathbf{E} \left[\left(\int_0^t \int_{\mathbb{R}} p'_{t-s}(x-y) V_\varepsilon(y, s) dy ds \right)^2 \right] \\ &= \mathbf{E} \left[\int_0^t \int_0^t \int_{\mathbb{R}} \int_{\mathbb{R}} p'_{t-s}(x-y) p'_{t-r}(x-z) V_\varepsilon(y, s) V_\varepsilon(z, r) dy dz ds dr \right] \\ &= \int_0^t \int_0^t \int_{\mathbb{R}} \int_{\mathbb{R}} p'_{t-s}(x-y) p'_{t-r}(x-z) \Phi_\varepsilon(y-z, s-r) dy dz ds dr \\ &= \int_0^t \int_0^t \int_{\mathbb{R}} \int_{\mathbb{R}} p'_s(y) p'_r(z) \Phi_\varepsilon(y-z, s-r) dy dz ds dr. \end{aligned}$$

It is easy to check that, for each $\varepsilon > 0$, this integral is a continuous function of t and that it converges, as $t \rightarrow +\infty$. Performing the change of variables $y' = \frac{y}{\varepsilon^{1/2+\alpha/4}}$, $z' = \frac{z}{\varepsilon^{1/2+\alpha/4}}$, $s' = \frac{s}{\varepsilon^{1+\alpha/2}}$, $r' = \frac{r}{\varepsilon^{1+\alpha/2}}$, renaming the new variables and setting $T_\varepsilon = \varepsilon^{-1-\alpha/2}t$, we obtain

$$\bar{V}_\varepsilon(t) = \frac{1}{16\pi} \int_0^{T_\varepsilon} \int_0^{T_\varepsilon} \int_{\mathbb{R}} \int_{\mathbb{R}} \frac{y}{s^{3/2}} \frac{z}{r^{3/2}} e^{-\frac{y^2}{4s} - \frac{z^2}{4r}} \Phi\left(\frac{y-z}{\varepsilon^{\frac{1-\alpha}{2}}}, \frac{s-r}{\varepsilon^{\frac{\alpha}{2}-1}}\right) dy dz ds dr.$$

We represent the integral on the right-hand side as

$$\bar{V}_\varepsilon(t) = \frac{1}{16\pi} \int_0^{T_\varepsilon} \int_0^{T_\varepsilon} \int_{\mathbb{R}} \int_{\mathbb{R}} \frac{z^2}{s^{3/2} r^{3/2}} e^{-\frac{z^2}{4s} - \frac{z^2}{4r}} \Phi\left(\frac{y-z}{\varepsilon^{\frac{1}{2}-\frac{\alpha}{4}}}, \frac{s-r}{\varepsilon^{\frac{\alpha}{2}-1}}\right) dy dz ds dr + r_\varepsilon(t). \quad (2.25)$$

The further analysis relies on the following limit relation:

$$\lim_{\varepsilon \rightarrow 0} \sup_{0 < t \leq +\infty} |r_\varepsilon(t)| = 0. \quad (2.26)$$

In order to justify it we denote $\varkappa = \frac{1}{2} - \frac{\alpha}{4}$ and $\varkappa_1 = \frac{\varkappa}{10}$, and divide the integration area into four parts as follows

$$\begin{aligned} \Pi_1 &= \{(y, z, s, r) \in \mathbb{R}^2 \times (\mathbb{R}^+)^2 : s \leq \varepsilon^{\varkappa_1}, r \leq \varepsilon^{\varkappa_1}\}, \\ \Pi_2 &= \{(y, z, s, r) \in \mathbb{R}^2 \times (\mathbb{R}^+)^2 : \varepsilon^{\varkappa_1} < s \leq T_\varepsilon, r \leq \varepsilon^{\varkappa_1}\}, \\ \Pi_3 &= \{(y, z, s, r) \in \mathbb{R}^2 \times (\mathbb{R}^+)^2 : s \leq \varepsilon^{\varkappa_1}, \varepsilon^{\varkappa_1} < r \leq T_\varepsilon\}, \\ \Pi_4 &= \{(y, z, s, r) \in \mathbb{R}^2 \times (\mathbb{R}^+)^2 : \varepsilon^{\varkappa_1} < s \leq T_\varepsilon, \varepsilon^{\varkappa_1} < r \leq T_\varepsilon\}. \end{aligned}$$

In Π_1 we have

$$\int_{\Pi_1} \frac{|y||z|}{s^{\frac{3}{2}} r^{\frac{3}{2}}} e^{-\frac{y^2}{4s} - \frac{z^2}{4r}} \Phi\left(\frac{y-z}{\varepsilon^{\frac{1}{2}-\frac{\alpha}{4}}}, \frac{s-r}{\varepsilon^{\frac{\alpha}{2}-1}}\right) dx dy ds dr \leq C^2 \int_0^{\varepsilon^{\varkappa_1}} \int_0^{\varepsilon^{\varkappa_1}} \frac{ds dr}{s^{\frac{1}{2}} r^{\frac{1}{2}}} = 4C^2 \varepsilon^{\varkappa_1}. \quad (2.27)$$

To estimate the integral over Π_2 we first notice that there exists a constant C_1 such that

$$\frac{|y|}{s^{\frac{1}{2}}} e^{-\frac{y^2}{4s}} \leq C_1$$

uniformly over $s > 0$ and $y \in \mathbf{R}$. Then,

$$\begin{aligned}
& \int_{\Pi_2} \frac{|y||z|}{s^{\frac{3}{2}}r^{\frac{3}{2}}} e^{-\frac{y^2}{4s}-\frac{z^2}{4r}} \Phi\left(\frac{y-z}{\varepsilon^{\frac{1}{2}-\frac{\alpha}{4}}}, \frac{s-r}{\varepsilon^{\frac{\alpha}{2}-1}}\right) dy dz ds dr \\
& \leq C_1 \int_{\varepsilon^{\kappa_1}}^{T_\varepsilon} \int_0^{\varepsilon^{\kappa_1}} \int_{\mathbb{R}} \frac{|z| dz dr ds}{s r^{\frac{3}{2}}} e^{-\frac{z^2}{r}} \int_{\mathbb{R}} \Phi\left(\frac{y-z}{\varepsilon^{\frac{1}{2}-\frac{\alpha}{4}}}, \frac{s-r}{\varepsilon^{\frac{\alpha}{2}-1}}\right) dy \\
& = C_1 \varepsilon^\varkappa \int_{\varepsilon^{\kappa_1}}^{T_\varepsilon} \int_0^{\varepsilon^{\kappa_1}} \int_{\mathbb{R}} e^{-\frac{z^2}{r}} \bar{\Phi}\left(\frac{s-r}{\varepsilon^{\frac{\alpha}{2}-1}}\right) \frac{|z| dz dr ds}{s r^{\frac{3}{2}}} \\
& = CC_1 \varepsilon^\varkappa \int_{\varepsilon^{\kappa_1}}^{T_\varepsilon} \int_0^{\varepsilon^{\kappa_1}} \bar{\Phi}\left(\frac{s-r}{\varepsilon^{\frac{\alpha}{2}-1}}\right) \frac{dr ds}{s r^{\frac{1}{2}}} \leq CC_1 \varepsilon^\varkappa \int_{\varepsilon^{\kappa_1}}^{T_\varepsilon} \int_0^{\varepsilon^{\kappa_1}} \widehat{\Phi}\left(\frac{s}{\varepsilon^{\frac{\alpha}{2}-1}}\right) \frac{dr ds}{s r^{\frac{1}{2}}} \\
& = 2CC_1 \varepsilon^\varkappa \varepsilon^{\frac{\kappa_1}{2}} \int_{\varepsilon^{\kappa_1}}^{T_\varepsilon} \widehat{\Phi}\left(\frac{s}{\varepsilon^{\frac{\alpha}{2}-1}}\right) \frac{ds}{s} \leq 2CC_1 \varepsilon^\varkappa \varepsilon^{\frac{\kappa_1}{2}} \int_{\varepsilon^{\kappa_1+2\kappa}}^{\infty} \widehat{\Phi}(s) \frac{ds}{s} \\
& \leq C_2(\varkappa_1 + 2\kappa) \varepsilon^{\varkappa + \frac{\kappa_1}{2}} |\log \varepsilon|;
\end{aligned} \tag{2.28}$$

here $\bar{\Phi}(t) = \int_{\mathbb{R}} \Phi(x, t) dx$, and $\widehat{\Phi}(t)$ stands for $\max\{\bar{\Phi}(s) : t-1 \leq s \leq t\}$. A similar estimate holds true for the integral over Π_3 . Therefore,

$$\lim_{\varepsilon \rightarrow 0} \sup_{0 < t \leq +\infty} \left| \int_{\Pi_1 \cup \Pi_2 \cup \Pi_3} \frac{y}{s^{3/2}} \frac{z}{r^{3/2}} e^{-\frac{y^2}{4s}-\frac{z^2}{4r}} \Phi\left(\frac{y-z}{\varepsilon^{\frac{1}{2}-\frac{\alpha}{4}}}, \frac{s-r}{\varepsilon^{\frac{\alpha}{2}-1}}\right) dy dz ds dr \right| = 0. \tag{2.29}$$

We also have

$$\begin{aligned}
& \int_{\Pi_1 \cup \Pi_2} \frac{z^2}{s^{3/2}r^{3/2}} e^{-\frac{z^2}{4s}-\frac{z^2}{4r}} \Phi\left(\frac{y-z}{\varepsilon^{\frac{1}{2}-\frac{\alpha}{4}}}, \frac{s-r}{\varepsilon^{\frac{\alpha}{2}-1}}\right) dy dz ds dr \\
& = C \varepsilon^\varkappa \int_0^{T_\varepsilon} \int_0^{\varepsilon^{\kappa_1}} \bar{\Phi}\left(\frac{s-r}{\varepsilon^{-2\kappa}}\right) \frac{ds dr}{(s+r)^{\frac{3}{2}}} = C \int_0^{\varepsilon^{2\kappa} T_\varepsilon} \int_0^{\varepsilon^{\kappa_1+2\kappa}} \bar{\Phi}(s-r) \frac{ds dr}{(s+r)^{\frac{3}{2}}} \\
& \leq C \int_0^1 \int_0^{\varepsilon^{\kappa_1+2\kappa}} \bar{\Phi}(s-r) \frac{ds dr}{(s+r)^{\frac{3}{2}}} + C \int_1^\infty \int_0^{\varepsilon^{\kappa_1+2\kappa}} \bar{\Phi}(s-r) \frac{ds dr}{(s+r)^{\frac{3}{2}}} \leq C_4 \varepsilon^\varkappa
\end{aligned} \tag{2.30}$$

Combining this estimate with a similar estimate for the integral over $\Pi_1 \cup \Pi_3$, we obtain

$$\lim_{\varepsilon \rightarrow 0} \sup_{0 < t \leq +\infty} \left| \int_{\Pi_1 \cup \Pi_2 \cup \Pi_3} \frac{z^2}{s^{3/2} r^{3/2}} e^{-\frac{z^2}{4s} - \frac{z^2}{4r}} \Phi\left(\frac{y-z}{\varepsilon^{\frac{1}{2}-\frac{\alpha}{4}}, \frac{s-r}{\varepsilon^{\frac{\alpha}{2}-1}}}\right) dy dz ds dr \right| = 0. \quad (2.31)$$

In order to justify (2.26) it remains to show that

$$\lim_{\varepsilon \rightarrow 0} \sup_{0 < t \leq +\infty} \int_{\Pi_4} \left| \frac{yz e^{-\frac{y^2}{4s} - \frac{z^2}{4r}} - z^2 e^{-\frac{z^2}{4s} - \frac{z^2}{4r}}}{s^{3/2} r^{3/2}} \Phi\left(\frac{y-z}{\varepsilon^{\frac{1}{2}-\frac{\alpha}{4}}, \frac{s-r}{\varepsilon^{\frac{\alpha}{2}-1}}}\right) dy dz ds dr \right| = 0 \quad (2.32)$$

We first estimate

$$\begin{aligned} J_\varepsilon(t) &:= \int_{\Pi_4} \frac{|yz|}{s^{3/2} r^{3/2}} e^{-\frac{z^2}{4r}} \left| e^{-\frac{y^2}{4s}} - e^{-\frac{z^2}{4s}} \right| \Phi\left(\frac{y-z}{\varepsilon^\alpha}, \frac{s-r}{\varepsilon^{-2\alpha}}\right) dy dz ds dr \\ &\leq \frac{1}{4} \int_{\Pi_4} \frac{|yz| |z^2 - y^2|}{s^{5/2} r^{3/2}} e^{-\frac{z^2}{4r}} \left(e^{-\frac{y^2}{4s}} + e^{-\frac{z^2}{4s}} \right) \Phi\left(\frac{y-z}{\varepsilon^\alpha}, \frac{s-r}{\varepsilon^{-2\alpha}}\right) dy dz ds dr \\ &\lesssim \varepsilon^\alpha \int_{\Pi_4} \left(\frac{|y|^3 + |y-z|^3}{s^{5/2} r^{3/2}} e^{-\frac{y^2}{4s}} + \frac{|z|^3 + |y-z|^3}{s^{5/2} r^{3/2}} e^{-\frac{z^2}{4s}} \right) e^{-\frac{z^2}{4r}} \Phi_1\left(\frac{y-z}{\varepsilon^\alpha}, \frac{s-r}{\varepsilon^{-2\alpha}}\right) dy dz ds dr \end{aligned} \quad (2.33)$$

with $\Phi_1(x, t) = |x| \Phi(x, t)$; here we have used the inequality $|e^a - e^b| \leq |b - a|(e^a + e^b)$ and the estimates $|yz||y + z| \leq C(|y|^3 + |y - z|^3)$ and $|yz||y + z| \leq C(|z|^3 + |y - z|^3)$ that follow from the Young inequality. Let us estimate the integral

$$\begin{aligned} &\varepsilon^\alpha \int_{\Pi_4} \frac{|y|^3}{s^{5/2} r^{3/2}} e^{-\frac{y^2}{4s}} e^{-\frac{z^2}{4r}} \Phi_1\left(\frac{y-z}{\varepsilon^\alpha}, \frac{s-r}{\varepsilon^{-2\alpha}}\right) dy dz ds dr \\ &\leq C_3 \varepsilon^\alpha \int_{\Pi_4} \frac{1}{sr^{3/2}} e^{-\frac{z^2}{4r}} \Phi_1\left(\frac{y-z}{\varepsilon^\alpha}, \frac{s-r}{\varepsilon^{-2\alpha}}\right) dy dz ds dr \leq C_4 \varepsilon^{2\alpha} \int_{\varepsilon^{\alpha_1}}^{\infty} \int_{\varepsilon^{\alpha_1}}^{\infty} \frac{1}{sr} \bar{\Phi}_1\left(\frac{s-r}{\varepsilon^{-2\alpha}}\right) ds dr \\ &= C_4 \varepsilon^{2\alpha} \int_{\varepsilon^{\alpha_1+2\alpha}}^{\infty} \int_{\varepsilon^{\alpha_1+2\alpha}}^{\infty} \bar{\Phi}_1(s-r) \frac{ds dr}{sr} \leq C_5 \varepsilon^{2\alpha} (\log \varepsilon)^2; \end{aligned}$$

here $C_3 = \max(x^3 e^{-x^2})$, and $\bar{\Phi}_1(t)$ stands for $\int_{\mathbb{R}} \Phi_1(x, t) dx$. Other terms on the right-hand side of (2.33) can be estimated in a similar way. Thus we obtain

$$\lim_{\varepsilon \rightarrow 0} \sup_{0 < t \leq \infty} J_\varepsilon(t) = 0. \quad (2.34)$$

The inequality

$$\int_{\Pi_4} \frac{|yz - z^2|}{s^{3/2}r^{3/2}} e^{-\frac{z^2}{4r}} e^{-\frac{z^2}{4s}} \Phi\left(\frac{y-z}{\varepsilon^\kappa}, \frac{s-r}{\varepsilon^{-2\kappa}}\right) dy dz ds dr \leq C\varepsilon^\kappa (\log \varepsilon)^2$$

can be obtained in the same way with a number of simplifications. This yields (2.26).

It remains to notice that

$$\begin{aligned} & \int_{\varepsilon^{\kappa_1}\varepsilon^{\kappa_1}}^{T_\varepsilon} \int_{\mathbb{R}} \int_{\mathbb{R}} \int_{\mathbb{R}} \frac{z^2}{s^{3/2}r^{3/2}} e^{-\frac{z^2}{4s} - \frac{z^2}{4r}} \Phi\left(\frac{y-z}{\varepsilon^\kappa}, \frac{s-r}{\varepsilon^{-2\kappa}}\right) dy dz ds dr \\ &= \varepsilon^\kappa \int_{\varepsilon^{\kappa_1}\varepsilon^{\kappa_1}}^{T_\varepsilon} \int_{\mathbb{R}} \int_{\mathbb{R}} \frac{z^2}{s^{3/2}r^{3/2}} e^{-\frac{z^2}{4s} - \frac{z^2}{4r}} \bar{\Phi}\left(\frac{s-r}{\varepsilon^{-2\kappa}}\right) dz ds dr \\ &= C_0 \varepsilon^\kappa \int_{\varepsilon^{\kappa_1}\varepsilon^{\kappa_1}}^{T_\varepsilon} \int_{\mathbb{R}} \bar{\Phi}\left(\frac{s-r}{\varepsilon^{-2\kappa}}\right) \frac{ds dr}{(s+r)^{3/2}} = C_0 \int_{\varepsilon^{\kappa_1+2\kappa}}^{\varepsilon^{-\alpha t}} \int_{\varepsilon^{\kappa_1+2\kappa}}^{\varepsilon^{-\alpha t}} \bar{\Phi}(s-r) \frac{ds dr}{(s+r)^{3/2}} \\ &= C_0 \int_0^\infty \int_0^\infty \bar{\Phi}(s-r) \frac{ds dr}{(s+r)^{3/2}} + R_\varepsilon(t) \end{aligned}$$

with $C_0 = \int_{\mathbb{R}} z^2 e^{-z^2/4} dz$, and

$$\lim_{\varepsilon \rightarrow 0} \sup_{1 \leq t \leq +\infty} |R_\varepsilon(t)| = 0.$$

Combining the last two relations with (2.25) and (2.26), we obtain the desired statement. \square

Lemma 2.11 *For any $T > 0$, any even integer $k \geq 2$, any $0 < \beta < 1/k$, any $p > k$ and any $\kappa > 0$, there exists a constant C such that for all $0 \leq t \leq T$, $\varepsilon > 0$,*

$$\begin{aligned} (\mathbf{E} \|Y^\varepsilon(t)\|_{0,p_\kappa}^p)^{1/p} &\leq C \varepsilon^{\frac{\alpha}{4}(1-\kappa)}, \quad (\mathbf{E} \|\partial_x Y^\varepsilon(t)\|_{0,p_\kappa}^p)^{1/p} \leq C \varepsilon^{-\kappa}, \\ (\mathbf{E} \|\partial_x Y^\varepsilon(t)\|_{\beta,p_\kappa}^p)^{1/p} &\leq C \varepsilon^{-\kappa}. \end{aligned}$$

Proof. We establish the estimates of the norms of $\partial_x Y^\varepsilon(t)$ only. The norm of $Y^\varepsilon(t)$ is estimated similarly. Let $q > 1$ and $p = qk$. For any $x < y$, we have the identity

$$|\partial_x Y^\varepsilon(t, y) - \partial_x Y^\varepsilon(t, x)|^k = k \int_x^y (\partial_x Y^\varepsilon(t, z) - \partial_x Y^\varepsilon(t, x))^{k-1} \partial_x^2 Y^\varepsilon(t, z) dz.$$

Raising this to the power q and taking expectations, we obtain

$$\begin{aligned} \mathbf{E}(|\partial_x Y^\varepsilon(t, y) - \partial_x Y^\varepsilon(t, x)|^p) &\leq k^q \left| \int_x^y (\partial_x Y^\varepsilon(t, z) - \partial_x Y^\varepsilon(t, x))^{k-1} \partial_x^2 Y^\varepsilon(t, z) dz \right|^q \\ &\lesssim (y-x)^{q-1} \int_x^y \mathbf{E}(|(\partial_x Y^\varepsilon(t, z) - \partial_x Y^\varepsilon(t, x))^{k-1} \partial_x^2 Y^\varepsilon(t, z)|^q) dz \\ &\lesssim (y-x)^q \sqrt{\mathbf{E}(|\partial_x Y^\varepsilon(t, x)|^{2q(k-1)}) \mathbf{E}(|\partial_x^2 Y^\varepsilon(t, x)|^{2q})} \lesssim (y-x)^q \varepsilon^{-q}, \end{aligned} \quad (2.35)$$

where we have used the stationarity (in z) of the processes $\partial_x Y^\varepsilon(t, z)$ and $\partial_x^2 Y^\varepsilon(t, z)$, as well as the estimates (2.16) and (2.17) from Lemma 2.8.

As a consequence of Kolmogorov's Lemma, there exists a stationary sequence of positive random variables $\{\xi_n\}_{n \in \mathbf{Z}}$ such that for every $n \in \mathbf{Z}$, the bound

$$\sup_{x \in [n, n+1]} |\partial_x Y^\varepsilon(t, x)| \leq \xi_n,$$

holds almost surely, and such that $(\mathbf{E} \xi_n^p)^{1/p} \lesssim \varepsilon^{-1/k}$ for every $p \geq 1$. The bound on $\|\partial_x Y^\varepsilon(t)\|_{0, p_\kappa}$ then follows at once.

The bound on $\|\partial_x Y^\varepsilon(t)\|_{\beta, p_\kappa}$ follows in virtually the same way, using the fact that (2.35) also yields the bound

$$\sup_{x, y \in [n-1, n+1]} \frac{|\partial_x Y^\varepsilon(t, x) - \partial_x Y^\varepsilon(t, y)|}{|x - y|^\beta} \leq \tilde{\xi}_n,$$

for some stationary sequence of random variables $\tilde{\xi}_n$ which has all of its moments bounded in the same way as the sequence $\{\xi_n\}$. \square

We further obtain the following bound on the ‘‘negative Hölder norm’’ of $\partial_x Y^\varepsilon$:

Corollary 2.12 *For any $T > 0$, k being any even integer, $p > k$ and $\kappa = 1/k$, there exists a constant $C_{T, p, \kappa}$ such that*

$$\left(\mathbf{E} \|\partial_x Y^\varepsilon(t)\|_{-\frac{1}{4}, p_\kappa}^p \right)^{1/p} \leq C_{T, p, \kappa} \varepsilon^{\alpha/16 - \kappa},$$

for all $0 \leq t \leq T$, $\varepsilon > 0$. \square

Proof. We note that

$$\|\partial_x Y^\varepsilon(t)\|_{-\frac{1}{4}, p_\kappa} = \sup_{|x-y| \leq 1} \frac{|Y^\varepsilon(t, x) - Y^\varepsilon(t, y)|}{p_\kappa(x) |x - y|^{3/4}},$$

$$\|Y^\varepsilon(t)\|_{0,p_\kappa} = \sup_x \frac{|Y^\varepsilon(t,x)|}{p_\kappa(x)}, \quad \|\partial_x Y^\varepsilon(t)\|_{0,p_\kappa} = \sup_x \frac{|\partial_x Y^\varepsilon(t,x)|}{p_\kappa(x)}.$$

We have, for $|x-y| \leq 1$,

$$\begin{aligned} \frac{|Y^\varepsilon(t,x) - Y^\varepsilon(t,y)|}{p_\kappa(x)|x-y|^{3/4}} &= \left(\frac{|Y^\varepsilon(t,x) - Y^\varepsilon(t,y)|}{p_\kappa(x)} \right)^{1/4} \left(\frac{|Y^\varepsilon(t,x) - Y^\varepsilon(t,y)|}{p_\kappa(x)|x-y|} \right)^{3/4} \\ &\leq \left(\frac{|Y^\varepsilon(t,x)|}{p_\kappa(x)} + C_\kappa \frac{|Y^\varepsilon(t,y)|}{p_\kappa(y)} \right)^{1/4} \left(C_\kappa \sup_{x \leq z \leq y} \frac{|\partial_x Y^\varepsilon(t,z)|}{p_\kappa(z)} \right)^{3/4}. \end{aligned}$$

It remains to take supremums and apply Hölder's inequality. \square

We have similar results for $Z^\varepsilon(x,t)$.

Lemma 2.13 *For each $p \geq 1$, there exists a constant C such that for all $\varepsilon > 0$, $t \geq 0$, $x \in \mathbf{R}$,*

$$[\mathbf{E}(|Z^\varepsilon(x,t)|^2)]^{1/2} \leq C(1+t^2)\varepsilon^{\gamma\alpha/2}.$$

Proof. The main ingredient in the proof is a bound on the correlation function of the right hand side of the equation for Z^ε , which we denote by

$$\Lambda_\varepsilon(z, z') = \mathbf{Cov}(|\partial_x Y^\varepsilon(z)|^2, |\partial_x Y^\varepsilon(z')|^2).$$

Inserting the definition of Y^ε , we obtain the identity

$$\Lambda_\varepsilon(z, z') = \int \cdots \int \tilde{P}(z-z_1)\tilde{P}(z-z_2)\tilde{P}(z'-z_3)\tilde{P}(z'-z_4)\Psi_\varepsilon^{(4)}(z_1, \dots, z_4) dz_1 \cdots dz_4,$$

where

$$\tilde{P}(z) = \tilde{P}(x,t) = \partial_x p_t(x),$$

with p_t the standard heat kernel and

$$\Psi_\varepsilon^{(4)}(z_1, \dots, z_4) = \varepsilon^{-2-\alpha}\Psi^{(4)}\left(\frac{x_1}{\varepsilon}, \dots, \frac{x_4}{\varepsilon}, \frac{t_1}{\varepsilon^\alpha}, \dots, \frac{t_4}{\varepsilon^\alpha}\right).$$

Here, we used the shorthand notation $z_i = (x_i, t_i)$, and integrals over z_i are understood to be shorthand for $\int_0^t \int_{\mathbf{R}} dx_i dt_i$. We now make use of Lemma 2.3, which allows to factor this integral as

$$|\Lambda_\varepsilon(z, z')| \lesssim \left(\varepsilon^{-1-\frac{\alpha}{2}} \int \int \tilde{P}(z-z_1)\tilde{P}(z'-z_3)\varrho_\varepsilon(z_1-z_3) dz_1 dz_3 \right)^2 \stackrel{\text{def}}{=} \tilde{\varrho}_\varepsilon^2(z, z'),$$

where we used the shorthand notation

$$\varrho_\varepsilon(x, t) = \varrho\left(\frac{x}{\varepsilon}, \frac{t}{\varepsilon^\alpha}\right).$$

We will show below that the following bound holds:

Lemma 2.14

$$\tilde{\varrho}_\varepsilon(z, z') \lesssim \left(1 \wedge \frac{\varepsilon^{\alpha\gamma/2}}{d_p^\gamma(z, z')}\right) + (1+t+t')\varepsilon^{\alpha/2} \stackrel{\text{def}}{=} \zeta_\varepsilon(z-z') + (1+t+t')\varepsilon^{\alpha/2},$$

where d_p denotes the parabolic distance given by

$$d_p(z, z')^2 = |x - x'|^2 + |t - t'|.$$

Taking this bound for granted, we write as in the proof of Lemma 2.8 $Z^\varepsilon = Z_-^\varepsilon + \sum_{n>0} Z_n^\varepsilon$ with

$$Z_n^\varepsilon(z) = 2^{-2n} \int \varphi_n(z-z') (|\partial_x Y^\varepsilon(z')|^2 - \bar{V}_\varepsilon(t')) dz',$$

and similarly for Z_-^ε . Squaring this expression and inserting the bound from Lemma 2.14, we obtain

$$\begin{aligned} \mathbf{E}|Z_n^\varepsilon(z)|^2 &\lesssim 2^{-4n} \int \int \varphi_n(z-z') \varphi_n(z-z'') (\zeta_\varepsilon^2(z'-z'') + (1+t'+t'')^2 \varepsilon^\alpha) dz' dz'' \\ &\lesssim 2^{-n} \int \zeta_\varepsilon^2(z') dz' + 2^{-4n} (1+t)^4 \varepsilon^\alpha, \end{aligned}$$

where we made use of the scaling of φ_n given by (2.19). Performing the corresponding bound for Z_-^ε , we similarly obtain

$$\mathbf{E}|Z_-^\varepsilon(z)|^2 \lesssim t \int \zeta_\varepsilon^2(z') dz' + (1+t)^4 \varepsilon^\alpha.$$

The claim now follows from the bounds

$$\int \zeta_\varepsilon^2(z') dz' \leq \int_0^t \int_{\mathbf{R}} \frac{\varepsilon^{\alpha\gamma}}{(|x|^2 + |s|)^\gamma} dx ds \lesssim \varepsilon^{\alpha\gamma} t^{\frac{3}{2}-\gamma}.$$

□

Proof of Lemma 2.14. Similarly to the proof of Lemma 2.8, we write

$$\tilde{\varrho}_\varepsilon(z, z') = \sum_{n_1 \geq 0} \sum_{n_2 \geq 0} \tilde{\varrho}_\varepsilon^{n_1, n_2}(z, z'),$$

with

$$\tilde{\varrho}_\varepsilon^{n_1, n_2}(z, z') = \varepsilon^{-1-\frac{\alpha}{2}} 2^{-n_1-n_2} \int \int \tilde{\varphi}_{n_1}(z-z_1) \tilde{\varphi}_{n_2}(z'-z_2) \varrho_\varepsilon(z_1-z_2) dz_1 dz_2.$$

Here, for $n \geq 1$, $\tilde{\varphi}_n$ is defined as in the proof of Lemma 2.8, whereas $\tilde{\varphi}_0$ is different from what it was there and is defined as

$$\tilde{\varphi}_0(x, t) = \partial_x p_t^-(x) .$$

By symmetry, we can restrict ourselves to the case $n_1 \geq n_2$, which we will do in the sequel. In the case where $n_2 > 0$, the above integral could be restricted to the set of pairs (z_1, z_2) , such that their parabolic distance satisfies

$$d_p(z_1, z_2) \geq (d_p(z, z') - 2^{2-n_2})_+ ,$$

where $(\dots)_+$ denotes the positive part of a number.

Replacing $\tilde{\varphi}_{n_2}$ by its supremum and integrating out $\tilde{\varphi}_{n_1}$ and ϱ_ε yields the bound

$$\tilde{\varrho}_\varepsilon^{n_1, n_2}(z, z') \lesssim (1 + \delta_{n_2, 0}(t + t')) 2^{2n_2 - n_1} \varepsilon^{\alpha/2} \int_{A_\varepsilon(n_2)} \varrho(z_3) dz_3 ,$$

where $A_\varepsilon(0) = \mathbf{R}^2$ and

$$A_\varepsilon(n_2) = \{z_3 : d_p(0, z_3) \geq \varepsilon^{-\alpha/2} (d_p(z, z') - 2^{2-n_2})_+\} ,$$

for $n_2 > 0$. (Remark that the prefactor $1 + t + t'$ is relevant only in the case $n_1 = n_2 = 0$.) It follows from the integrability of ϱ that one always has the bound

$$\tilde{\varrho}_\varepsilon^{n_1, n_2}(z, z') \lesssim (1 + \delta_{n_2, 0}(t + t')) 2^{2n_2 - n_1} \varepsilon^{\alpha/2} . \quad (2.36)$$

Moreover, we deduce from Assumption 1.4 that, whenever $n_2 > 0$ and $d(z, z') \geq 2^{3-n_2}$, one has the improved bound: for any $\gamma > 0$,

$$\tilde{\varrho}_\varepsilon^{n_1, n_2}(z, z') \lesssim 2^{2n_2 - n_1} \varepsilon^{\alpha/2} \left(1 \wedge \frac{\varepsilon^{\alpha\gamma/2}}{d_p^\gamma(z, z')} \right) . \quad (2.37)$$

The bound (2.36) is sufficient for our needs in the case $n_2 = 0$, so we assume that $n_2 > 0$ from now on.

We now obtain a second bound on $\tilde{\varrho}_\varepsilon^{n_1, n_2}(z, z')$ which will be useful in the regime where n_2 is very large. Since the integral of $\tilde{\varphi}_{n_1}$ is bounded independently of n_1 , we obtain

$$\tilde{\varrho}_\varepsilon^{n_1, n_2}(z, z') \lesssim \varepsilon^{-1 - \frac{\alpha}{2}} 2^{-n_1 - n_2} \sup_{d_p(z_1, z) \leq 2^{1-n_1}} \int \tilde{\varphi}_{n_2}(z' - z_2) \varrho_\varepsilon(z_1 - z_2) dz_2 . \quad (2.38)$$

We now distinguish between three cases, which depend on the size of $z - z'$.

Case 1: $d_p(z, z') \leq \varepsilon^{\alpha/2}$. In this case, we proceed as in the proof of Lemma 2.6, which yields

$$\begin{aligned} \tilde{\varrho}_\varepsilon^{n_1, n_2}(z, z') &\lesssim \varepsilon^{-1-\frac{\alpha}{2}} 2^{-n_1-n_2} \sup_{z_1} \int \tilde{\varphi}_{n_2}(z_2) \varrho_\varepsilon(z_2 - z_1) dz_2 \\ &\lesssim \varepsilon^{-1-\frac{\alpha}{2}} 2^{-n_1-n_2} \sup_{x_1} \int_{\mathbf{R}} \sup_s \varrho_\varepsilon(x_2 - x_1, s) \int_0^t \tilde{\varphi}_{n_2}(x_2, t_2) dt_2 dx_2 \\ &\lesssim \varepsilon^{-1-\frac{\alpha}{2}} 2^{-n_1} \int_{\mathbf{R}} \sup_s \varrho_\varepsilon(x_2, s) dx_2 \lesssim \varepsilon^{-\frac{\alpha}{2}} 2^{-n_1}. \end{aligned} \quad (2.39)$$

Case 2: $|x - x'| \geq d_p(z, z')/2 \geq \varepsilon^{\alpha/2}/2$. Note that in (2.38), the argument of ϱ_ε can only ever take values with $|x_1 - x_2| \in B_\varepsilon(n_2)$ where

$$B_\varepsilon(n_2) = \{\bar{x} : |\bar{x}| \geq (|x - x'| - 2^{2-n_2})\}.$$

As a consequence, we obtain the bound

$$\tilde{\varrho}_\varepsilon^{n_1, n_2}(z, z') \lesssim \varepsilon^{-1-\frac{\alpha}{2}} 2^{-n_1-n_2} \sup_{\bar{x} \in B_\varepsilon(n_2)} \sup_{s \in \mathbf{R}} \varrho_\varepsilon(\bar{x}, s).$$

The case of interest to us for this bound will be $2^{6-n_2} \leq \varepsilon^{\alpha/2}$, in which case we deduce from this calculation and Assumption 1.4 that

$$\tilde{\varrho}_\varepsilon^{n_1, n_2}(z, z') \lesssim \varepsilon^{-1-\frac{\alpha}{2}} 2^{-n_1-n_2} \left(\frac{\varepsilon}{d_p(z, z')} \right)^\gamma,$$

where γ is an arbitrarily large exponent. Choosing γ large enough, we conclude that one also has the bound

$$\tilde{\varrho}_\varepsilon^{n_1, n_2}(z, z') \lesssim \varepsilon^{-\frac{\alpha}{2}} 2^{-n_1} \left(1 \wedge \frac{\varepsilon^{\alpha/2}}{d_p(z, z')} \right)^\gamma, \quad (2.40)$$

which will be sufficient for our needs.

Case 3: $|t - t'| \geq d_p^2(z, z')/2 \geq \varepsilon^\alpha/2$. Similarly, we obtain

$$\tilde{\varrho}_\varepsilon^{n_1, n_2}(z, z') \lesssim \varepsilon^{-\frac{\alpha}{2}} 2^{-n_1} \int_{\mathbf{R}} \sup_{s \in B'_\varepsilon(n_2)} \varrho_\varepsilon(x_2, s) dx_2,$$

where

$$B'_\varepsilon(n_2) = \{s : |s| \geq \varepsilon^{-\alpha} (|t - t'| - 2^{8-2n_2})\}.$$

Restricting ourselves again to the case $2^{6-n_2} \leq \varepsilon^{\alpha/2}$, this yields as before

$$\tilde{\varrho}_\varepsilon^{n_1, n_2}(z, z') \lesssim \varepsilon^{-\frac{\alpha}{2}} 2^{-n_1} \left(1 \wedge \frac{\varepsilon^{\alpha/2}}{d_p(z, z')} \right)^\gamma. \quad (2.41)$$

It now remains to sum over all values $n_1 \geq n_2 \geq 0$.

For $n_2 = 0$, we sum the bound (2.36), which yields

$$\sum_{n_1 \geq 0} \tilde{\varrho}_\varepsilon^{n_1, 0}(z, z') \leq (1 + t + t') \varepsilon^{\alpha/2}.$$

In order to sum the remaining terms, we first consider the case $d_p(z, z') < \varepsilon^{\alpha/2}$. In this case, we use (2.36) and (2.39) to deduce that

$$\sum_{n_1 \geq n_2} \tilde{\varrho}_\varepsilon^{n_1, n_2}(z, z') \lesssim 2^{n_2} \varepsilon^{\alpha/2} \wedge 2^{-n_2} \varepsilon^{-\alpha/2},$$

so that in this case $\tilde{\varrho}_\varepsilon(z, z') \lesssim 1 + (1 + t + t') \varepsilon^{\alpha/2}$.

It remains to consider the case $d_p(z, z') \geq \varepsilon^{\alpha/2}$. For this, we break the sum over n_2 in three pieces:

$$\begin{aligned} N_1 &= \{n_2 \geq 1 : 2^{-n_2} \geq d(z, z')/8\}, \\ N_2 &= \{n_2 \geq 1 : 2^{-6} \varepsilon^{\alpha/2} \leq 2^{-n_2} < d(z, z')/8\}, \\ N_3 &= \{n_2 \geq 1 : 2^{-n_2} < 2^{-6} \varepsilon^{\alpha/2}\}. \end{aligned}$$

For $n_2 \in N_1$, we only make use of the bound (2.36). Summing first over $n_1 \geq n_2$ and then over $n_2 \in N_1$, we obtain

$$\sum_{n_2 \in N_1} \sum_{n_1 \geq n_2} \tilde{\varrho}_\varepsilon^{n_1, n_2}(z, z') \lesssim \frac{\varepsilon^{\alpha/2}}{d_p(z, z')}.$$

For $n_2 \in N_2$, we only make use of the bound (2.37). Summing again first over $n_1 \geq n_2$ and then over $n_2 \in N_1$, we obtain

$$\sum_{n_2 \in N_2} \sum_{n_1 \geq n_2} \tilde{\varrho}_\varepsilon^{n_1, n_2}(z, z') \lesssim \frac{\varepsilon^{\alpha\gamma/2}}{d_p^\gamma(z, z')}.$$

In the last case, we similarly use either (2.40) or (2.41), depending on whether $|x - x'| \geq d_p(z, z')/2$ or $|t - t'| \geq d_p^2(z, z')/2$, which yields again

$$\sum_{n_2 \in N_3} \sum_{n_1 \geq n_2} \tilde{\varrho}_\varepsilon^{n_1, n_2}(z, z') \lesssim \frac{\varepsilon^{\alpha\gamma/2}}{d_p^\gamma(z, z')}.$$

Combining the above bounds, the claim follows. \square

Lemma 2.15 *For any $T > 0$, $p \geq 1$, $\kappa > 0$, $0 < \beta < 1$, there exists a constant $C_{T,p,\kappa,\beta}$ such that for all $0 \leq t \leq T$, $\varepsilon > 0$,*

$$(\mathbf{E}\|\partial_x Z^\varepsilon(t)\|_{\beta,p\kappa}^p)^{1/p} \leq C_{T,p,\kappa,\beta}\varepsilon^{-\kappa}.$$

Proof. This is a corollary of Lemma 2.11 and Proposition 2.2. \square

As a Corollary we deduce

Corollary 2.16 *For any $T > 0$, $p \geq 1$, $\kappa > 0$, there exists a constant $C_{T,\kappa}$ such that for all $0 \leq t \leq T$, $\varepsilon > 0$,*

$$|\mathbf{E}\|Z^\varepsilon(t)\|_{0,p\kappa}^p|^{1/p} \leq C_{T,\kappa}\varepsilon^{\alpha/2-\kappa}, \quad |\mathbf{E}\|\partial_x Z^\varepsilon(t)\|_{0,p\kappa}^p|^{1/p} \leq C_{T,\kappa}\varepsilon^{\alpha/4-\kappa}.$$

We will need moreover

Corollary 2.17 *As $\varepsilon \rightarrow 0$, $Z^\varepsilon(x, t) \rightarrow 0$ in probability, locally uniformly in (x, t) .*

Proof. It follows from estimate (2.16) that for any $p > 1$ and any bounded subset $K \subset \mathbb{R} \times \mathbb{R}^+$, there exists a constant $C_{p,K}$ such that

$$\mathbf{E}\left(\int_K (|\partial_x Y^\varepsilon(x, t)|^2 - \bar{V}^\varepsilon)^p dx dt\right) \leq C_{p,K}.$$

Then, by the Nash estimate, we obtain

$$\mathbf{E}\|Z^\varepsilon\|_{C^\gamma(K)} \leq C_K, \tag{2.42}$$

where the Hölder exponent $\gamma > 0$ and C_K do not depend on ε . As a consequence of the first estimate of Corollary 2.16 we have

$$\mathbf{E}\|Z^\varepsilon\|_{L^p(K)}^p \leq C_{p,K}\varepsilon^{p(\alpha/2-\kappa)}. \tag{2.43}$$

Combining (2.42) and (2.43) one can easily derive the required convergence. \square

3 Proof of the main result

Before concluding with the proof of our main theorem, we prove a result for a parabolic heat equation with coefficients which live in spaces of weighted Hölder continuous functions.

We consider an abstract evolution equation of the type

$$\partial_t u = \partial_x^2 u + F \partial_x u + G u , \quad (3.1)$$

where F and G are measurable functions of time, taking values in $\mathcal{C}_{p\kappa}^{-\beta}$ for some suitable $\kappa > 0$ and $\beta < \frac{1}{2}$. The main result of this section is the following:

Theorem 3.1 *Let β and κ be positive numbers such that $\beta + \kappa < \frac{1}{2}$ and let F and G be functions in $L_{\text{loc}}^p(\mathbf{R}_+, \mathcal{C}_{p\kappa}^{-\beta})$ for every $p \geq 1$.*

Let furthermore $\ell \in \mathbf{R}$ and $u_0 \in \mathcal{C}_{e_\ell}^{3/2}$. Then, there exists a unique global mild solution to (3.1). Furthermore, this solution is continuous with values in $\mathcal{C}_{e_m}^{3/2}$ for every $m < \ell$ and the map $(u_0, F, G) \mapsto u$ is jointly continuous in these topologies.

Proof. We will show a slightly stronger statement, namely that for every $\delta > 0$ sufficiently small, the mild solution has the property that $u_t \in \mathcal{C}_{e_{\ell-\delta t}}^\gamma$ for $t \in [0, T]$ for arbitrary values of $T > 0$. We fix T, δ and ℓ from now on.

We then write

$$\|u\|_{\delta, \ell, T} \stackrel{\text{def}}{=} \sup_{t \in [0, T]} \|u_t\|_{\frac{3}{2}, e_{\ell-\delta t}} ,$$

and we denote by $\mathcal{B}_{\delta, \ell, T}$ the corresponding Banach space. With this notation at hand, we define a map $\mathcal{M}_T: \mathcal{B}_{\delta, \ell, T} \rightarrow \mathcal{B}_{\delta, \ell, T}$ by

$$(\mathcal{M}_T u)_t = \int_0^t P_{t-s} (F_s \partial_x u_s + G_s u_s) ds , \quad t \in [0, T] .$$

It follows from Proposition 2.2 that we have the bound

$$\|(\mathcal{M}_T u)_t\|_{\frac{3}{2}, e_{\ell-\delta t}} \leq C \int_0^t (t-s)^{-\frac{3+2\beta}{4}} \|F_s \partial_x u_s + G_s u_s\|_{-\beta, e_{\ell-\delta t}} ds .$$

Combining Proposition 2.1 with (2.3) and (2.4), we furthermore obtain the bound

$$\begin{aligned} \|F_s \partial_x u_s\|_{-\beta, e_{\ell-\delta t}} &\leq C (\delta |t-s|)^{-\kappa} \|F_s\|_{-\beta, p\kappa} \|\partial_x u_s\|_{\frac{1}{2}, e_{\ell-\delta s}} \\ &\leq C (\delta |t-s|)^{-\kappa} \|F_s\|_{-\beta, p\kappa} \|u\|_{\delta, \ell, T} , \end{aligned}$$

where the proportionality constant C is uniformly bounded for $\delta \in (0, 1]$ and bounded ℓ and s . A similar bound holds for $G_s u_s$ so that, combining these bounds and using Hölder's inequality for the integral over t , we obtain the existence of constants $\zeta > 0$ and $p > 1$ such that the bound

$$\|\mathcal{M}_T u\|_{\delta, \ell, T} \leq C \delta^{-\kappa} T^\zeta (\|F\|_{L^p(\mathcal{C}_{p\kappa}^{-\beta})} + \|G\|_{L^p(\mathcal{C}_{p\kappa}^{-\beta})}) \|u\|_{\delta, \ell, T} ,$$

holds. Since the norm of this operator is strictly less than 1 provided that T is small enough, the short-time existence and uniqueness of solutions follow from Banach's fixed point theorem. The existence of solutions up to the final time T follows by iterating this argument, noting that the interval of short-time existence restarting from $u(t)$ at time t can be bounded from below by a constant that is uniform over all $t \in [0, T]$, as a consequence of the linearity of the equation.

Actually, we obtain the bound

$$\|u_t\|_{\frac{3}{2}, e_{\ell-\delta t}} \lesssim \exp(Ct(\|F\|_{L^p(\mathcal{C}_{p\kappa}^{-\beta})} + \|G\|_{L^p(\mathcal{C}_{p\kappa}^{-\beta})})^{1/\zeta}) \|u_0\|_{\frac{3}{2}, e_\ell} ,$$

where the constants C and ζ depend on the choice of ℓ and δ .

The solutions are obviously linear in u_0 since the equation is linear in u . It remains to show that the solutions also depend continuously on F and G . Let \bar{u} be the solution to the equation

$$\partial_t \bar{u} = \partial_x^2 \bar{u} + \bar{F} \partial_x \bar{u} + \bar{G} \bar{u} , \quad (3.2)$$

and write $\varrho = u - \bar{u}$. The difference ϱ then satisfies the equation

$$\partial_t \varrho = \partial_x^2 \varrho + F \partial_x \varrho + G \varrho + (F - \bar{F}) \partial_x \bar{u} + (G - \bar{G}) \bar{u} ,$$

with zero initial condition. Similarly to before, we thus have

$$\varrho_t = (\mathcal{M}_T \varrho)_t + \int_0^t P_{t-s} ((F_s - \bar{F}_s) \partial_x \bar{u}_s + (G_s - \bar{G}_s) \bar{u}_s) ds .$$

It follows from the above bounds that

$$\|\varrho\|_{\delta, \ell, T} \lesssim \|\mathcal{M}_T \varrho\|_{\delta, \ell, T} + C \delta^{-\kappa} T^\zeta (\|F - \bar{F}\|_{L^p(\mathcal{C}_{p\kappa}^{-\beta})} + \|G - \bar{G}\|_{L^p(\mathcal{C}_{p\kappa}^{-\beta})}) \|\bar{u}\|_{\delta, \ell, T} .$$

Over short times, the required continuity statement thus follows at once. Over fixed times, it follows as before by iterating the argument. \square

Remark 3.2 In principle, one could obtain a similar result for less regular initial conditions, but this does not seem worth the additional effort in this context.

We now have finally all the ingredients in place to give the proof of our main result.

Proof of Theorem 1.8. We apply Theorem 3.1 with $\beta = \frac{1}{4}$ and $\kappa = \frac{1}{8}$. Note that the equation (1.8) for v^ε is precisely of the form (3.1) with

$$F = 2\partial_x Y^\varepsilon + 2\partial_x Z^\varepsilon, \quad G = |\partial_x Z^\varepsilon|^2 + 2\partial_x Z^\varepsilon \partial_x Y^\varepsilon.$$

It follows from Corollaries 2.12 and 2.16 that, for every $p > 0$ and $\delta > 0$, one has the bound

$$\left| \mathbf{E} \int_0^T \|F\|_{\beta, p_\kappa}^p dt \right|^{1/p} \lesssim \varepsilon^{\frac{\alpha}{16} - \delta},$$

say. Similarly, it follows from Lemma 2.11 and Corollary 2.16 that one actually has the bound

$$\left| \mathbf{E} \int_0^T \|G\|_{0, p_\kappa}^p dt \right|^{1/p} \lesssim \varepsilon^{\frac{\alpha}{4} - \delta},$$

which is stronger than what we required. As a consequence of Theorem 3.1, this shows immediately that $v^\varepsilon \rightarrow u$ in probability, locally uniformly both in space and in time. We conclude by recalling that from Corollary 2.9 and 2.17, the correctors Y^ε and Z^ε themselves converge locally uniformly to 0 in probability. \square

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